

# CP Violation, B Physics and Future Facilities

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## Introduction/Review/Motivation

Complex phases in the weak interaction:  $V_{td}$  and  $V_{ts}$  and associated CPV asymmetries

Connections to the charged Higgs

## Flavor Physics, The Next Generation: Belle II and the LHCb upgrade

Apologies: In the limited time, I cannot cover all the new results from BaBar, Belle, LHCb, CMS, ATLAS, Tevatron ... I have borrowed slides from many excellent physicists

# Amplitudes and Phases in the Weak Interaction

N. Cabibbo



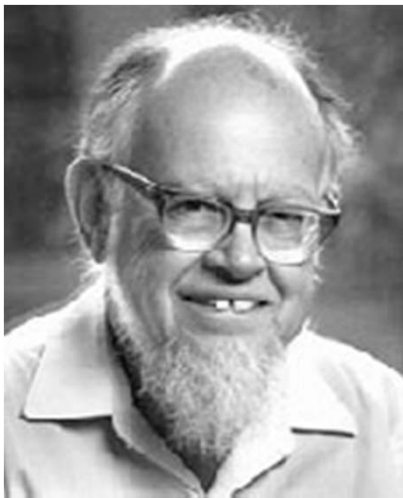
M. Kobayashi



T. Maskawa



$$V_{\text{CKM}} \equiv V_L^u V_L^{d\dagger} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ \underline{V_{td}} & V_{ts} & V_{tb} \end{pmatrix}$$

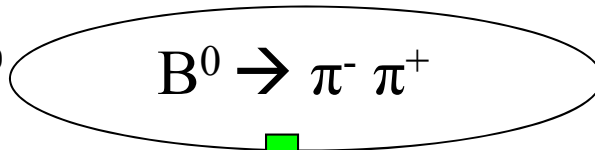


L. Wolfenstein

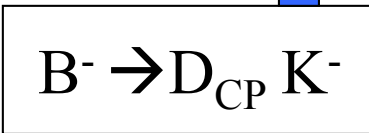
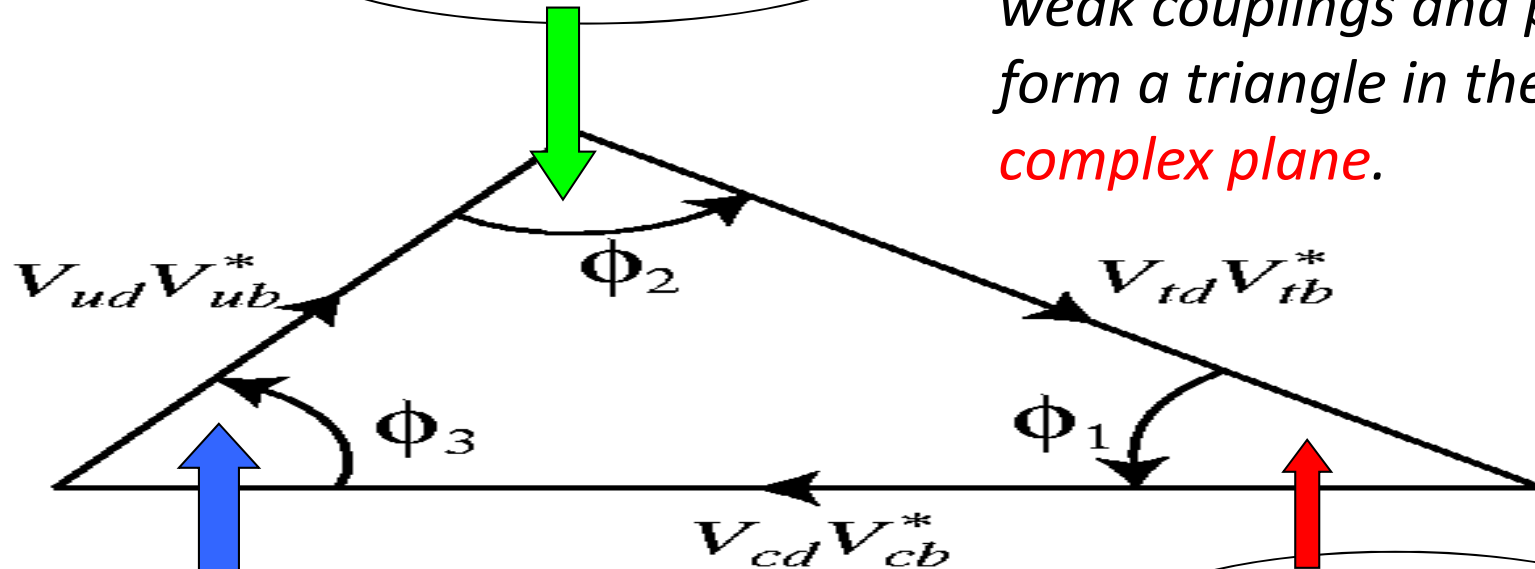
$$V = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & \underline{A\lambda^3(\rho - i\eta)} \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ \underline{A\lambda^3(1 - \rho - i\eta)} & -A\lambda^2 & 1 \end{pmatrix}$$

