

Recent measurements of CP violation

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The Mystery of the Matter Asymmetry

- According to the Big Bang theory, equal amounts of matter and antimatter were initially created in the early Universe. . However, in the current Universe, almost all things consist of only matter.
- CP violation is required to explain the matter- antimatter asymmetry of the Universe.
- The study of CP violation (CPV) is experiment-driven. CP violation was discovered in 1964 in the decays of neutral kaons, which was completely unexpected then.
- To date, CPV has been confirmed in K , B and D meson decays.
- In the Standard Model (SM), CPV is present in the weak interactions via CKM mechanism, it is insufficient to explain the observed matter-antimatter asymmetry in the Universe.
- The SM is incomplete and other sources of CPV must exist.



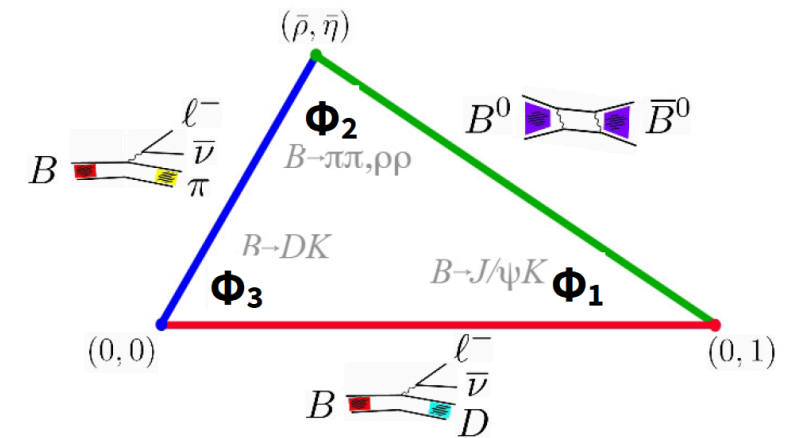
CP violation and the Unitarity Triangle

- In SM, the Kobayashi-Maskawa mechanism provides an elegant and simple explanation of CP violation.
- Wolfenstein parametrization: CKM matrix described by 4 parameters λ, A, ρ, η .

$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$

$$\approx \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$



- η is the source of CPV in the SM.
- Key test of the SM: determining accurately the Unitarity Triangle
 - Magnitudes: branching fractions or mixing frequencies
 - Angles $(\phi_1/\beta, \phi_2/\alpha, \phi_3/\gamma)$: CP violation measurements in B decays.

Recent CPV measurements



Time-dependent:

- $B^0 \rightarrow J/\psi K_S^0$ [arXiv:2302.12898](#)
- $B^0 \rightarrow \phi K_S^0$ Moriond 2023
- $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ Moriond 2023
- $B^0 \rightarrow K_S^0 \pi^0$ [arXiv:2305.07555](#)

Time-integrated:

- $B^+ \rightarrow \pi^+ \pi^0$ Moriond 2023
- $B^0 \rightarrow \pi^0 \pi^0$ [arXiv:2303.08354](#)

γ/ϕ_3 measurement :

- $B^+ \rightarrow D(K_S^0 K^+ \pi^-) K^+$ [arXiv: 2306.02940](#)
- $B^+ \rightarrow D_{CP} K^+$ Moriond 2023

Charm

- T-odd CPV in $D_{(s)}^+$ four-body decays
[PRD 107, 052001 \(2023\)](#)
[arXiv:2305.12806](#) [arXiv:2305.11405](#)
- $\Lambda_c^+ \rightarrow \Lambda h^+$ and $\Lambda_c^+ \rightarrow \Xi^0 h^+$
[Sci.Bull. 68 \(2023\) 583](#)

- Recent CPV measurements in the quark sector.

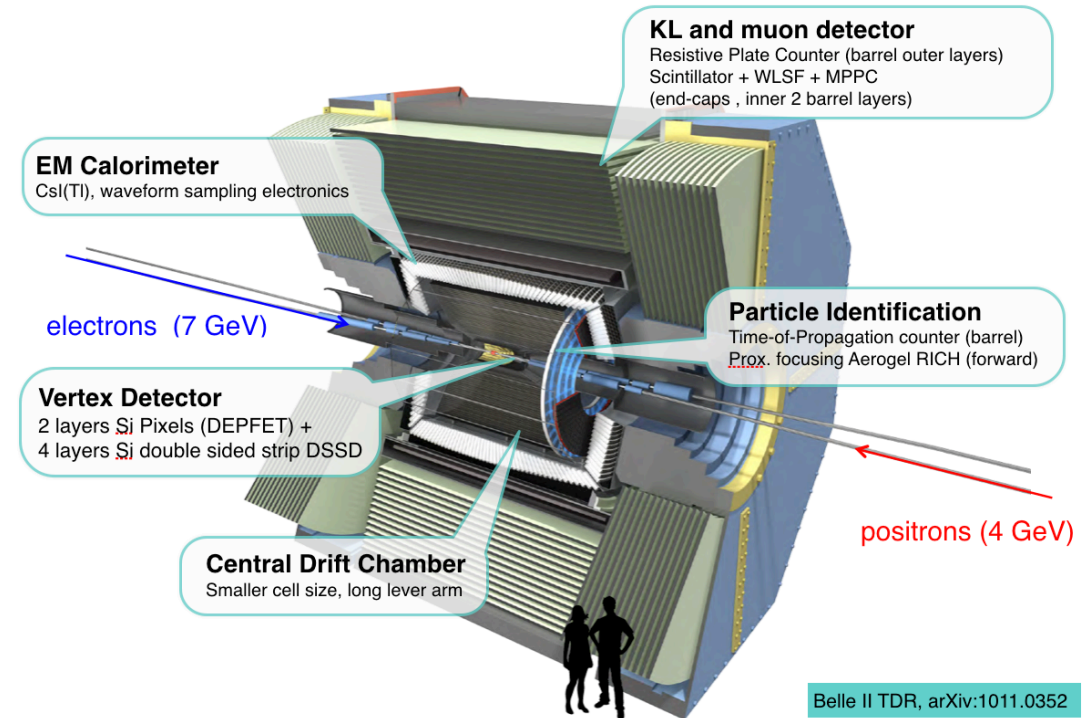
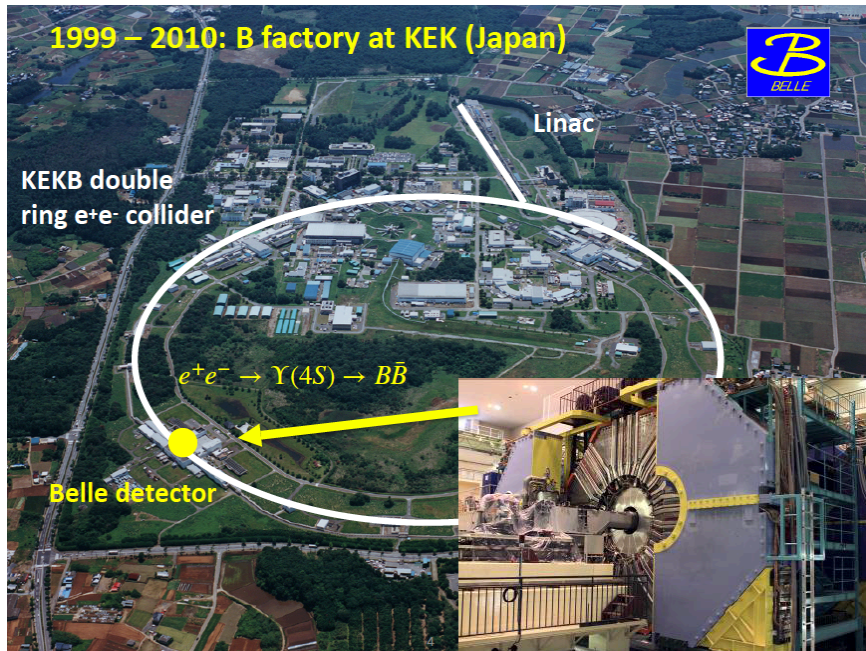


- $B^0 \rightarrow \psi K_S^0$ LHCb-Paper-2023-013
- $B^+ \rightarrow D(K^- \pi^+ \pi^+ \pi^-) h^+$
[arXiv:2209.03692](#)
- $B^+ \rightarrow D(h^+ h^- \pi^+ \pi^-) h^+$
[arXiv:2301.10328](#)
- $B_S^0 \rightarrow J/\psi K^+ K^-$ LHCb-Paper-2023-016
- $B_S^0 \rightarrow \phi \phi$ [arXiv:2304.06198](#)
- $D^0 \rightarrow K^+ K^-$ [arXiv:2209.03179](#)
- $D_{(s)} \rightarrow K^- K^+ K^+$ [arXiv:2303.04062](#)
- $D^0 \rightarrow \pi^- \pi^+ \pi^0$ [arXiv:2306.12746](#)



- $J/\psi \rightarrow \Xi^- \Xi^+$
- [Nature 606, 64 \(2022\)](#)
- $J/\psi \rightarrow \Sigma^+ \bar{\Sigma}^-$
- [arXiv:2304.14655](#)
- $J/\psi \rightarrow \Xi^0 \bar{\Xi}^0$
- [arXiv:2305.09218](#)
- $\psi(3686) \rightarrow \Xi^0 \bar{\Xi}^0$
- [arXiv:2302.09767](#)
- $J/\psi \rightarrow \Lambda \bar{\Lambda}$
- [PRL 129 \(2022\) 131801](#)

Belle & Belle II Experiments

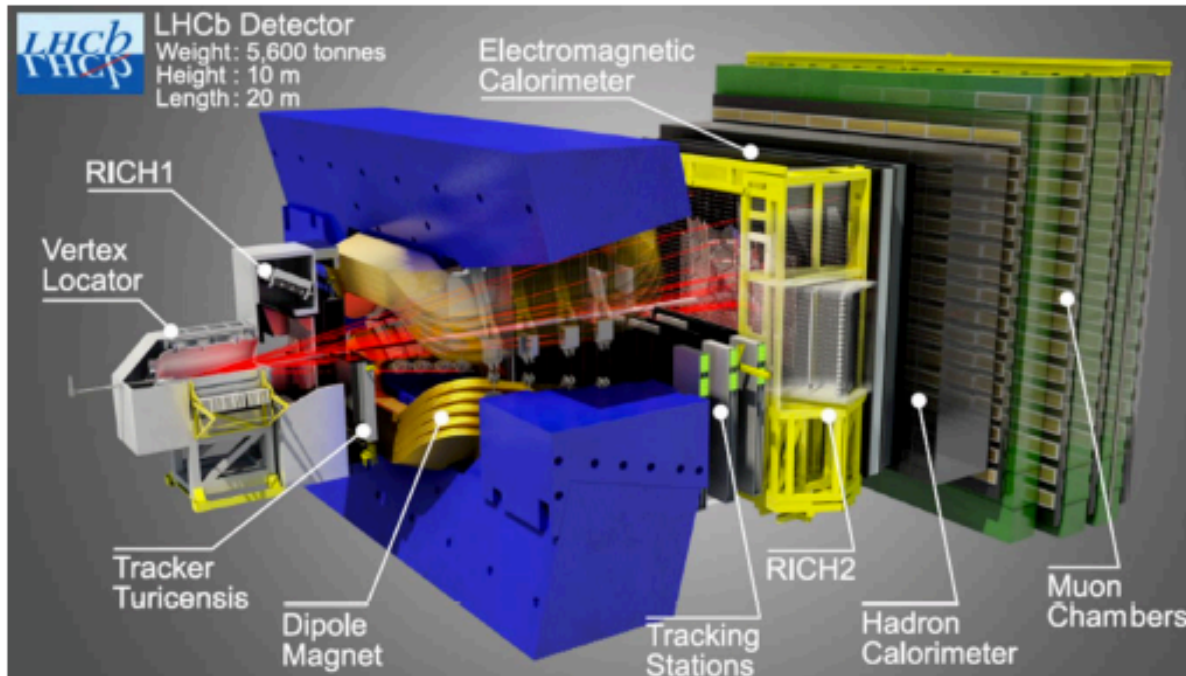


- Asymmetric e^+e^- colliders. B factories, also charm and τ factories.
- Belle \rightarrow Belle II: e^+ (3.5 GeV) e^- (8 GeV) $\rightarrow e^+$ (4 GeV) e^- (7 GeV). Improved vertex resolution allows lower boost.
- 428 fb^{-1} (362 fb^{-1} at $\Upsilon(4S)$) collected at Belle II so far. Goal: 50 ab^{-1} .
- Record peak luminosity $4.7 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ at Belle II.

- Silicon detectors for vertex measurement
- Drift chamber for p and dE/dx measurement
- TOP and ARICH counters for particle identification
- EM Calorimeter (ECL) for e^+ and γ reconstruction
- KLM detector for K_L and μ detection

The LHCb Experiment

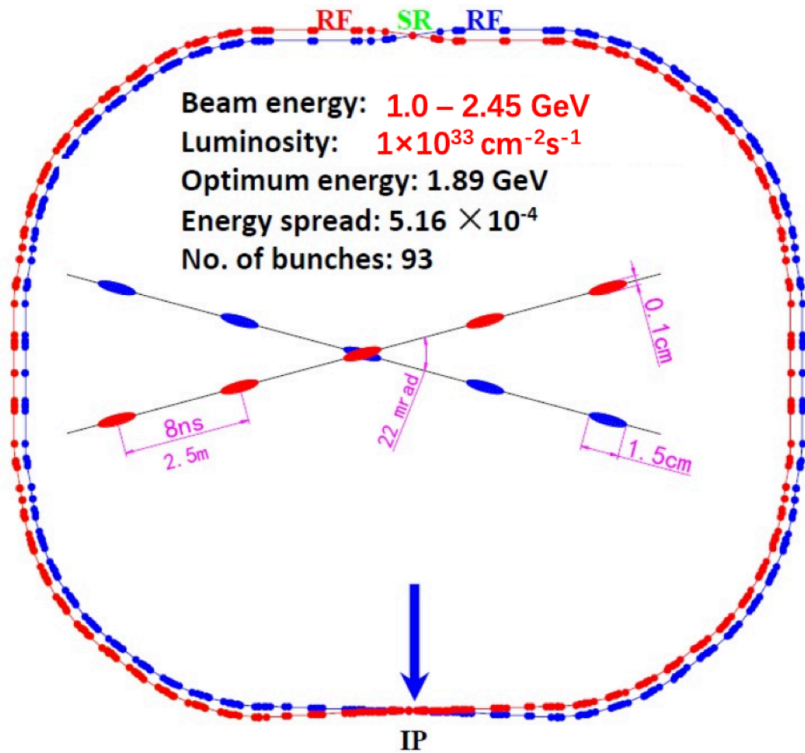
LHCb @ LHC (CERN): pp collisions



Run 1-2 detector

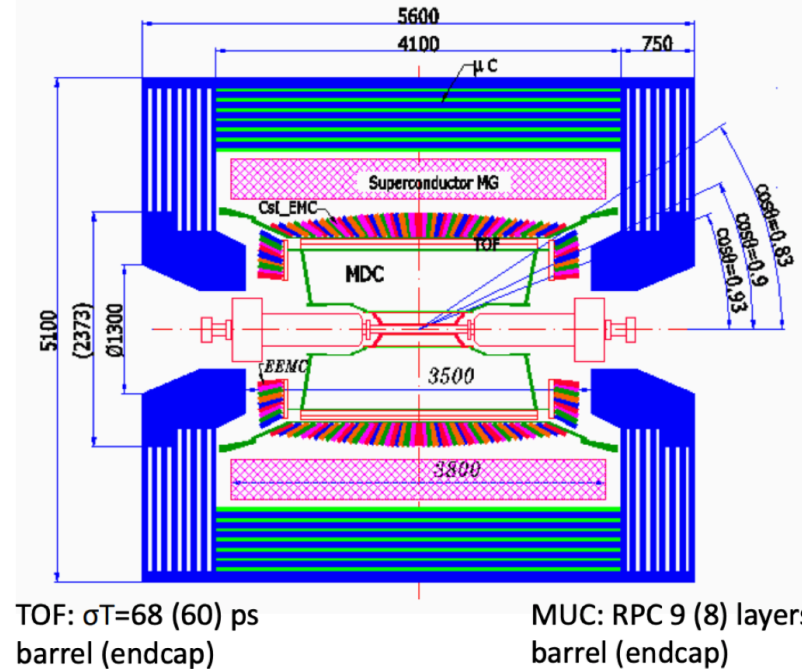
- Forward detector optimized for b and c meson studies.
- Huge b cross-section.
- Excellent vertex resolution and particle identification.
- Events with high multiplicity, reconstruction of neutrals is challenging.
- 9 fb^{-1} accumulated during Run1-2 (2010-2018)
- Run 3 started in 2022 with an upgraded LHCb detector, goal 50 fb^{-1} .