# Three Angles: $(\varphi_1, \varphi_2, \varphi_3)$ or $(\beta, \alpha, \gamma)$

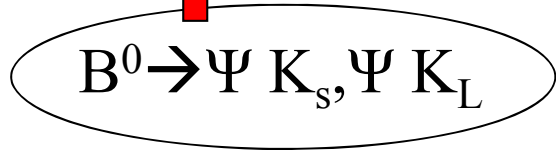
ICHEP 2014: New  
Belle result on  $\pi^0 \pi^0$



Unitarity implies that the weak couplings and phases form a triangle in the complex plane.



ICHEP2014: LHCb results on CPV in  $B_s \rightarrow D_s^- K^+$   
(A. Dziurda)



Big Questions: *Are determinations of angles consistent with determinations of the sides of the triangle? Are angle determinations from loop and tree decays consistent?*

Time-dependent  $CP$  violation is  
“A Double-Slit experiment” with particles and antiparticles

QM interference between two diagrams

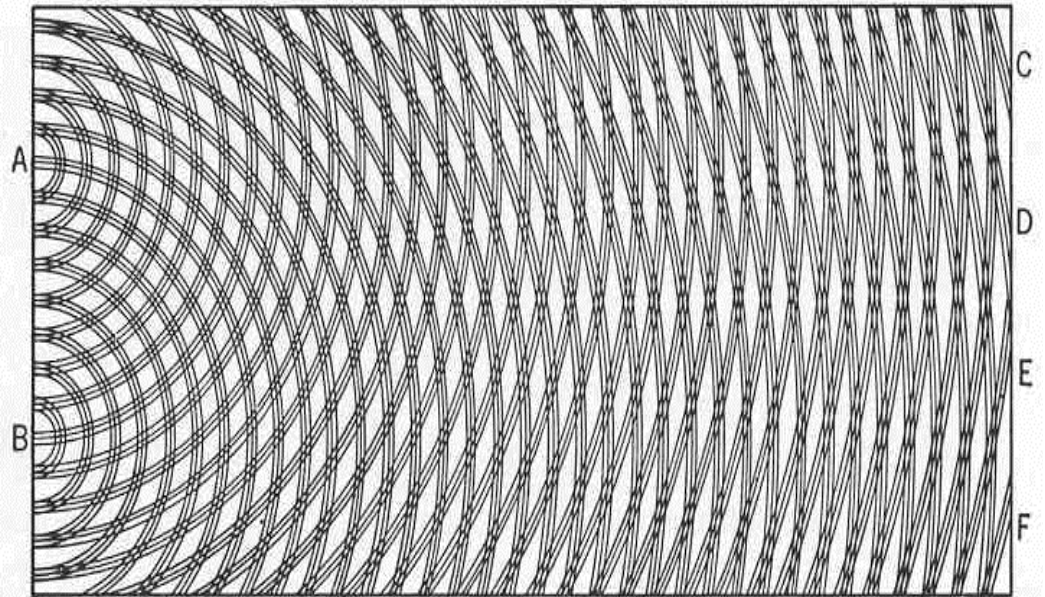
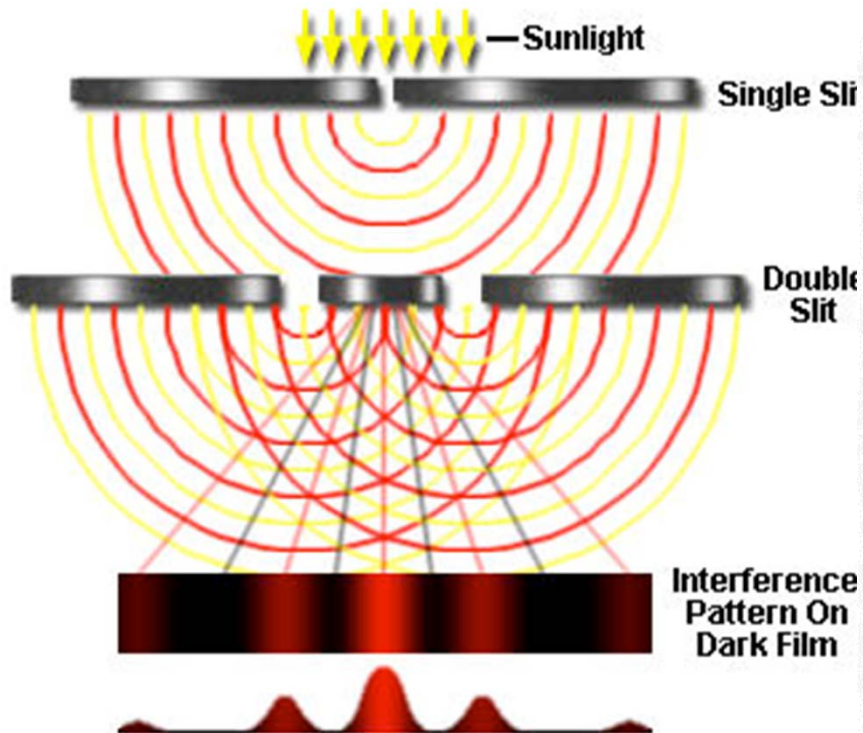
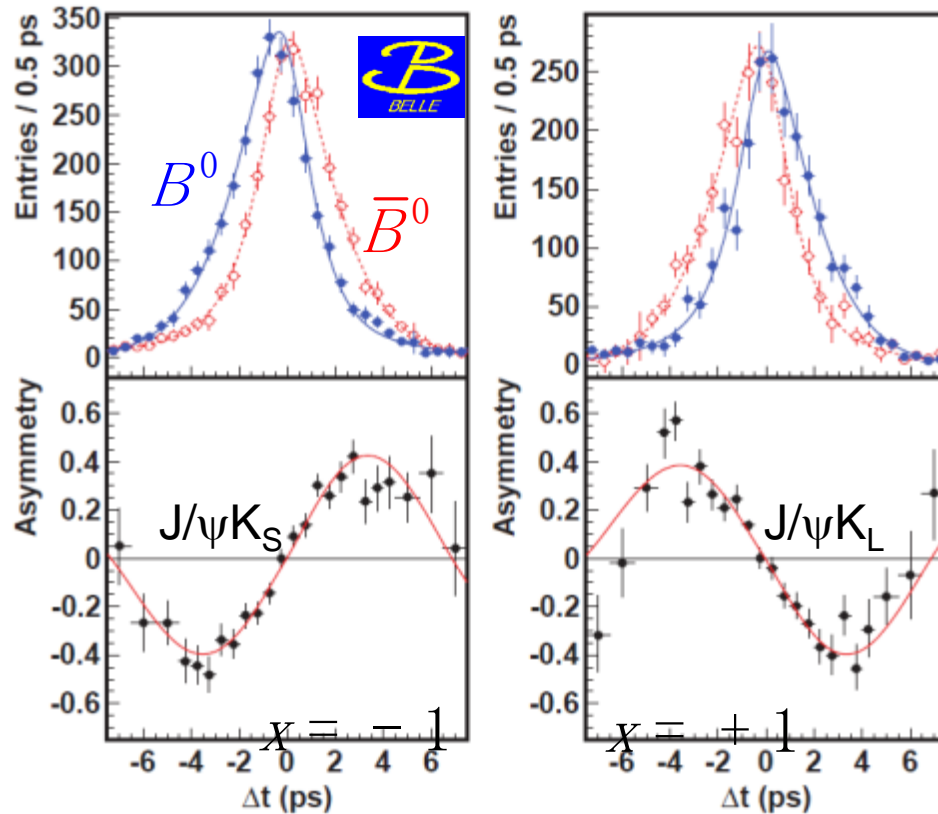