The BESIII Experiment



MDC: $\Delta p/p = 0.5\%$ @ 1GeV/c
 $dE/dx: \sim 6\%$

EMC: CsI (TI) 2.5% (5.0%)
 barrel (endcap) @ 1GeV



- Symmetric e^+e^- collider, low background. Charm and τ factory.
- Multipurpose detector with good resolution, near 4π angular coverage.
- World's largest charmonia data sample.

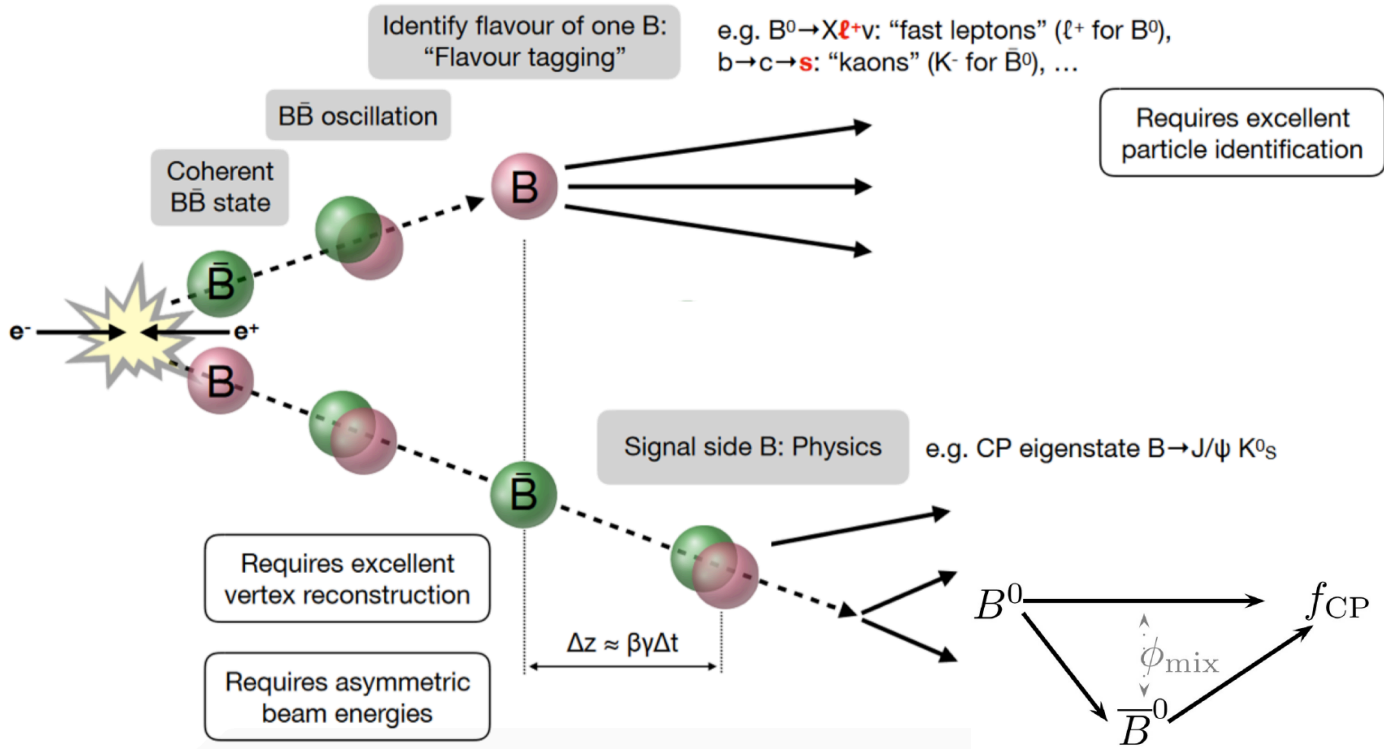
Time Dependent (TD) CPV measurement at Belle(II)

- The CP asymmetry, $A_{CP}(\Delta t)$, of the proper time difference distribution Δt of the two B mesons

$$A_{CP}(\Delta t) = \frac{\mathcal{B}(\bar{B}^0 \rightarrow f_{CP})(\Delta t) - \mathcal{B}(B^0 \rightarrow f_{CP})(\Delta t)}{\mathcal{B}(\bar{B}^0 \rightarrow f_{CP})(\Delta t) + \mathcal{B}(B^0 \rightarrow f_{CP})(\Delta t)} = S \sin(\Delta m_d \Delta t) + A \cos(\Delta m_d \Delta t)$$

The SM predicts $A = 0$ and $S = -\eta \sin 2\phi_1$

Mixing-induced CPV Direct CPV



- The $B^0\bar{B}^0$ pairs from $\Upsilon(4S)$ are produced in a coherent, entangled quantum mechanical state. Tagging the flavor the other side B allows to know the flavor of signal side B.
- The distance between signal and tag-side B decay vertices in z direction (Δz) allows Δt to be measured.

TD CPV in $B^0 \rightarrow J/\psi K_S^0$

- Tree-dominated $b \rightarrow c\bar{c}s$, golden mode
 - small theoretical uncertainty and clean experimental signature.
 - contribution of penguin diagrams with a different CKM phase is expected to be at less than 1% level.
- Fit ΔE distribution and subtract background.
- Fit the sWeighted Δt distribution to extract $\sin(2\beta)$ and A_{CP} .

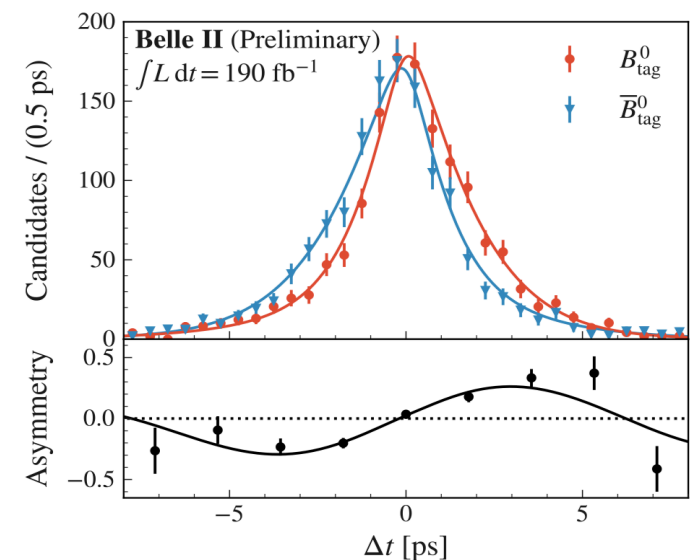
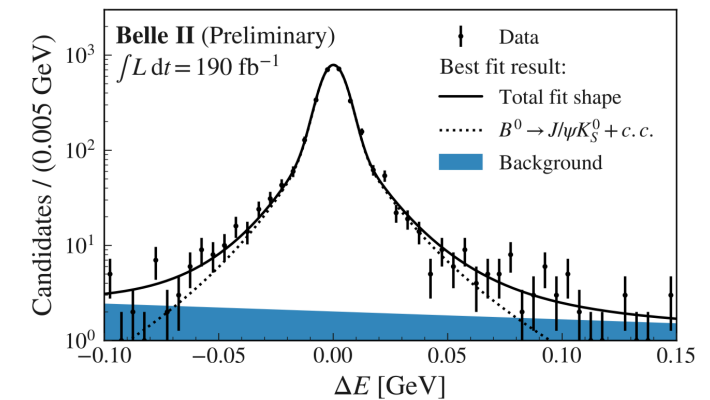
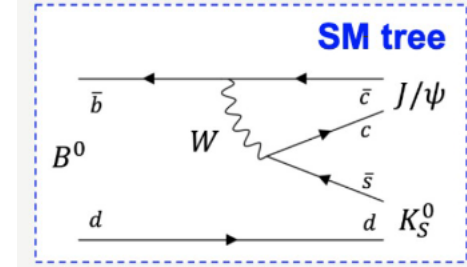
$$\mathcal{A}(\Delta t) = S_{J/\psi K_S^0} \sin(\Delta m_d \Delta t) + A_{J/\psi K_S^0} \cos(\Delta m_d \Delta t)$$

$$S_{CP} = 0.720 \pm 0.062(\text{stat}) \pm 0.016(\text{syst}),$$

$$A_{CP} = 0.094 \pm 0.044(\text{stat}) \begin{matrix} +0.042 \\ -0.017 \end{matrix}(\text{syst}),$$

- Consistent with the world-average. Statistical uncertainty is twice that of the current most precise determination (from Belle), consistent with a four-times smaller data set.

$$S_{J/\psi K_S^0} = \sin 2\phi_1$$

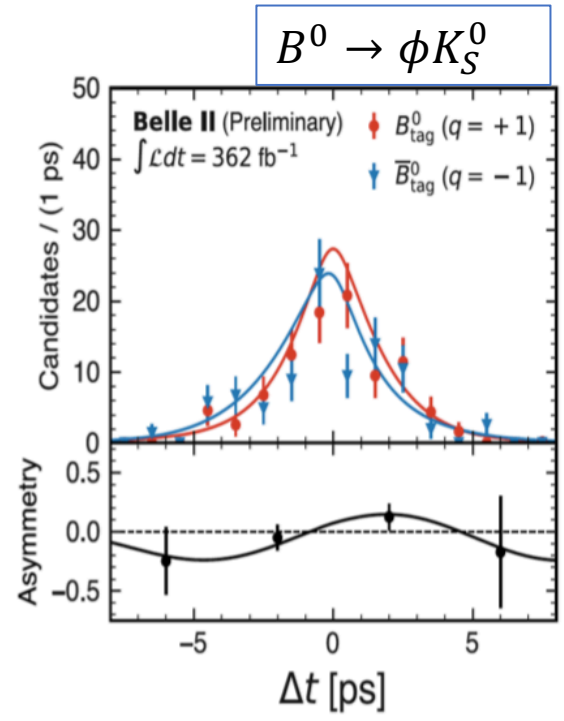
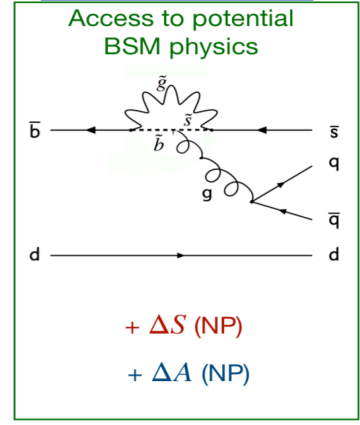
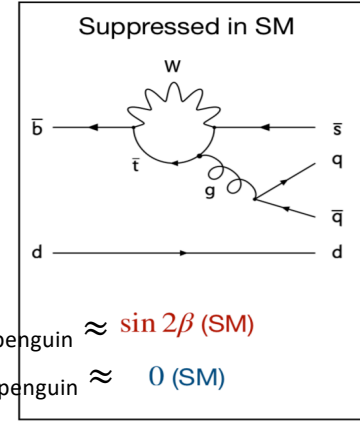


TD CPV in $b \rightarrow sq\bar{q}$

- Penguin-dominated modes $b \rightarrow sq\bar{q}$: $B^0 \rightarrow \phi K_S^0$ and $B^0 \rightarrow K_S^0 K_S^0 K_S^0$.

$$\mathcal{A}(\Delta t) = S_{sq\bar{q}} \sin(\Delta m_d \Delta t) + A_{sq\bar{q}} \cos(\Delta m_d \Delta t)$$

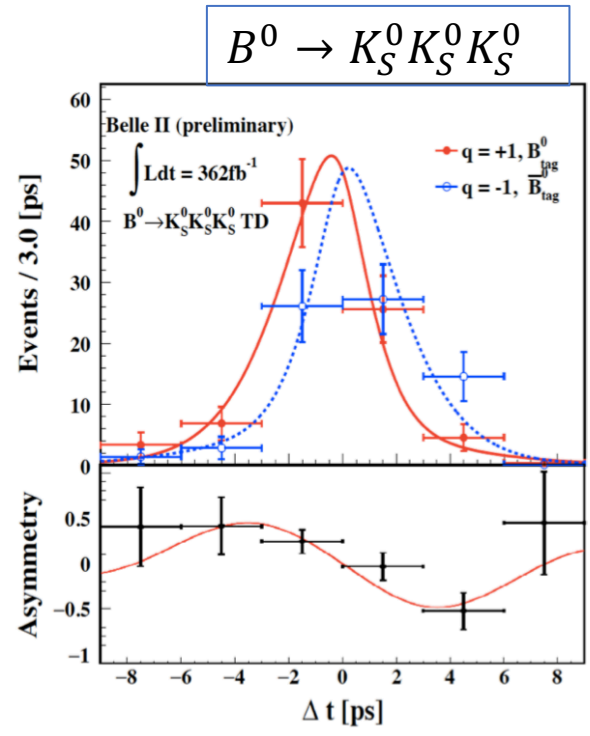
- $S_{sq\bar{q}} \approx -\eta_{CP} S_{J/\psi K_S^0}$ in the SM; discrepancy can arise if NP particle propagated in the $b \rightarrow sq\bar{q}$ loop.



$S_{K_S\phi} = +0.54 \pm 0.26^{+0.06}_{-0.08}$
 $A_{K_S\phi} = +0.31 \pm 0.20^{+0.05}_{-0.06}$

- $S_{K_S\phi}$ is consistent with $-\eta_{CP} S_{J/\psi K_S^0}$
- Belle results:

535M $B\bar{B}$ pairs
 $S_{K_S\phi} = 0.50 \pm 0.21 \pm 0.06$
657M $B\bar{B}$ pairs
 $A_{K_S\phi} = 0.04 \pm 0.20 \pm 0.10$



- Challenge: no prompt tracks from vertex
- Decay vertex reconstruction using K_S^0 trajectory and profile of the interaction point.