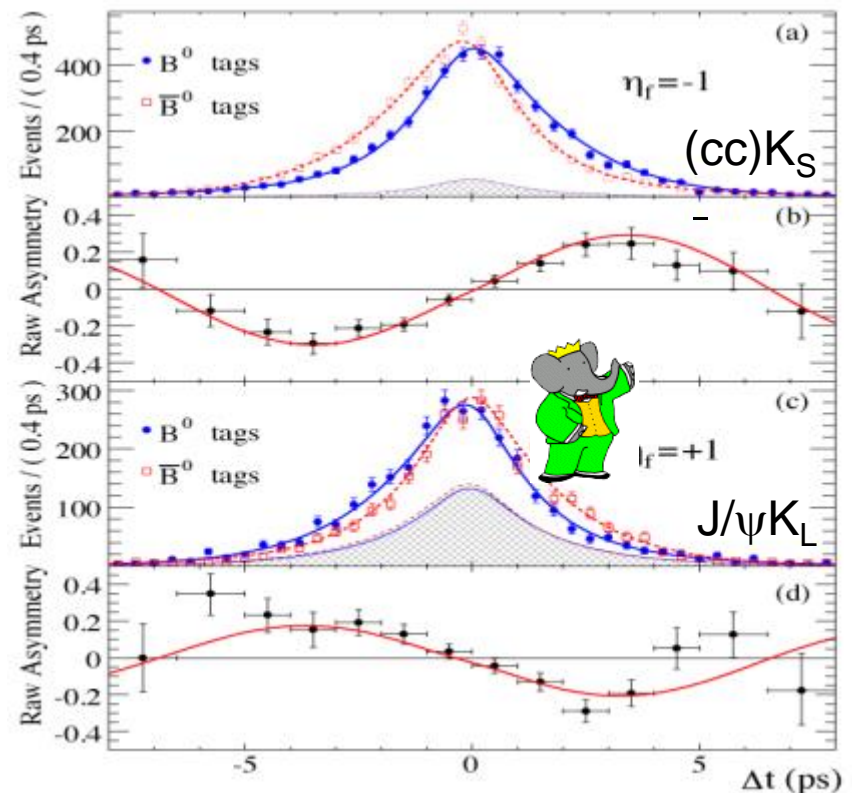
FIG. 1

*Measures the phase of  $V_{td}$  or equivalently the phase of  $B_d$ -anti  $B_d$  mixing.*

# Measurement of $\sin(2\phi_1)/\sin(2\beta)$ in Charmonium $K^0$ modes



$\sin 2\phi_1 = 0.667 \pm 0.023 \pm 0.012$   
 $A_f = 0.006 \pm 0.016 \pm 0.012$   
**PRL108,171802 (2012)**

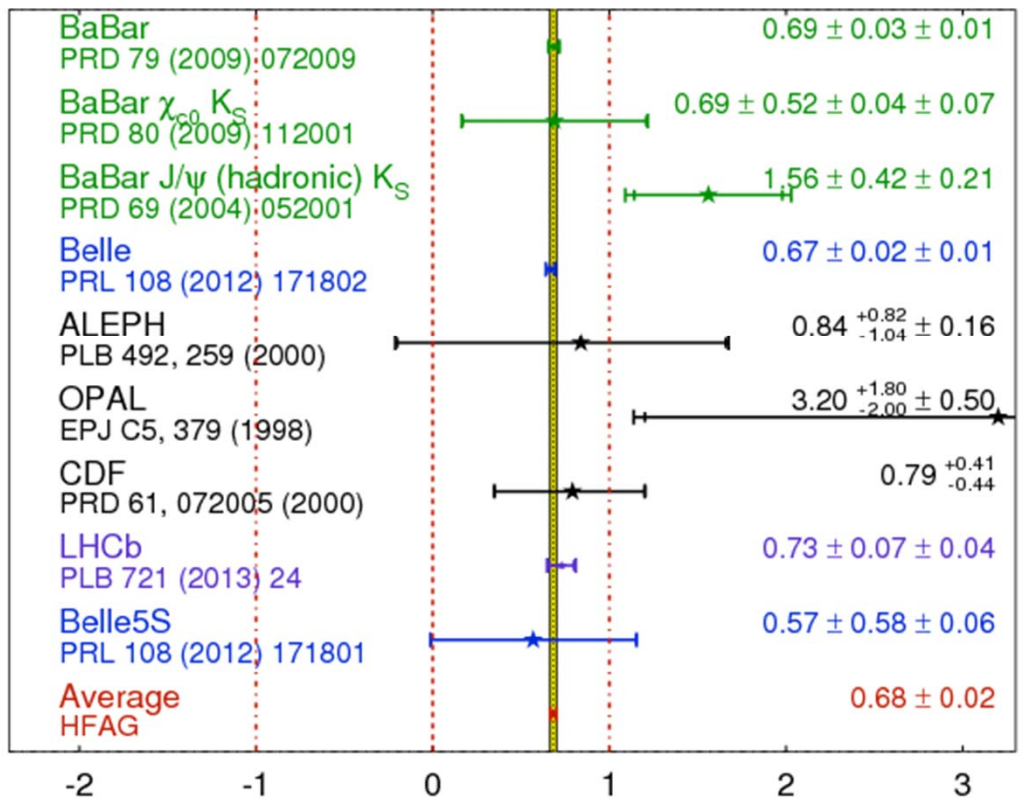


$\sin 2\phi_1 = 0.687 \pm 0.028 \pm 0.012$   
 $A_f = -0.024 \pm 0.020 \pm 0.016$   
**PRD79,072009 (2009)**

Overpowering evidence for CP violation (matter-antimatter asymmetries). >>>> **The phase of  $V_{td}$**  is in good agreement with Standard Model expectations

$$\sin(2\beta) \equiv \sin(2\phi_1)$$

**HFAG**  
Moriond 2014  
PRELIMINARY



B factories: High precision CPV measurement and a calibration for NP.