$S_{3K_S} = -1.37^{+0.35}_{-0.45} \pm 0.03$
 $A_{3K_S} = +0.07^{+0.15}_{-0.20} \pm 0.02$

- S_{3K_S} is consistent with $-\eta_{CP} S_{J/\psi K_S^0}$
- Belle (772M $B\bar{B}$ pairs) :
 $S_{3K_S} = -0.71 \pm 0.23 \pm 0.05$
 $A_{3K_S} = 0.12 \pm 0.16 \pm 0.05$

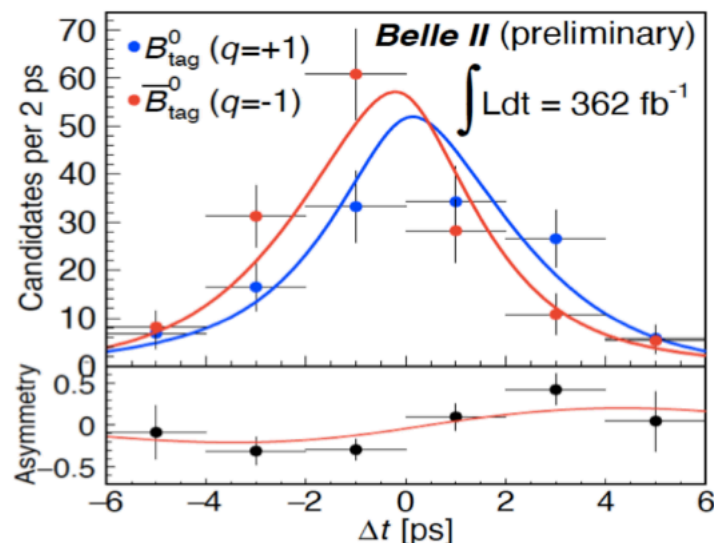
CPV in $B^0 \rightarrow K_S^0 \pi^0$

- $B^0 \rightarrow K_S^0 \pi^0$ proceeds mainly via $b \rightarrow sd\bar{d}$ loop amplitude.
- Sensitive to effective value of $\sin 2\phi_1$ and providing inputs to isospin sum-rule

$$I_{K\pi} = \text{Br}(K^+\pi^-) \mathcal{A}_{K^+\pi^-} + \mathcal{A}_{K^0\pi^+} \text{Br}(K^0\pi^+) \frac{\tau_{B^0}}{\tau_{B^+}} - 2\mathcal{A}_{K^+\pi^0} \text{Br}(K^+\pi^0) \frac{\tau_{B^0}}{\tau_{B^+}} - 2\mathcal{A}_{K^0\pi^0} \text{Br}(K^0\pi^0); \quad I_{K\pi}^{SM} = 0$$

M. Gronau, Phys. Lett. B 627, 82(2005).

- Time dependent $B^0 \rightarrow K_S^0 \pi^0$: main challenge is the decay vertex reconstruction.
- Validated on $B^0 \rightarrow J/\psi K_S^0$ events reconstructed w/o J/ψ vertex.
- Simultaneous time-integrated (40% events without Δt) and time-dependent fit to maximize the sensitivity on A_{CP} .
- Competitive with world's best results with less (60-80%) luminosity.



$$A_{CP} = 0.04 \pm 0.15 \pm 0.05 \quad A_{CP}^{w.a.} = 0.00 \pm 0.13$$

$$S_{CP} = 0.75^{+0.20}_{-0.23} \pm 0.04 \quad S_{CP}^{w.a.} = 0.58 \pm 0.17$$

combine time-integrated and time-dependent results:

$$A_{CP}^{K_S^0 \pi^0} = -0.01 \pm 0.12 \pm 0.05 \quad w.a. = -0.0 \pm 0.13$$

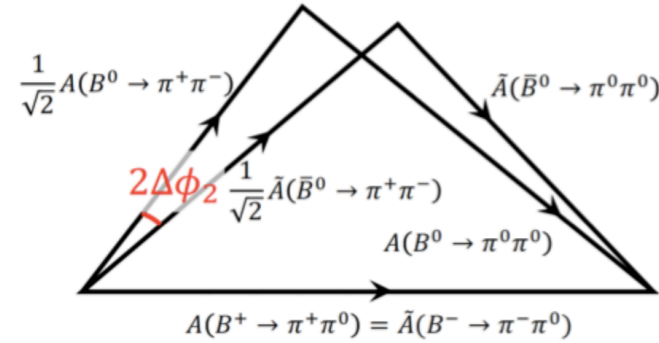
- Combining all $B \rightarrow K\pi$ final states at Belle II:

$$I_{K\pi} = -0.03 \pm 0.13 \pm 0.05 \quad w.a. = 0.13 \pm 0.11$$

CPV in $B^+ \rightarrow \pi^+ \pi^0$ and $B^0 \rightarrow \pi^0 \pi^0$

- The CKM angle ϕ_2 measurement is based on $b \rightarrow u$ processes. Significant contributions from penguins ($b \rightarrow d$): $\phi_2^{\text{eff}} = \phi_2 + \Delta\phi_2$.
- Penguin and tree contributions can be disentangled using the isospin relations

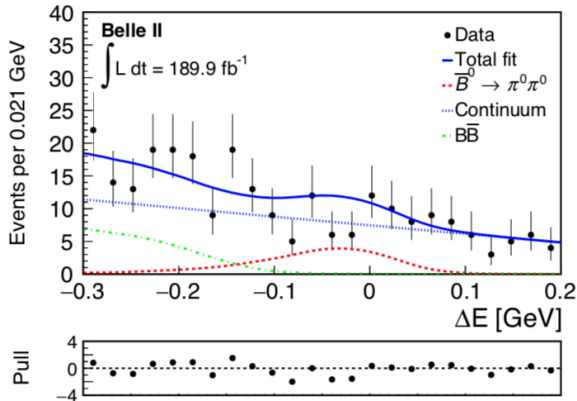
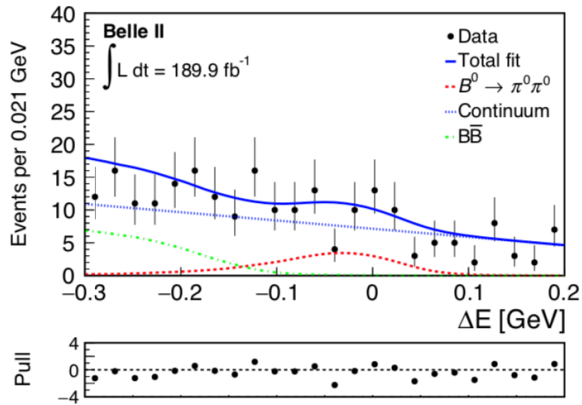
M. Gronau and D. London, Phys Rev. Lett. **65**, 3381 (1990).



$$A^{+0} = \frac{1}{\sqrt{2}} A^{+-} + A^{00} \quad \text{and} \quad \bar{A}^{-0} = \frac{1}{\sqrt{2}} \bar{A}^{+-} + \bar{A}^{00},$$

[arXiv:2303.08354](https://arxiv.org/abs/2303.08354)

Belle II 189 fb⁻¹

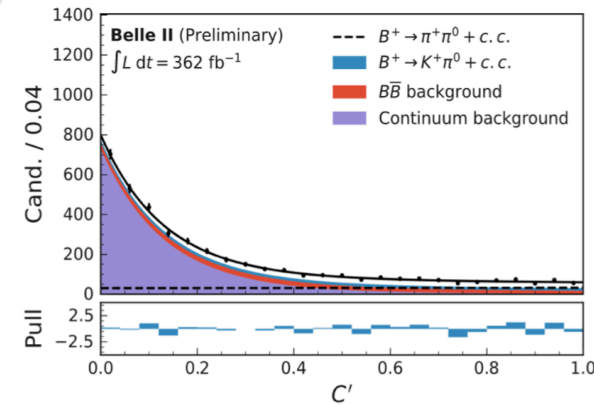
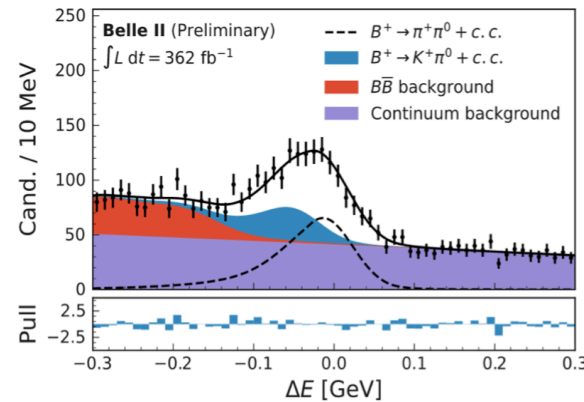


$$A_{CP}(B^0 \rightarrow \pi^0 \pi^0) = 0.14 \pm 0.46 \pm 0.07.$$

Belle (772M $B\bar{B}$): $A_{CP} = +0.14 \pm 0.36$ (stat.) ± 0.10 (syst.)

$B^+ \rightarrow \pi^+ \pi^0$

Belle II 362 fb⁻¹



$$A_{CP}(B^+ \rightarrow \pi^+ \pi^0) = -0.08 \pm 0.05 \pm 0.01$$

Belle (772M $B\bar{B}$): $A_{CP} = +0.025 \pm 0.043^{12} \pm 0.007$

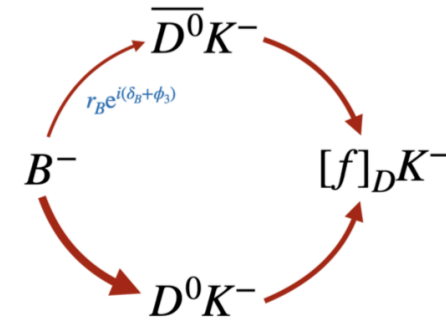
CPV in $B^+ \rightarrow DK^+$ (towards ϕ_3/γ)

- ϕ_3/γ is the phase between $b \rightarrow u$ and $b \rightarrow c$ transitions.
- The absence of the loop contribution allows extremely clean theoretical prediction of ϕ_3
- Various approaches with different D final states:
 - Self-conjugate final states $D^0 \rightarrow K_S h^+ h^-$
 - singly Cabibbo-suppressed $D^0 \rightarrow K_S K^+ \pi^-$.
 - CP eigenstates: $D^0 \rightarrow K^+ K^-$, $K_S^0 \pi^0$
 - CF and DCS decays: $D^0 \rightarrow K^+ \pi^-$

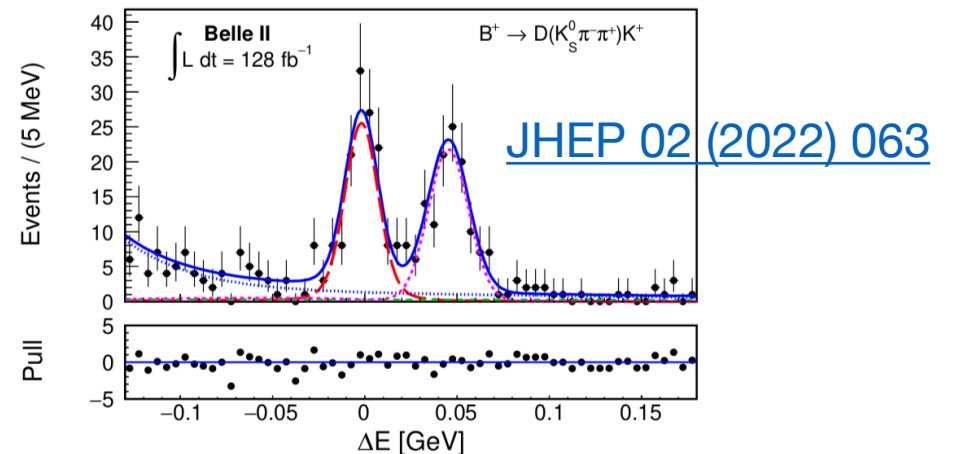
Current WA dominated by LHCb:

$$\gamma[^\circ] = 65.9 \begin{matrix} + 3.3 \\ - 3.5 \end{matrix} \quad \perp$$

Interference between two decays to same final state gives access to phase:



$$\frac{\mathcal{A}^{\text{suppr.}}(B^- \rightarrow \bar{D}^0 K^-)}{\mathcal{A}^{\text{favor.}}(B^- \rightarrow D^0 K^-)} = r_B e^{i(\delta_B + \gamma)}$$



$$\phi_3 = (78.4 \pm 11.4 \pm 0.5 \pm 1.0)^\circ$$

$$B^+ \rightarrow D(K_S^0 K^+ \pi^-)K^+ / \pi^+$$

Belle II 362 fb⁻¹
+ Belle 711 fb⁻¹

- $B^\pm \rightarrow D(K_S^0 K^\pm \pi^\mp)K^\pm / \pi^\pm$. SS: same sign, OS: opposite sign.
- 2D fit (ΔE , CS') of 8 categories: (+, -) x (SS, OS) x (DK, D π) in full D phase space and in interference-enhanced $D \rightarrow K^* K$ region.
- Model-independent information on the ϕ_3 .
- Results consistent with LHCb, but not competitive. Contribute to constrain γ/ϕ_3 in combination with other measurements.

$$A_{SS}^{DK} = -0.089 \pm 0.091 \pm 0.011,$$

$$A_{OS}^{DK} = 0.109 \pm 0.133 \pm 0.013,$$

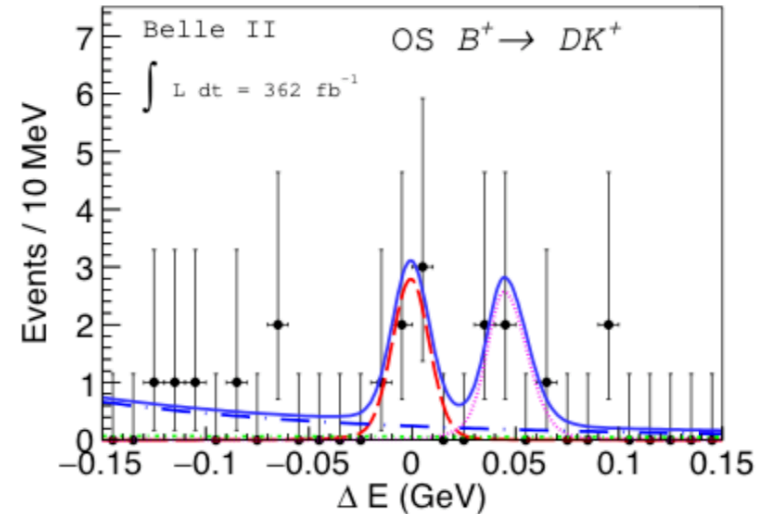
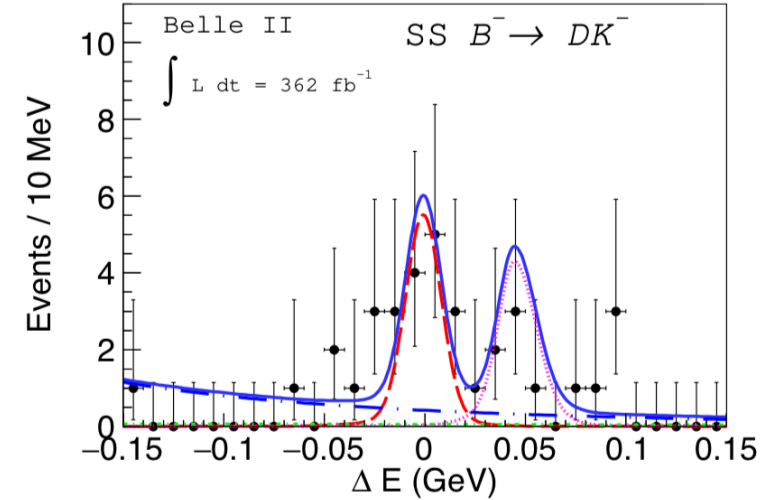
$$A_{SS}^{D\pi} = 0.018 \pm 0.026 \pm 0.009,$$

$$A_{OS}^{D\pi} = -0.028 \pm 0.031 \pm 0.009,$$

$$R_{SS}^{DK/D\pi} = 0.122 \pm 0.012 \pm 0.004,$$

$$R_{OS}^{DK/D\pi} = 0.093 \pm 0.013 \pm 0.003,$$

$$R_{SS/OS}^{D\pi} = 1.428 \pm 0.057 \pm 0.002,$$



$B^+ \rightarrow D_{CP} K^+$

Belle II 189 fb⁻¹
+ Belle 711 fb⁻¹

- $B^+ \rightarrow D_{CP} K^+, D \rightarrow K^+ K^-$ (CP even), $D \rightarrow K_S^0 \pi^0$ (CP odd).
- 2D fit ($\Delta E, CS'$) of 6 categories: $(DK, D\pi) \times (K^+ K^-, K_S^0 \pi^0, K^+ \pi^-)$.

$$- \mathcal{R}_{CP\pm} = \frac{Br(B^- \rightarrow D_{CP\pm} K^-) + Br(B^+ \rightarrow D_{CP\pm} K^+)}{Br(B^- \rightarrow D_{flav} K^-) + Br(B^+ \rightarrow D_{flav} K^+)}$$

$$= 1 + r_B^2 + 2r_B \cos \delta_B \cos \phi_3$$

$$- \mathcal{A}_{CP\pm} \equiv \frac{\Gamma(B^- \rightarrow D_{CP\pm} K^-) - \Gamma(B^+ \rightarrow D_{CP\pm} K^+)}{\Gamma(B^- \rightarrow D_{CP\pm} K^-) + \Gamma(B^+ \rightarrow D_{CP\pm} K^+)}$$

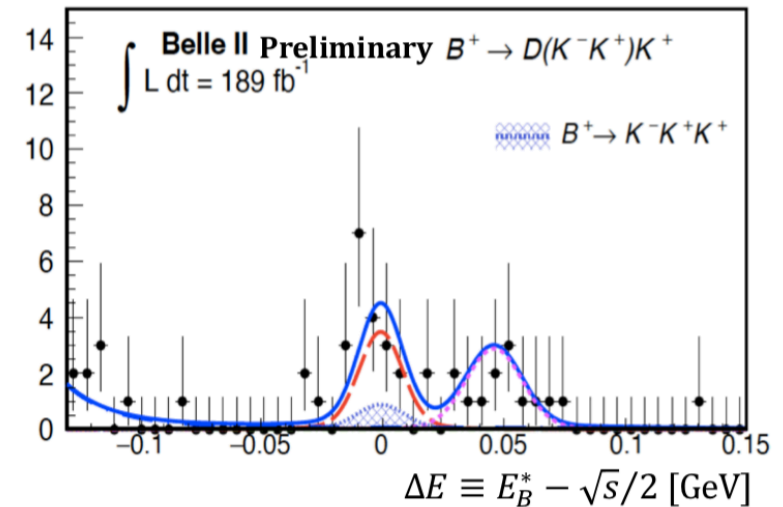
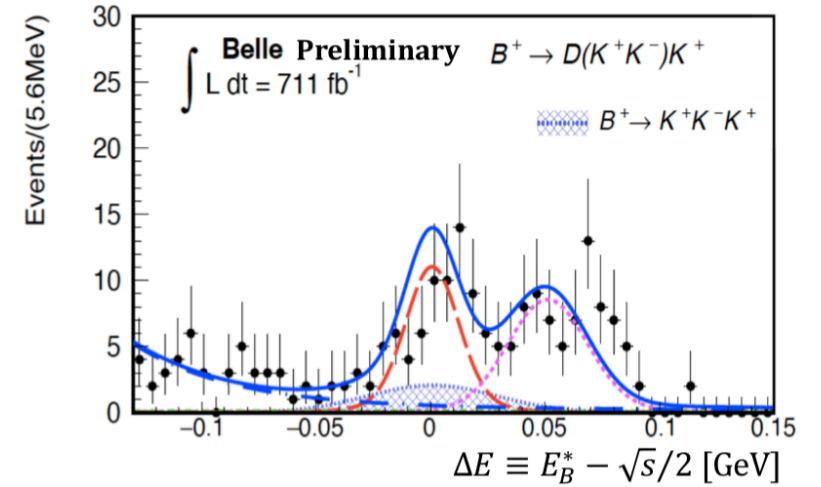
$$= \pm 2r_B \sin \delta_B \sin \phi_3 / \mathcal{R}_{CP\pm}$$

$$\mathcal{R}_{CP+} = 1.164 \pm 0.081 \pm 0.036 \quad \mathcal{R}_{CP-} = 1.151 \pm 0.074 \pm 0.019$$

$$\mathcal{A}_{CP+} = (+12.5 \pm 5.8 \pm 1.4)\% \quad \mathcal{A}_{CP-} = (-16.7 \pm 5.7 \pm 0.6)\%$$

... best \mathcal{A}_{CP-} measurement

$$4.7^\circ < \phi_3 < 175.8^\circ, \quad 0.069 < r_B < 0.560 \quad \dots @ 95.4\% \text{ CL}$$



TD CPV in $B^0 \rightarrow \psi K_S^0$

- Three modes with full Run 2 data
 - $B^0 \rightarrow J/\psi(\rightarrow \mu\mu)K_S^0(\rightarrow \pi^+\pi^-)$ (82%)
 - $B^0 \rightarrow J/\psi(\rightarrow ee)K_S^0(\rightarrow \pi^+\pi^-)$ (12%)
 - $B^0 \rightarrow \psi(2S)(\rightarrow \mu\mu)K_S^0(\rightarrow \pi^+\pi^-)$ (6%)
- From mass fit, *sWeights* are obtained for effective background subtraction in CP fit.

$$P_{B^0,(\bar{B}^0)}(t) \propto (1 \mp \alpha)(1 \mp \Delta\epsilon)(1 \mp \mathbf{S} \sin(\Delta m_d t) \pm \mathbf{C} \cos(\Delta m_d t))$$

- The most precise single measurement of $\sin(2\phi_1)$ to date.

Combined fit result

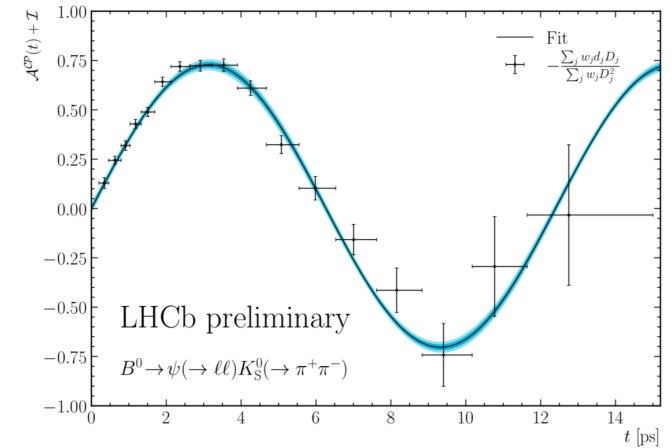
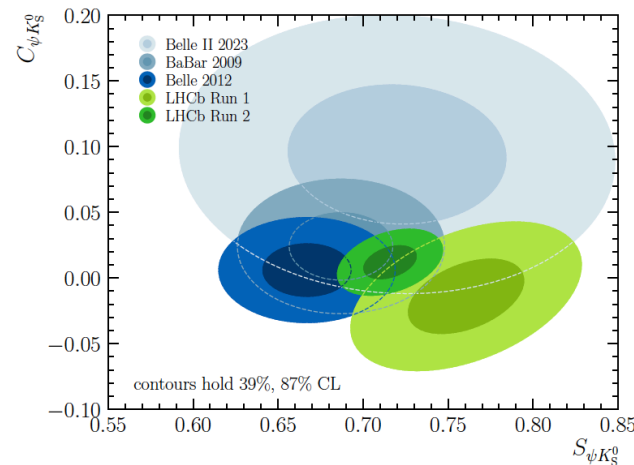
$$S_{\psi K_S^0}^{\text{Run 2}} = 0.716 \pm 0.013 (\text{stat}) \pm 0.008 (\text{syst})$$

$$C_{\psi K_S^0}^{\text{Run 2}} = 0.012 \pm 0.012 (\text{stat}) \pm 0.003 (\text{syst})$$

New total LHCb combination

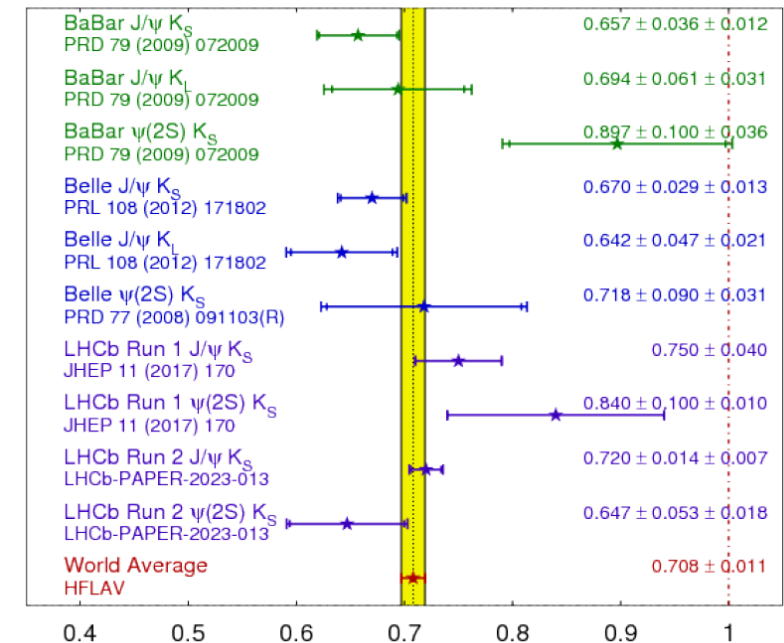
$$S_{\psi K_S^0}^{\text{Run 1+2}} = 0.723 \pm 0.014 (\text{stat+syst})$$

$$C_{\psi K_S^0}^{\text{Run 1+2}} = 0.007 \pm 0.012 (\text{stat+syst})$$



$$\sin(2\beta) \equiv \sin(2\phi_1)$$

HFLAV
Summer 2023
PRELIMINARY



$$B^+ \rightarrow D(K^- \pi^+ \pi^+ \pi^-) h^+$$

- Improved sensitivity can be achieved with binned measurement.

- 4 bins chosen according to LHCb amplitude analysis [EPJC 78 \(2018\) 6, 443](#)

- using external inputs for

- Hadronic decay parameters from model-independent determinations by CLEO-c and BESIII [JHEP 05 \(2021\) 164](#)

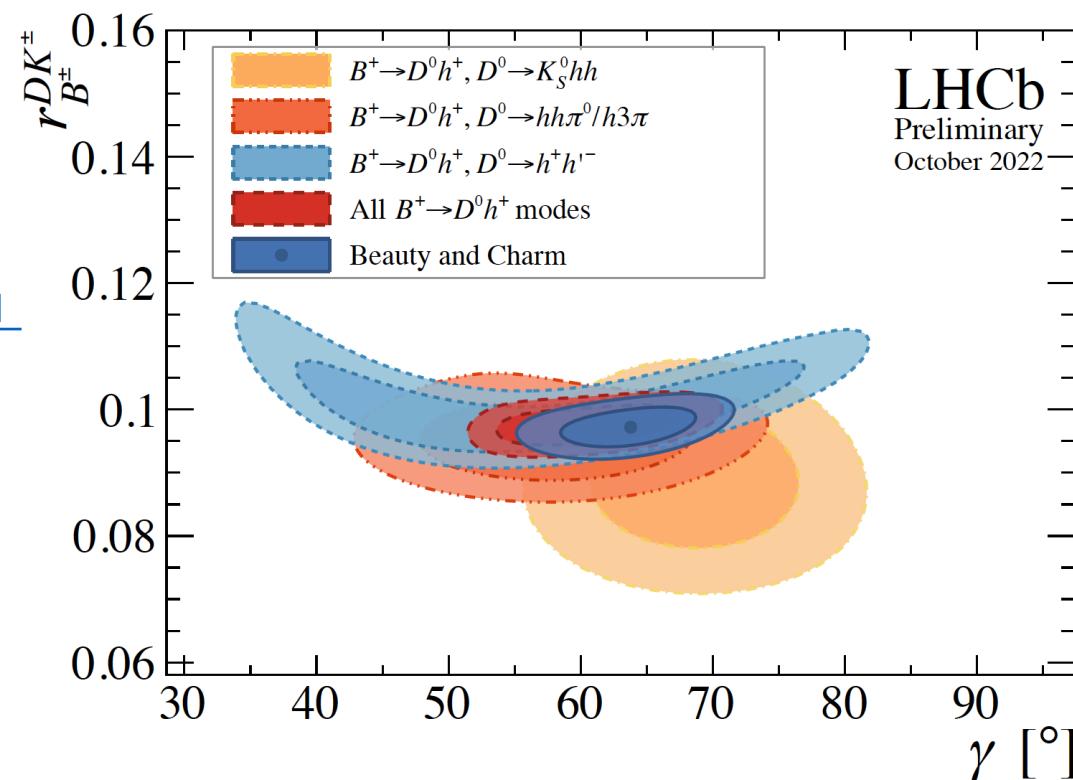
- $D^0 - \bar{D}^0$ mixing parameters [PRL 116 \(2016\) 241801](#) by LHCb

- 2nd most precise determination of γ from single D-decay mode

$$\gamma = \left(54.8^{+6.0+0.6+6.7}_{-5.8-0.6-4.3} \right)^\circ$$

- New LHCb combination:

$$\gamma = \left(63.8^{+3.5}_{-3.7} \right)^\circ$$



TD CPV in $B_s^0 \rightarrow J/\psi K^+ K^-$

- Mixing-induced CPV phase ϕ_s in B_s^0 decays through $b \rightarrow c\bar{c}s$ transitions.
- Golden mode : $B_s^0 \rightarrow J/\psi\phi(K^+K^-)$
- Flavor-tagged time-dependent angular analysis.