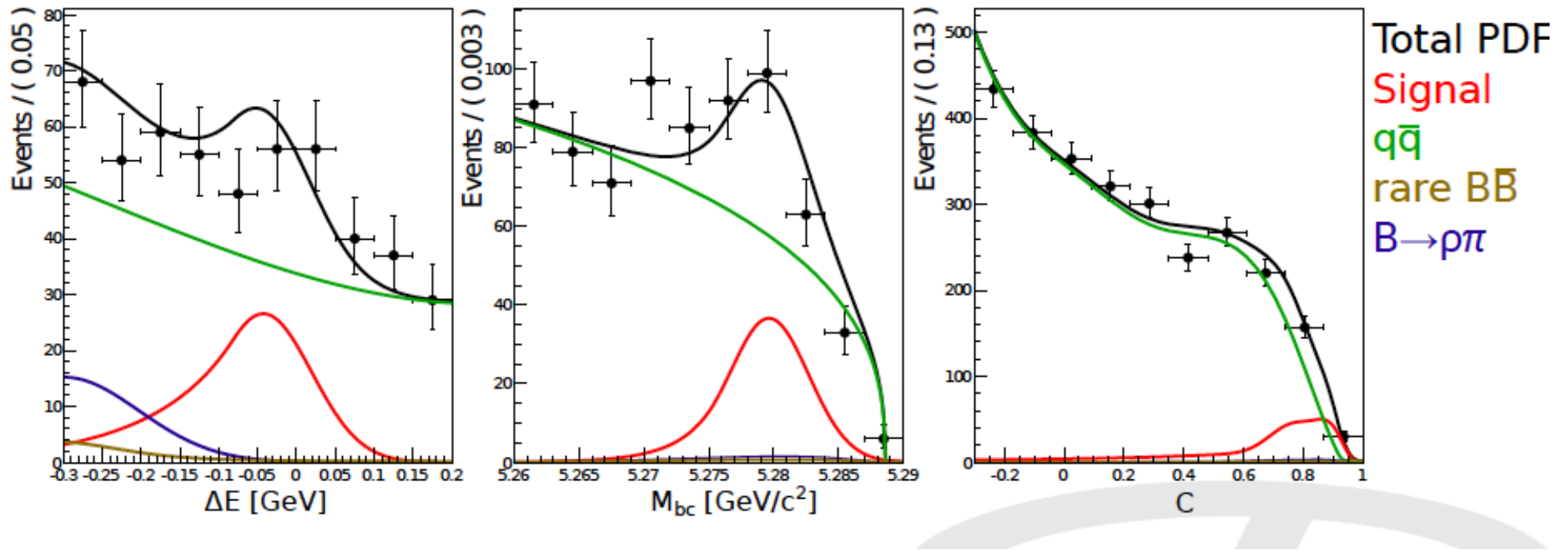
2013: LHCb has joined the game.



J. Bernabeu *et al* point out that by comparing decays with  $t_{tag} > t_{decay}$  and vice versa, one could check for T violation at the B factories. N.B. *This does not assume the CPT theorem.*

BaBar (Valencia) did this and found a  $14\sigma$  signal for T violation !

$B^0 \rightarrow \pi^0 \pi^0$  – fit results (Preliminary)