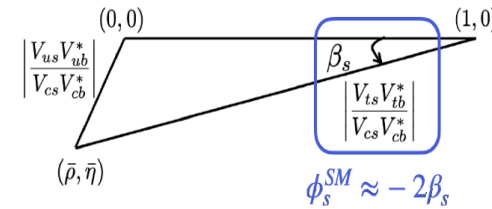
$$A_{CP}(t) = \frac{\Gamma(\bar{B}_s^0 \rightarrow J/\psi KK) - \Gamma(B_s^0 \rightarrow J/\psi KK)}{\Gamma(\bar{B}_s^0 \rightarrow J/\psi KK) + \Gamma(B_s^0 \rightarrow J/\psi KK)} = \eta_f \cdot \sin \phi_s^{\text{obs}} \cdot \sin(\Delta m_s t)$$

- CP eigenvalue of the final state $\eta_f = (-1)^L$
- A mixture of CP-even & CP-odd components \rightarrow angular analysis

$$\begin{array}{ll} \phi_s \text{ [rad]} & -0.039 \pm 0.022 \pm 0.006 \\ |\lambda| & 1.001 \pm 0.011 \pm 0.005 \end{array}$$

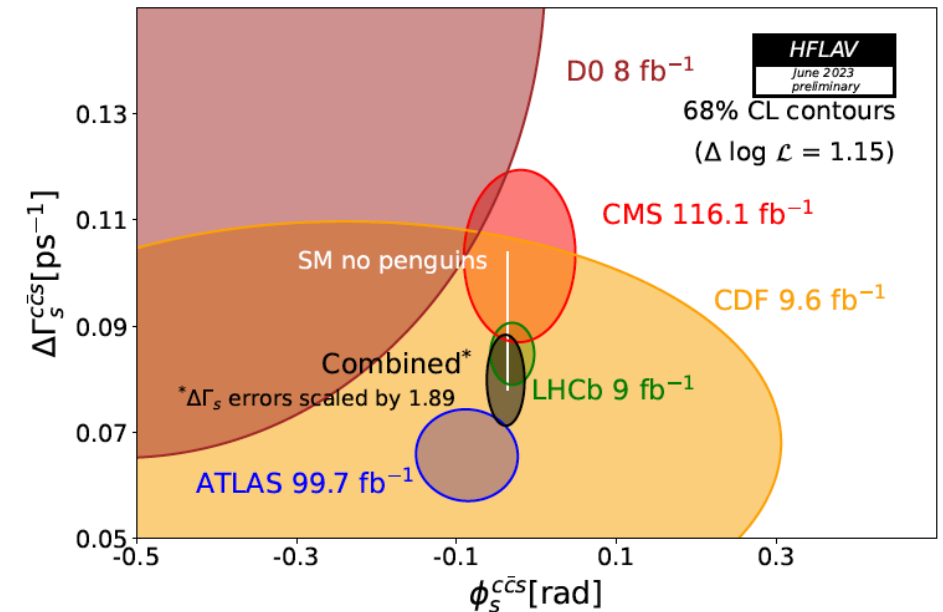
- The most precise measurement of ϕ_s to date.

$$V_{us}V_{ub}^* + V_{cs}V_{cb}^* + V_{ts}V_{tb}^* = 0$$



- Combination with all measurements:
 - $\phi_s^{J/\psi KK} = -0.050 \pm 0.017$ rad \rightarrow improved by 23%
 - $\phi_s^{c\bar{c}s} = -0.039 \pm 0.016$ rad \rightarrow improved by 15%
 - Consistent with the prediction of Global fits assuming SM:³

$$\phi_s^{\text{tree}} \approx (-0.0368_{-0.0009}^{+0.0006}) \text{ rad}$$



TD CPV in $B_S^0 \rightarrow \phi\phi$

- Penguin-dominated decay $b \rightarrow s\bar{s}s$.
 - NP could contribute to mixing and penguin processes.
 - Three linear polarization states for $\phi\phi$.
 - Flavor-tagged time-dependent angular analysis.
-
- The most precise measurement of $\phi_s^{s\bar{s}s}$ in penguin dominated decays.
 - Agreement with the SM.
 - Polarization-dependent CP-violation parameters are measured for the first time.

- polarization independent result:

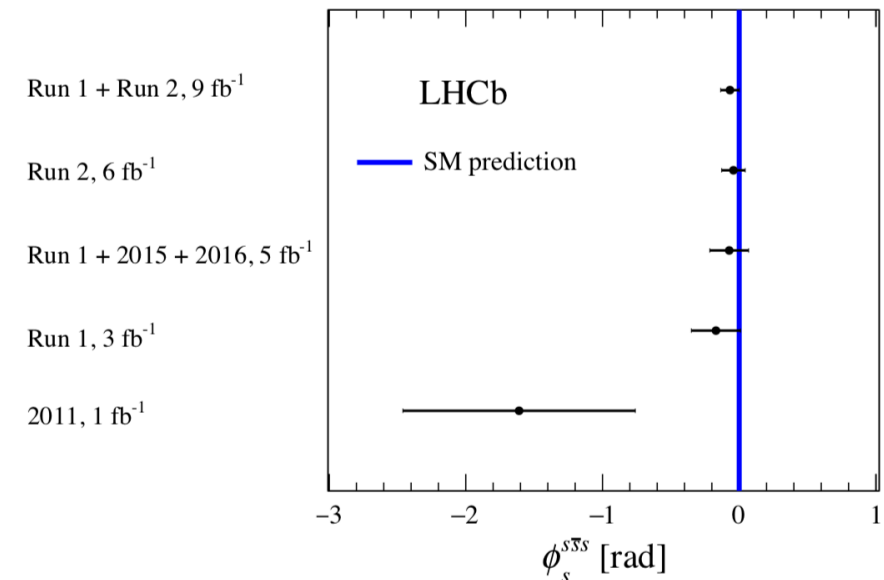
$$\phi_s^{s\bar{s}s} = -0.042 \pm 0.075 \pm 0.009 \text{ rad},$$

$$|\lambda| = 1.004 \pm 0.030 \pm 0.009,$$

Combination with Run 1 gives:

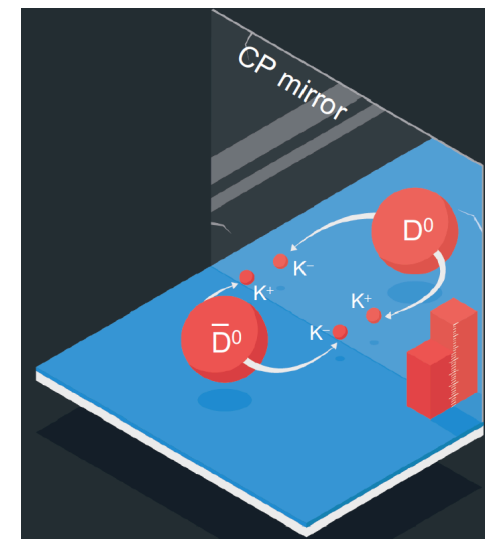
$$\phi_s^{s\bar{s}s} = (-0.074 \pm 0.069) \text{ rad}$$

$$|\lambda| = 1.009 \pm 0.030$$



CP violation in charm sector

- Charm provides unique gate for study of CP violation in up-type quark decays.
- Due to smallness of involved CKM elements and GIM mechanism, CP violation in charm is predicted to be small: 10^{-4} - 10^{-3} .
- SM calculation dominated by long distance contributions.
- LHCb reported direct CP violation in $D^0 \rightarrow h^+ h^-$ in March 2019, using the difference between two time-integrated CP-violating asymmetries of Cabibbo-suppressed D^0 decays [Phys. Rev. Lett. 122 \(2019\) 211803](#)
- The result challenges predictions based on first- principle QCD dynamics
- Further measurements are needed in charm sector.



$D^0 \rightarrow K^+ K^-$

- Time-integrated CPV measurement in the Cabibbo-suppressed decay $D^0 \rightarrow K^+ K^-$.
- D^0 mesons originate from promptly produced $D^{*+} \rightarrow D^0 \pi^+$.

$$A(K^- K^+) \approx \boxed{\mathcal{A}_{CP}(K^- K^+)} + \boxed{A_P(D^{*+})} + \boxed{A_D(\pi_{\text{tag}}^+)}$$

Physical CP asymmetry Production asy. Detection asy. of tagging pions

$$A_{CP}(D^0 \rightarrow K^- K^+) = +A(D^{*+} \rightarrow (D^0 \rightarrow K^- K^+) \pi_{\text{soft}}^+) - A(D^{*+} \rightarrow (D^0 \rightarrow K^- \pi^+) \pi_{\text{soft}}^+) + A(D^+ \rightarrow K^- \pi^+ \pi^+) - [A(D^+ \rightarrow \bar{K}^0 \pi^+) - A(\bar{K}^0)]$$

Original method (used also in Run 1)

$$A_{CP}(D^0 \rightarrow K^- K^+) = +A(D^{*+} \rightarrow (D^0 \rightarrow K^- K^+) \pi_{\text{soft}}^+) - A(D^{*+} \rightarrow (D^0 \rightarrow K^- \pi^+) \pi_{\text{soft}}^+) + A(D_s^+ \rightarrow \phi \pi^+) - [A(D_s^+ \rightarrow \bar{K}^0 K^+) - A(\bar{K}^0)]$$

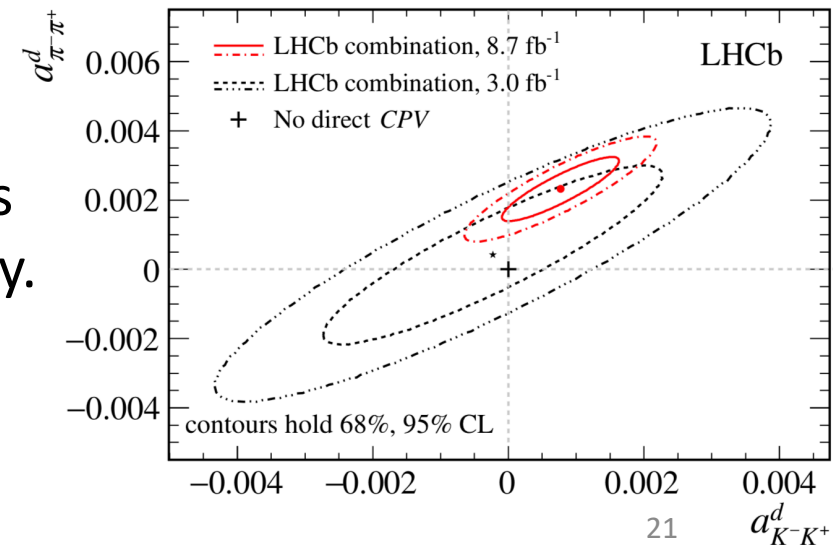
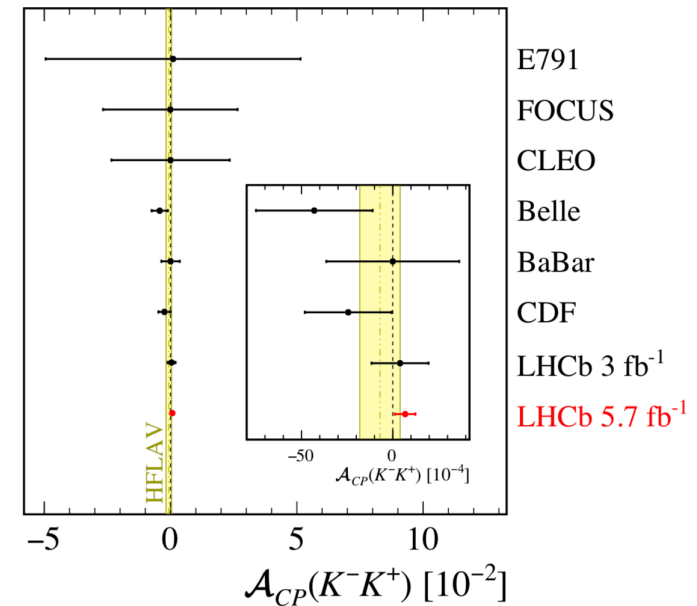
New method

- A combination with the previous LHCb measurements shows the first evidence for direct CP violation in $D^0 \rightarrow \pi^+ \pi^-$ decay.

$$\mathcal{A}_{CP}(K^- K^+) = [6.8 \pm 5.4 (\text{stat}) \pm 1.6 (\text{syst})] \times 10^{-4}$$

$$a_{K^- K^+}^d = (7.7 \pm 5.7) \times 10^{-4},$$

$$a_{\pi^- \pi^+}^d = (23.2 \pm 6.1) \times 10^{-4},$$



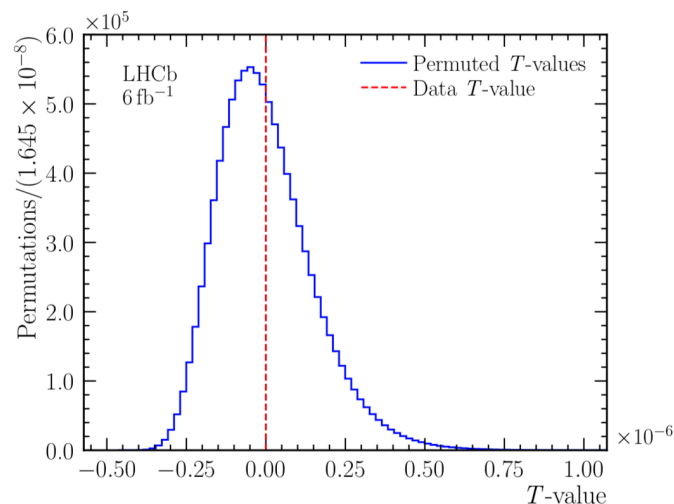
$$D^0 \rightarrow \pi^- \pi^+ \pi^0$$

- $D^0 \rightarrow \pi^- \pi^+ \pi^0$ is similar with $D^0 \rightarrow \pi^- \pi^+$. CS decay. Sensitive to interference between tree and penguin diagrams.
- Three-body phase space can enrich local CPV effects.
- The energy test computes the T-value over all D^0 - D^0 pairs (n) and $\overline{D^0}$ - $\overline{D^0}$ pairs (\bar{n}) and between D^0 - $\overline{D^0}$. [PhysRevD.84.054015](#)

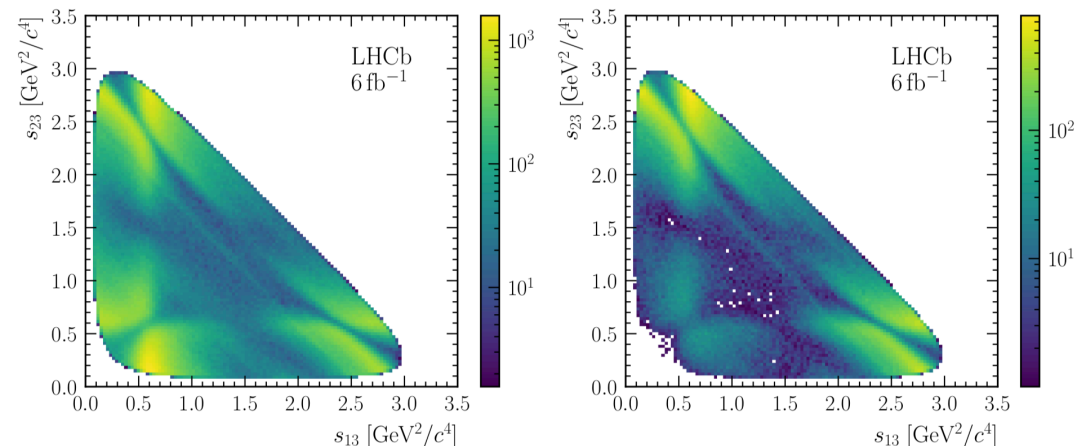
$$T \equiv \frac{1}{2n(n-1)} \sum_{i,j \neq i}^n \psi_{ij} + \frac{1}{2\bar{n}(\bar{n}-1)} \sum_{i,j \neq i}^{\bar{n}} \psi_{ij} - \frac{1}{n\bar{n}} \sum_{i,j}^{n,\bar{n}} \psi_{ij}. \quad \psi_{ij} = e^{-d_{ij}^2/2\delta^2}$$

$$d_{ij}^2 = [(\Delta s_{12})_{ij}^2 + (\Delta s_{13})_{ij}^2 + (\Delta s_{23})_{ij}^2]$$

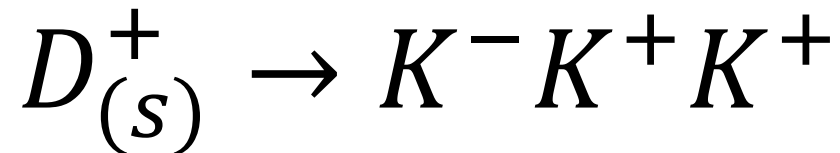
- The function ψ_{ij} gives a weighted distance between pairs in phase space.
- The tunable parameter δ is optimized to maximize sensitivity to CP violation.



The results are consistent with CP symmetry.



- Background-subtracted signal samples for the resolved (left) and merged (right) π^0 .

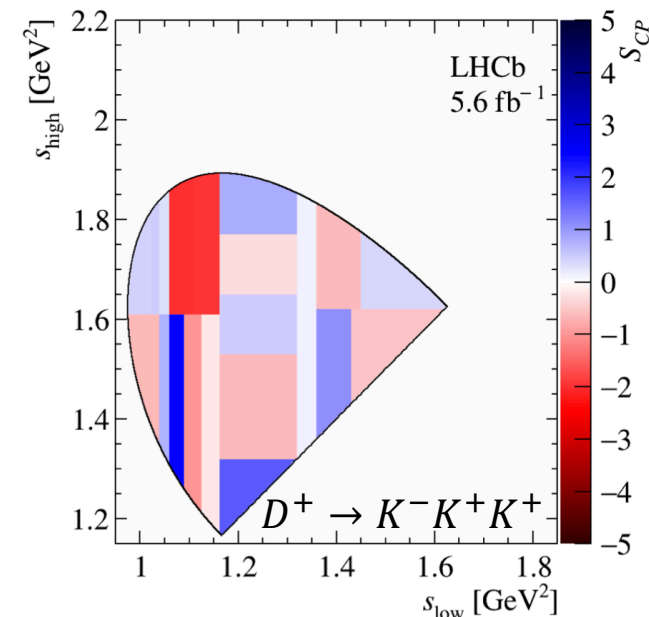
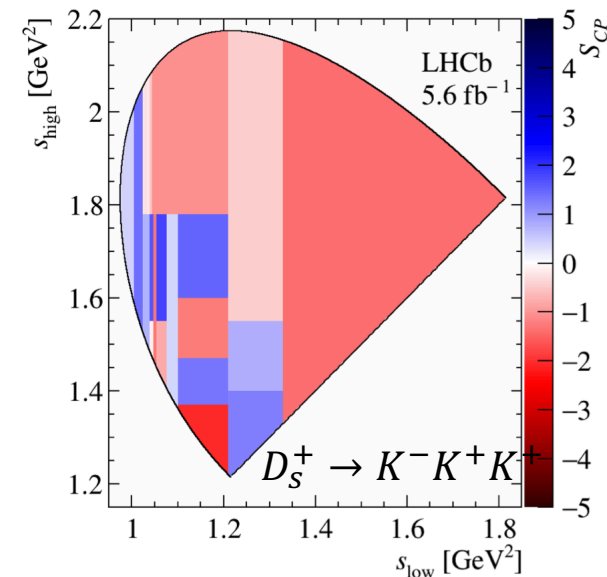


- SCS decay $D_s^+ \rightarrow K^- K^+ K^+$; DCS decay $D^+ \rightarrow K^- K^+ K^+$, no CPV in SM.
- Model-independent binned Dalitz-plot analysis.
- Dalitz plot divided in 21 bins that reproduce the pattern of the main resonances (\approx constant strong phase).
- Fit-per-bin method. Variation of the original Miranda technique: test χ^2 to compare Dalitz distributions of $N^i(D_s^+)$ and $N^i(D_s^-)$ (signal yields obtained by mass fit in each bin)

α takes into account global nuisance asymmetries

$$S_{CP}^i = \frac{N^i(D_{(s)}^+) - \alpha N^i(D_{(s)}^-)}{\sqrt{\alpha (\delta_{N^i(D_{(s)}^+)}^2 + \delta_{N^i(D_{(s)}^-)}^2)}}, \quad \alpha = \frac{\sum_i N^i(D_{(s)}^+)}{\sum_i N^i(D_{(s)}^-)}, \quad \chi^2 = \sum_i (S_{CP}^i)^2$$

- No evidence for CP violation is observed in these decays.



T-odd CPV in $D_{(s)}^+$ four-body decays

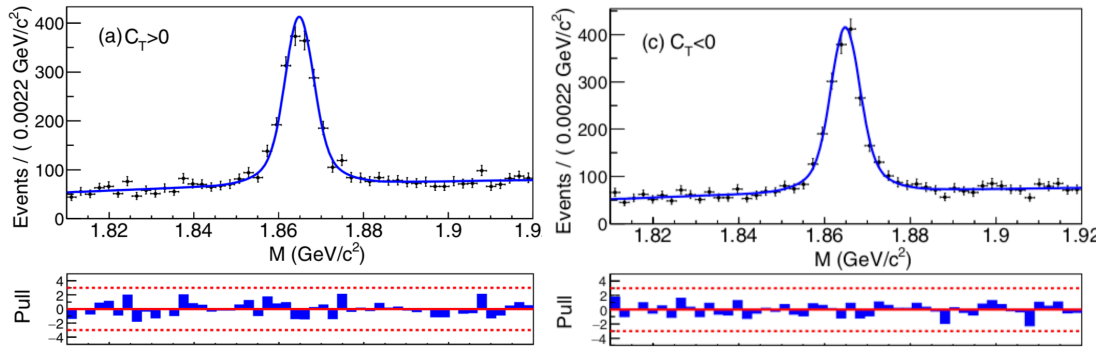
- T -odd correlations provides a powerful tool to indirectly search for CP violation under CPT symmetry conservation.
- C_T observable defined by a triple mixed product $C_T = (\vec{p}_1 \times \vec{p}_2) \cdot \vec{p}_3$, satisfying $CP(C_T) = -C(C_T) = -\bar{C}_T$

• All results consistent with zero CP Violation.

$$A_T = \frac{\Gamma_+(C_T > 0) - \Gamma_+(C_T < 0)}{\Gamma_+(C_T > 0) + \Gamma_+(C_T < 0)} \quad \bar{A}_T = \frac{\Gamma_-(-\bar{C}_T > 0) - \Gamma_-(-\bar{C}_T < 0)}{\Gamma_-(-\bar{C}_T > 0) + \Gamma_-(-\bar{C}_T < 0)}$$

- T-odd CP-violating asymmetry is defined as $a_{CP}^{T\text{-odd}} = \frac{1}{2}(A_T - \bar{A}_T)$

$D^0 \rightarrow K_S^0 K_S^0 \pi^+ \pi^-$ [PRD 107, 052001 \(2023\)](#)

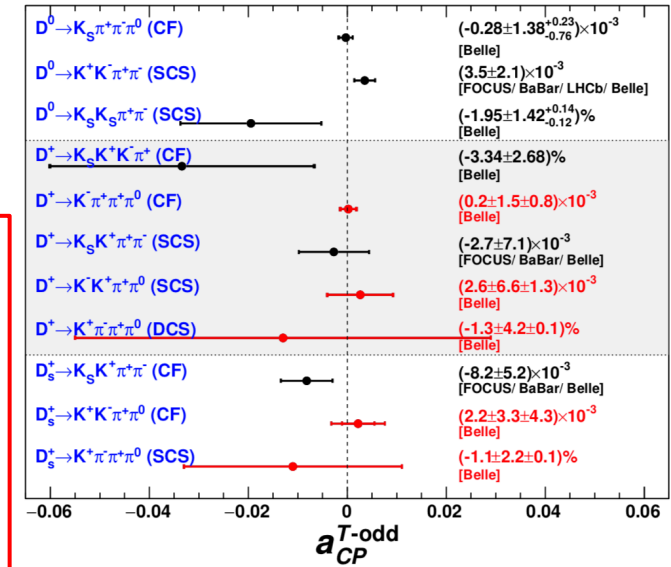


$$A_{CP}(D^0 \rightarrow K_S^0 K_S^0 \pi^+ \pi^-) = [-2.51 \pm 1.44(\text{stat})_{-0.10}^{+0.11}(\text{syst})]\%$$

$$a_{CP}^T(D^0 \rightarrow K_S^0 K_S^0 \pi^+ \pi^-) = [-1.95 \pm 1.42(\text{stat})_{-0.12}^{+0.14}(\text{syst})]\%$$

- α_{CP}^T is proportional to $\sin\phi\cos\delta$, ϕ and δ are the weak and strong phase differences.
- A_{CP} (α_{CP}^T) is sensitive to $\delta = \pi(0)$.

$D_{(s)}^+ \rightarrow K^\pm h^\mp \pi^+ \pi^0$ [arXiv:2305.12806](#)



$D_{(s)}^+ \rightarrow K^+ K_S^0 h^- h^-$ [arXiv:2305.11405](#)

$$a_{CP}^{T\text{-odd}}(D^+ \rightarrow K^+ K_S^0 \pi^+ \pi^-) = (0.34 \pm 0.87 \pm 0.32)\%$$

$$a_{CP}^{T\text{-odd}}(D_s^+ \rightarrow K^+ K_S^0 \pi^+ \pi^-) = (-0.46 \pm 0.63 \pm 0.38)\%$$

$$a_{CP}^{T\text{-odd}}(D^+ \rightarrow K^+ K^- K_S^0 \pi^+) = (-3.34 \pm 2.66 \pm 0.35)\%$$



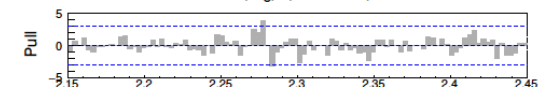
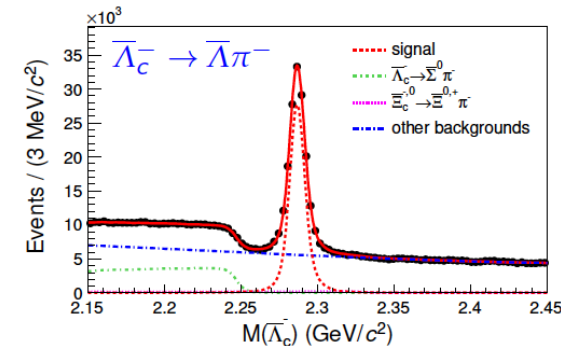
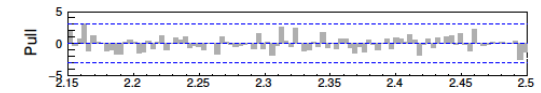
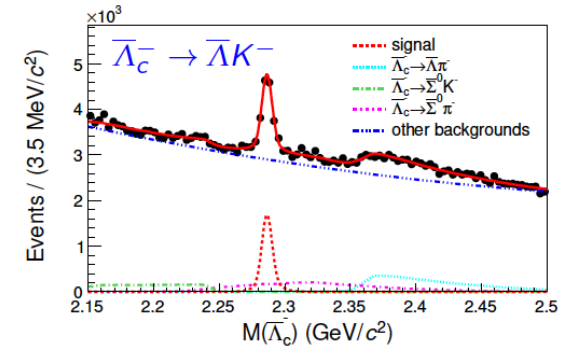
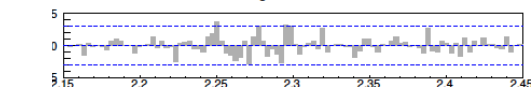
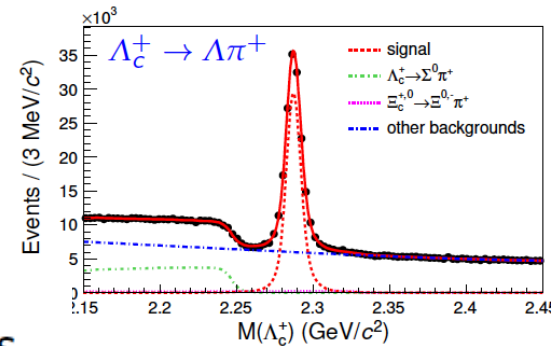
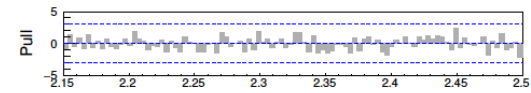
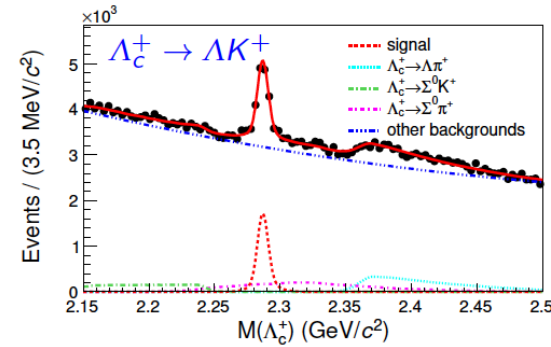
CPV in $\Lambda_c^+ \rightarrow \Lambda h^+$ and $\Lambda_c^+ \rightarrow \Sigma^0 h^+$ ($h = K, \pi$)

- Baryogenesis, the process by which the baryon-antibaryon asymmetry of the universe developed, is directly related to baryon CPV.
- To date, CPV has been confirmed in the open-flavored meson sector (K, D and B mesons), but not in the baryon sector.

$$A_{\text{raw}}(\Lambda_c^+ \rightarrow \Lambda K^+) \approx A_{CP}^{\Lambda_c^+ \rightarrow \Lambda K^+} + A_{CP}^{\Lambda \rightarrow p\pi^-} + A_\epsilon^\Lambda + A_\epsilon^{K^+} + A_{\text{FB}}^{\Lambda_c^+}$$

$$\begin{aligned} & A_{\text{raw}}^{\text{corr}}(\Lambda_c^+ \rightarrow \Lambda K^+) - A_{\text{raw}}^{\text{corr}}(\Lambda_c^+ \rightarrow \Lambda \pi^+) \\ &= A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Lambda K^+) - A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Lambda \pi^+). \end{aligned}$$

- Simultaneous fit on the $A_\epsilon^{h^+}$ -weighted Λ_c^\pm samples gives
 $A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Lambda K^+) = (+2.1 \pm 2.6 \pm 0.1)\%$
 $A_{CP}^{\text{dir}}(\Lambda_c^+ \rightarrow \Sigma^0 K^+) = (+2.5 \pm 5.4 \pm 0.4)\%$,
first CPV result of charmed baryon SCS two-body decays.



A_{CP}^α in $\Lambda_c^+ \rightarrow \Lambda h^+, \Sigma^0 h^+$

- The decay asymmetry parameter α was introduced by Lee and Yang to study the parity-violating and parity-conserving amplitudes in weak hyperon decays.

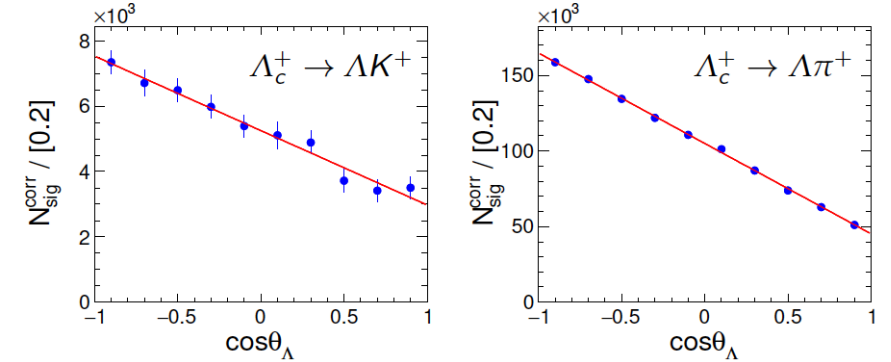
$$\frac{dN(\Lambda_c^+ \rightarrow \Lambda h^+)}{d \cos \theta_\Lambda} \propto 1 + \alpha_{\Lambda_c^+} \alpha_- \cos \theta_\Lambda$$

$$\frac{dN(\Lambda_c^+ \rightarrow \Sigma^0 h^+)}{d \cos \theta_{\Sigma^0} d \cos \theta_\Lambda} \propto 1 - \alpha_{\Lambda_c^+} \alpha_- \cos \theta_{\Sigma^0} \cos \theta_\Lambda$$

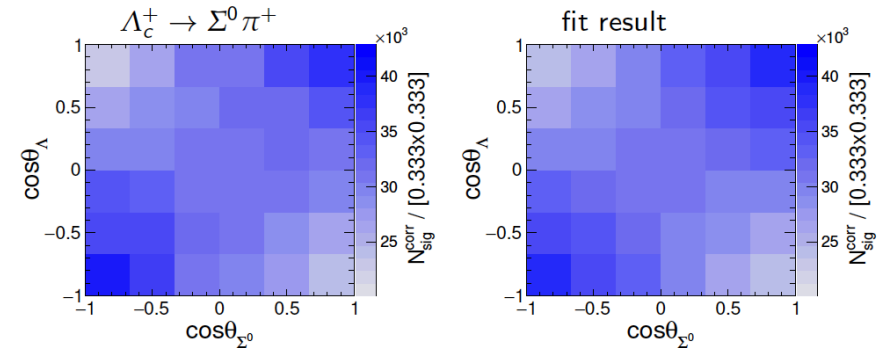
- The α -induced CP asymmetry:

$$A_{CP}^\alpha \equiv \frac{\alpha_{\Lambda_c^+} - \widehat{CP} \alpha_{\Lambda_c^+} \widehat{CP}^\dagger}{\alpha_{\Lambda_c^+} + \widehat{CP} \alpha_{\Lambda_c^+} \widehat{CP}^\dagger} = \frac{\alpha_{\Lambda_c^+} + \alpha_{\Lambda_c^-}}{\alpha_{\Lambda_c^+} - \alpha_{\Lambda_c^-}}$$

- $\cos \theta_\Lambda$ distributions of $\Lambda_c^+ \rightarrow \Lambda h^+$ after efficiency correction, fitted with $1 + \alpha_{\Lambda_c^+} \alpha_- \cos \theta_\Lambda$



- $(\cos \theta_\Lambda, \cos \theta_{\Sigma^0})$ 2D distributions of $\Lambda_c^+ \rightarrow \Sigma^0 h^+$ after efficiency correction, fitted with $1 - \alpha_{\Lambda_c^+} \alpha_- \cos \theta_{\Sigma^0} \cos \theta_\Lambda$



Channel	$k = \alpha_{\Lambda_c^+} \alpha_-$	$\bar{k} = \alpha_{\Lambda_c^-} \alpha_+$	$\alpha_{\Lambda_c^+}$	$\alpha_{\Lambda_c^-}$	A_{CP}^α	W.A. A_{CP}^α	our $A_{CP}^\alpha (\Lambda \rightarrow p \pi^-)$
$\Lambda_c^+ \rightarrow \Lambda K^+$	-0.418 ± 0.053	-0.442 ± 0.053	$-0.566 \pm 0.071 \pm 0.028$	$0.592 \pm 0.070 \pm 0.079$	$-0.023 \pm 0.086 \pm 0.071$	—	—
$\Lambda_c^+ \rightarrow \Lambda \pi^+$	-0.582 ± 0.006	-0.565 ± 0.006	$-0.784 \pm 0.008 \pm 0.006$	$0.754 \pm 0.008 \pm 0.018$	$+0.020 \pm 0.007 \pm 0.013$	-0.07 ± 0.22	$+0.017 \pm 0.007 \pm 0.012$
$\Lambda_c^+ \rightarrow \Sigma^0 K^+$	-0.43 ± 0.18	-0.37 ± 0.21	$-0.58 \pm 0.24 \pm 0.09$	$0.49 \pm 0.28 \pm 0.14$	$+0.08 \pm 0.35 \pm 0.14$	—	—
$\Lambda_c^+ \rightarrow \Sigma^0 \pi^+$	-0.340 ± 0.016	-0.358 ± 0.017	$-0.452 \pm 0.022 \pm 0.023$	$0.473 \pm 0.023 \pm 0.035$	$-0.023 \pm 0.034 \pm 0.030$	—	$-0.026 \pm 0.034 \pm 0.030$

combined: $+0.013 \pm 0.007 \pm 0.011$

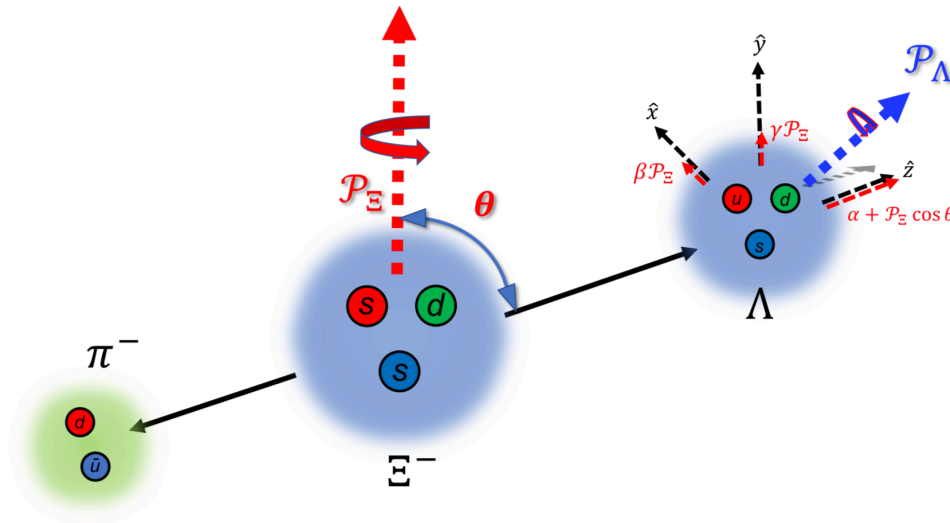
- No evidence of CPV in baryon (Λ_c^+ and Λ) decays is found.

CP symmetry test with entangled baryons

- $J/\psi \rightarrow \Xi^- \Xi^+ \rightarrow \Lambda \pi^- \Lambda \pi^+$

- Self analyzing hyperons: daughter particles are emitted according to polarization of mother hyperon.
- From decay amplitudes one can construct CP-odd decay parameters $\alpha_{\Xi}, \beta_{\Xi}, \gamma_{\Xi}$

$$\beta_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2 \sin \phi_{\Xi}}; \gamma_{\Xi} = \sqrt{1 - \alpha_{\Xi}^2 \cos \phi_{\Xi}}$$



$$A_{\text{CP}}^Y = \frac{\alpha_Y + \bar{\alpha}_Y}{\alpha_Y - \bar{\alpha}_Y}, \quad \Delta\phi_{\text{CP}} = \frac{\phi_Y + \bar{\phi}_Y}{2}$$

$$A_{\text{CP}}^{\Xi} \approx -\tan(\delta_P - \delta_S) \tan(\xi_P - \xi_S)$$

strong phase diff. weak phase diff.

Longitudinal polarization of Λ governed by α_{Ξ} ; ϕ_{Ξ} gives rotation of P_{Λ} with respect to P_{Ξ}

ϕ_{Ξ} allows for new CP test: $\Delta\phi_{\Xi} = (\phi_{\Xi} + \bar{\phi}_{\Xi})/2$

$$J/\psi \rightarrow \Xi^- \Xi^+ \rightarrow \Lambda \pi^- \Lambda \pi^+$$

- Nine-dimensional phase space given by nine helicity angles

$$\xi = (\theta, \theta_\Lambda, \varphi_\Lambda, \theta_{\bar{\Lambda}}, \varphi_{\bar{\Lambda}}, \theta_p, \varphi_p, \theta_{\bar{p}}, \varphi_{\bar{p}}).$$

- Eight free global parameters determined by maximum log likelihood method: $\omega = (\alpha_\psi, \Delta\Phi, \alpha_\Xi, \phi_\Xi, \bar{\alpha}_\Xi, \bar{\phi}_\Xi, \alpha_\Lambda, \bar{\alpha}_\Lambda)$

- The formalism exploits **polarization**, **entanglement** and **sequential decays**

$$\mathcal{W}(\xi; \omega) = \sum_{\mu, \nu=0}^3 \underbrace{C_{\mu\nu}}_{\text{blue circle}} \sum_{\mu', \nu'=0}^3 \underbrace{a_{\mu\mu'}^\Xi a_{\nu\nu'}^{\bar{\Xi}} a_{\mu'0}^\Lambda a_{\nu'0}^{\bar{\Lambda}}}_{\text{green oval}} \quad C_{\mu\nu}(\theta; \alpha_\psi, \Delta\Phi) :$$

Parameter	This work	Previous result
A_{CP}^Ξ	$(6.0 \pm 13.4 \pm 5.6) \times 10^{-3}$	–
$\Delta\phi_{\text{CP}}^\Xi$	$(-4.8 \pm 13.7 \pm 2.9) \times 10^{-3}$ rad	–
A_{CP}^Λ	$(-3.7 \pm 11.7 \pm 9.0) \times 10^{-3}$	$(-6 \pm 12 \pm 7) \times 10^{-3}$ [14]

Other CP symmetry test in hyperon decays

- $J/\psi \rightarrow \Lambda \bar{\Lambda}$ [Phys.Rev.Lett. 129 \(2022\) 131801](#)

- 1×10^{10} J/ψ events

$$A_{CP} = (\alpha_- + \alpha_+) / (\alpha_- - \alpha_+) = -0.0025 \pm 0.0046 \pm 0.0012,$$

- $J/\psi \rightarrow \Sigma^+ \bar{\Sigma}^-$ [arXiv:2304.14655](#)

- 1×10^{10} J/ψ events

- First study to test CP symmetry in the hyperon to neutron decay, and result is consistent with CP-conservation.

- $J/\psi (\psi(3686)) \rightarrow \Xi^0 \bar{\Xi}^0$

- 1×10^{10} J/ψ events; 4×10^8 $\psi(3686)$ events

- Results are consistent with CP-conservation.

[arXiv:2305.09218](#)

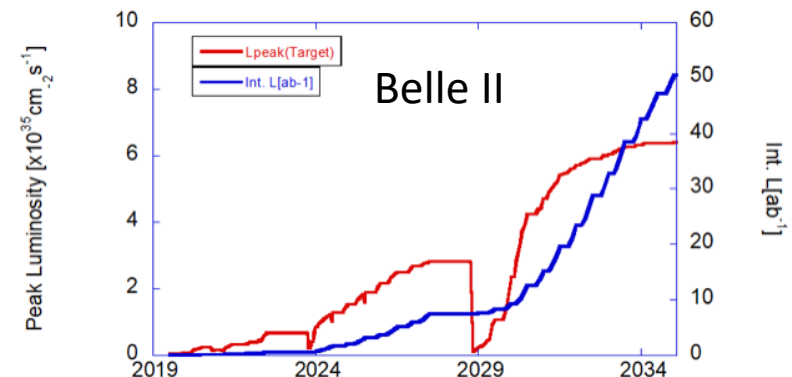
A_{CP}^{Ξ}	$(-5.4 \pm 6.5 \pm 3.1) \times 10^{-3}$
$\Delta\phi_{CP}^{\Xi}(\text{rad})$	$(-0.1 \pm 6.9 \pm 0.9) \times 10^{-3}$
A_{CP}^{Λ}	$(6.9 \pm 5.8 \pm 1.8) \times 10^{-3}$

[arXiv:2302.09767](#)

A_{CP}^{Ξ}	$-0.007 \pm 0.082 \pm 0.025$
$\Delta\phi_{CP}^{\Xi}$	$-0.079 \pm 0.082 \pm 0.010$

Summary

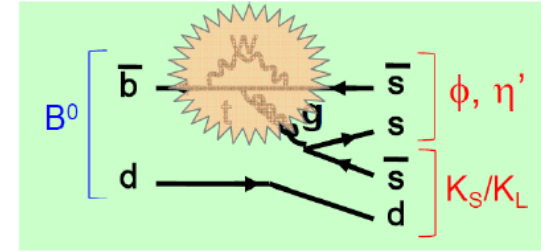
- CP violation is being tested in broad physics programs at different experiments.
- We are in the era of seeking for new physics corrections to SM CPV.
- No evidence for new CPV physics so far.
- More results from Belle II/LHCb/BESIII (etc.) are anticipated.
 - Belle II: 428 fb⁻¹ recorded. Goal: 50 ab⁻¹.
 - LHCb: 9 fb⁻¹ in Run1-2. More than 5 times luminosity during Runs 3–4.
 - BESIII: collected $1.0 \times 10^{10} J/\psi$ events, rich program using entangled hyperon pairs.
- Large datasets will allow for unprecedented reach on precise CPV measurements.



Extra

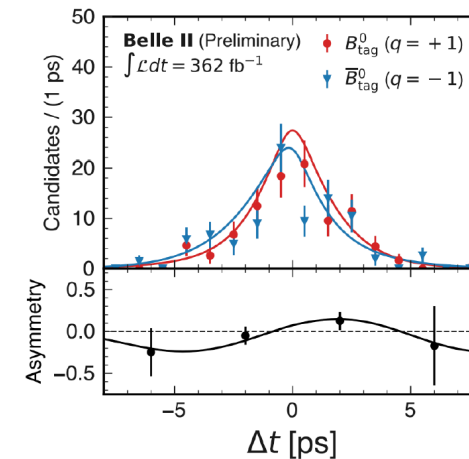
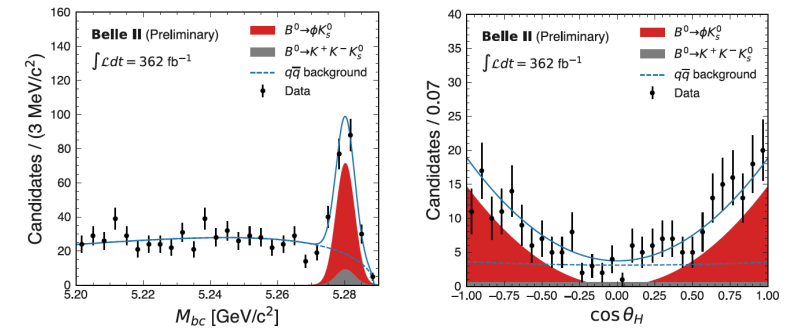


$b \rightarrow s$ ($B \rightarrow \phi K^0, \eta' K^0$)



TD CPV in $B^0 \rightarrow \phi K_S^0$

- Penguin-dominated modes $b \rightarrow q\bar{q}s$.
- Clean experimental signature with similar Δt resolution as $B^0 \rightarrow J/\psi K_S$.
- Main challenge: dilution from non-resonant decays with opposite CP
- Non-resonant $B^0 \rightarrow K^+K^-K_S^0$ component disentangled in $\cos\theta$.
- Simultaneous Δt fit to extract the CP asymmetries.
 - $B^- \rightarrow K^+K^-K_S$ fixed from HFLAV.
 - Validated on the B^- control sample (null asymmetry).
- 162 ± 17 signal yields.

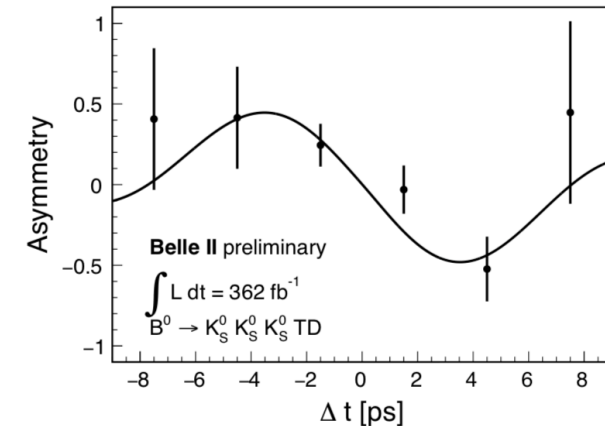
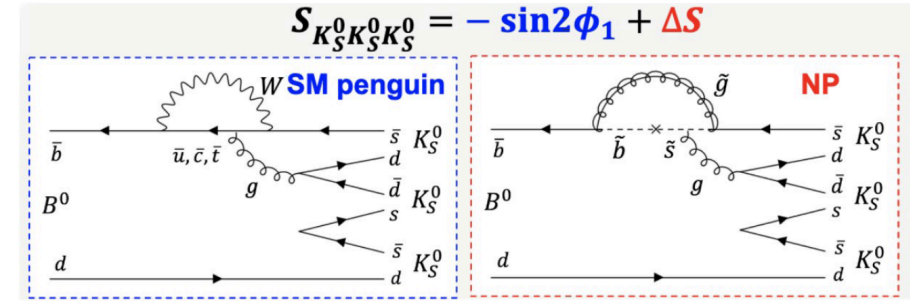


$$A_{CP} = 0.31 \pm 0.20^{+0.05}_{-0.06} \quad A_{CP}^{w.a.} = -0.01 \pm 0.14$$

$$S_{CP} = 0.54 \pm 0.26^{+0.06}_{-0.08} \quad S_{CP}^{w.a.} = 0.59 \pm 0.14$$

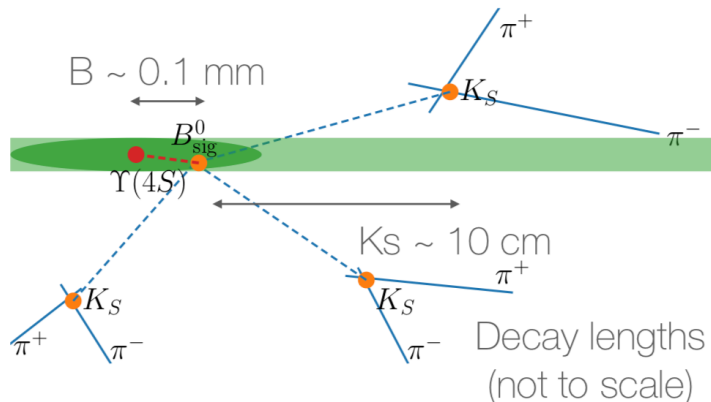
$B^0 \rightarrow K_S^0 K_S^0 K_S^0$

- Penguin process suppressed in the SM, sensitive to new physics.
- Same underlying quark transition as $B^0 \rightarrow \phi K_S^0$, no contribution from opposite-CP backgrounds.
- Challenge: no prompt tracks to from a vertex. K_S^0 flies 10 cm on average.
- Reconstruct $B^+ \rightarrow K_S^0 K_S^0 K^+$ as control mode, remove vertex information from K^+ .
- Simultaneous fit on control sample to constrain background shapes and Δt resolution function.
- 3D fit (invariant mass, CS output and Mbc) to extract signal yields.
- Dataset divided into events with (TD) and without (TI) information from the vertex detector.
- 158_{-13}^{+14} (TD) and 62 ± 9 (TI) $B^0 \rightarrow K_S^0 K_S^0 K_S^0$ signal evens with 387M $B\bar{B}$ pairs.
- TI events used only to constrain the time-integrated asymmetry A_{CP} .



$$\sin(2\beta)^{\text{eff}} = 1.36_{-0.45}^{+0.35} \text{ (stat.)} \pm 0.03 \text{ (syst.)}$$

$$A_{CP} = 0.07_{-0.20}^{+0.15} \text{ (stat.)} \pm 0.02 \text{ (syst.)}$$



Test iso-spin symmetry in $B \rightarrow K\pi$

- From the iso-spin symmetry in the SM:

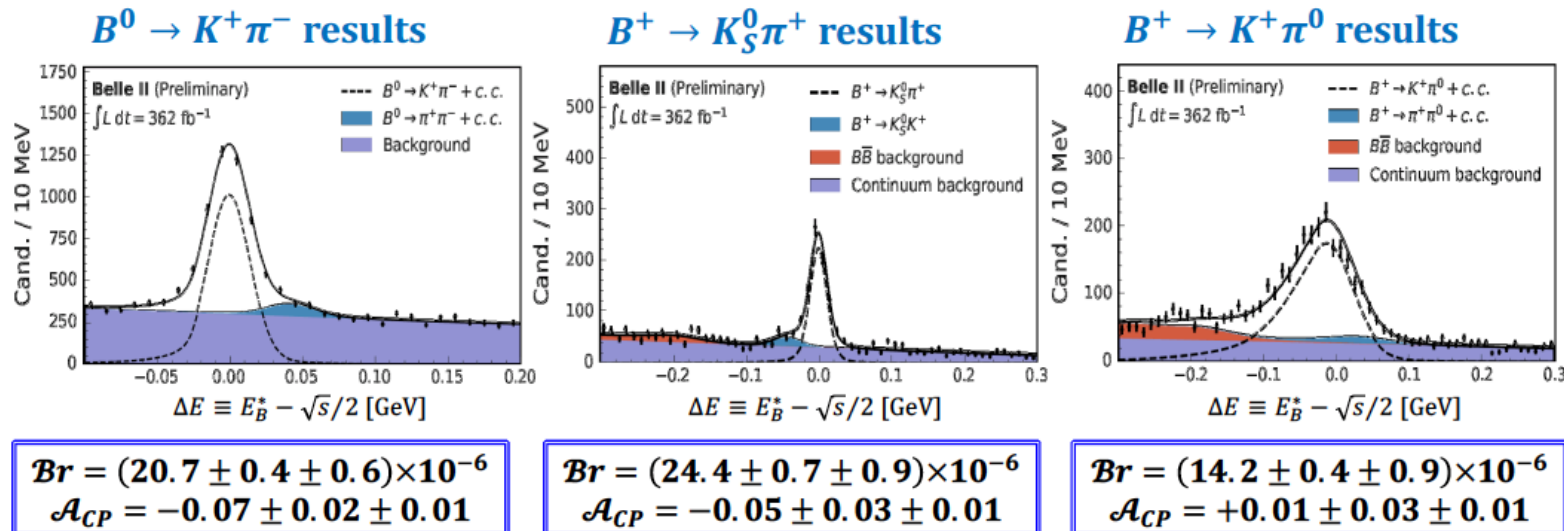
Phys.Lett.B 627 (2005) 82

Isospin sum-rule: relation among the products of $\mathcal{B}r$ and \mathcal{A}_{CP} for $B \rightarrow K\pi$

$$I_{K\pi} = \mathcal{B}r(K^+\pi^-)\mathcal{A}_{K^+\pi^-} + \mathcal{A}_{K^0\pi^+}\mathcal{B}r(K^0\pi^+) \frac{\tau_{B^0}}{\tau_{B^+}} - 2\mathcal{A}_{K^+\pi^0}\mathcal{B}r(K^+\pi^0) \frac{\tau_{B^0}}{\tau_{B^+}} - 2\mathcal{A}_{K^0\pi^0}\mathcal{B}r(K^0\pi^0); \quad I_{K\pi}^{SM} = 0$$

M. Gronau, Phys. Lett. B 627, 82(2005).

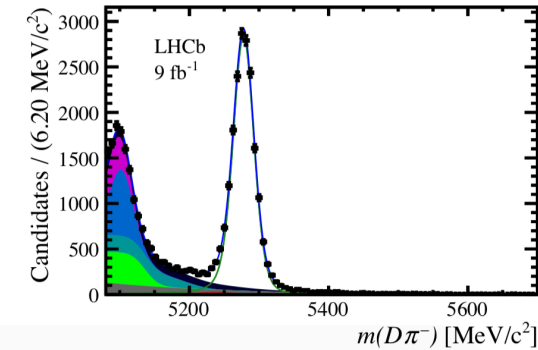
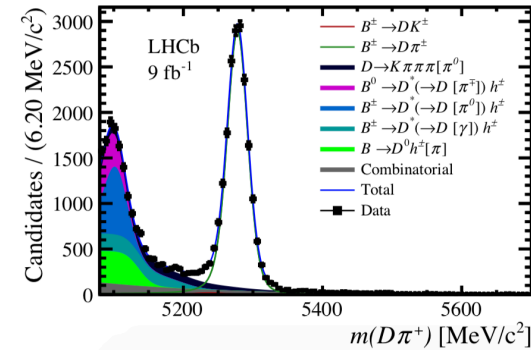
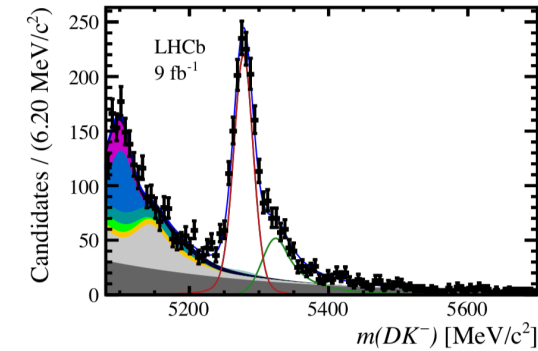
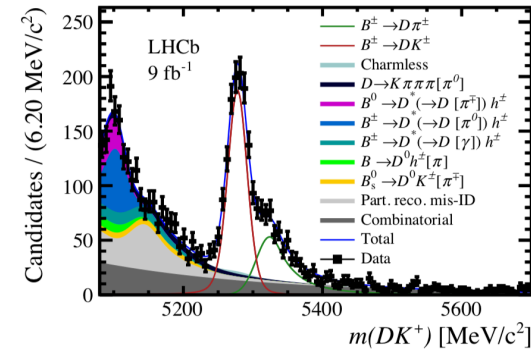
- Belle II can measure all the observables.



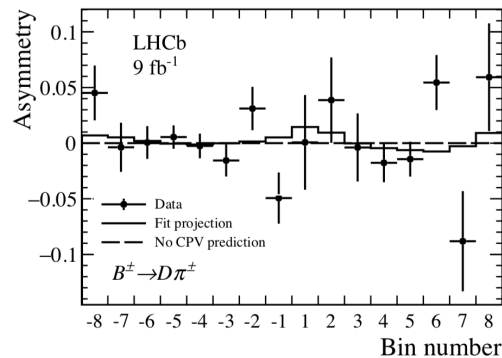
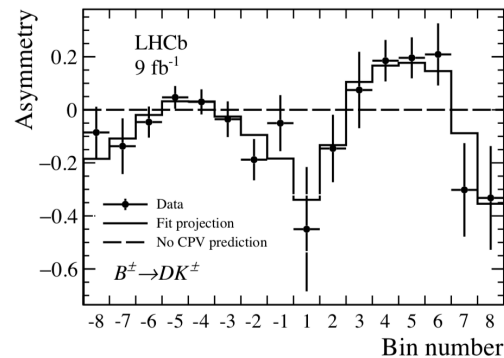
CPV in $B^+ \rightarrow D(h^+ h^- \pi^+ \pi^-) h^+$

- First study of CP violation in $B^+ \rightarrow D(K^+ K^- \pi^+ \pi^-) h^+$ and updated for $B^+ \rightarrow D(\pi^+ \pi^- \pi^+ \pi^-) h^+$.
- Global phase-space integrated measurement for $B^+ \rightarrow [K^+ K^- \pi^+ \pi^-]_D h^+$ and $B^+ \rightarrow [\pi^+ \pi^- \pi^+ \pi^-]_D h^+$.
- Binned measurement for $[K^+ K^- \pi^+ \pi^-]_D$. Sensitive to the local CP asymmetries.
- External information on charm-decay parameters, taken from LHCb amplitude analysis. [JHEP 02 \(2019\) 126](#)
- Precision will improve after charm model-independent measurements.

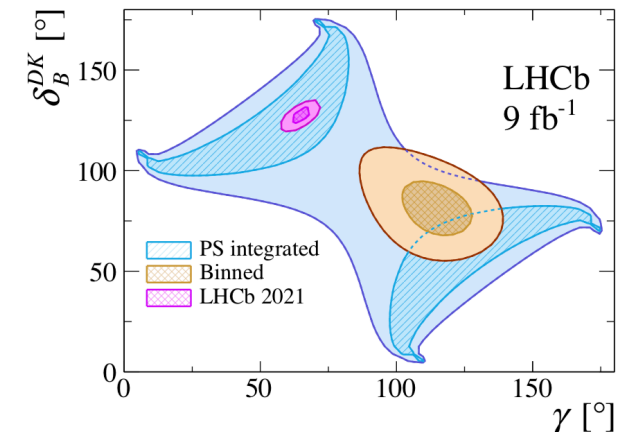
LHCb Run 1 + Run 2 data sample



Fractional bin asymmetries:

model-dependent γ :

$$\gamma = (116^{+12}_{-14})^{\circ}$$



Other topics

- CP violation in the lepton sector
 - neutrino mass eigenstates mismatch weak eigenstates
 - analogue to CKM matrix: PMNS matrix

In the 3-Flavor PMNS model:

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta_{CP}} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

- CP violation in QCD
 - QCD Lagrangian can contain CP violating relative phase $\bar{\theta}$
 - measurement of neutron electric dipole moment
 - Axion ALPs search