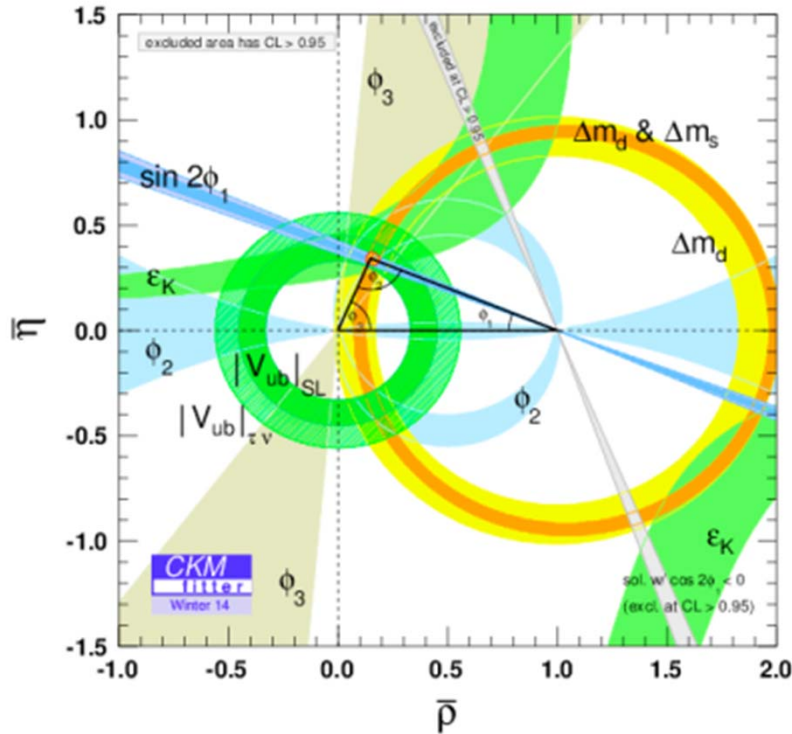


$$\mathcal{B}(B^0 \rightarrow \pi^0 \pi^0) = (0.90 \pm 0.12 \pm 0.10) \times 10^{-6} \quad (6.7\sigma)$$

A difficult mode for hadron colliders

# Results from Global Fits to Data (CKMFitter Group)

Great progress on  $\varphi_3$  or  $\gamma$  (first from B factories and now in the last two years from LHCb (several new results at ICHEP2014)). These measure the phase of  $V_{ub}$

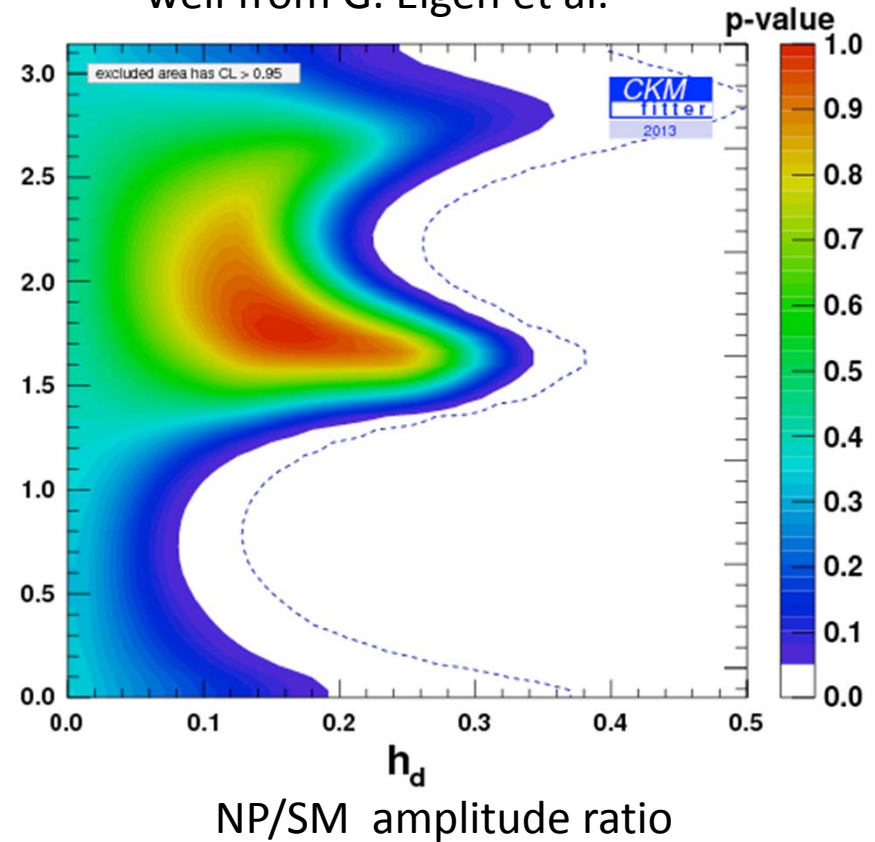


Looks good  
(except for an issue with  $|V_{ub}|$ )

ICHEP2014: Similar results from UTFIT (D. Derkach) as well from G. Eigen et al.

NP  
Phase

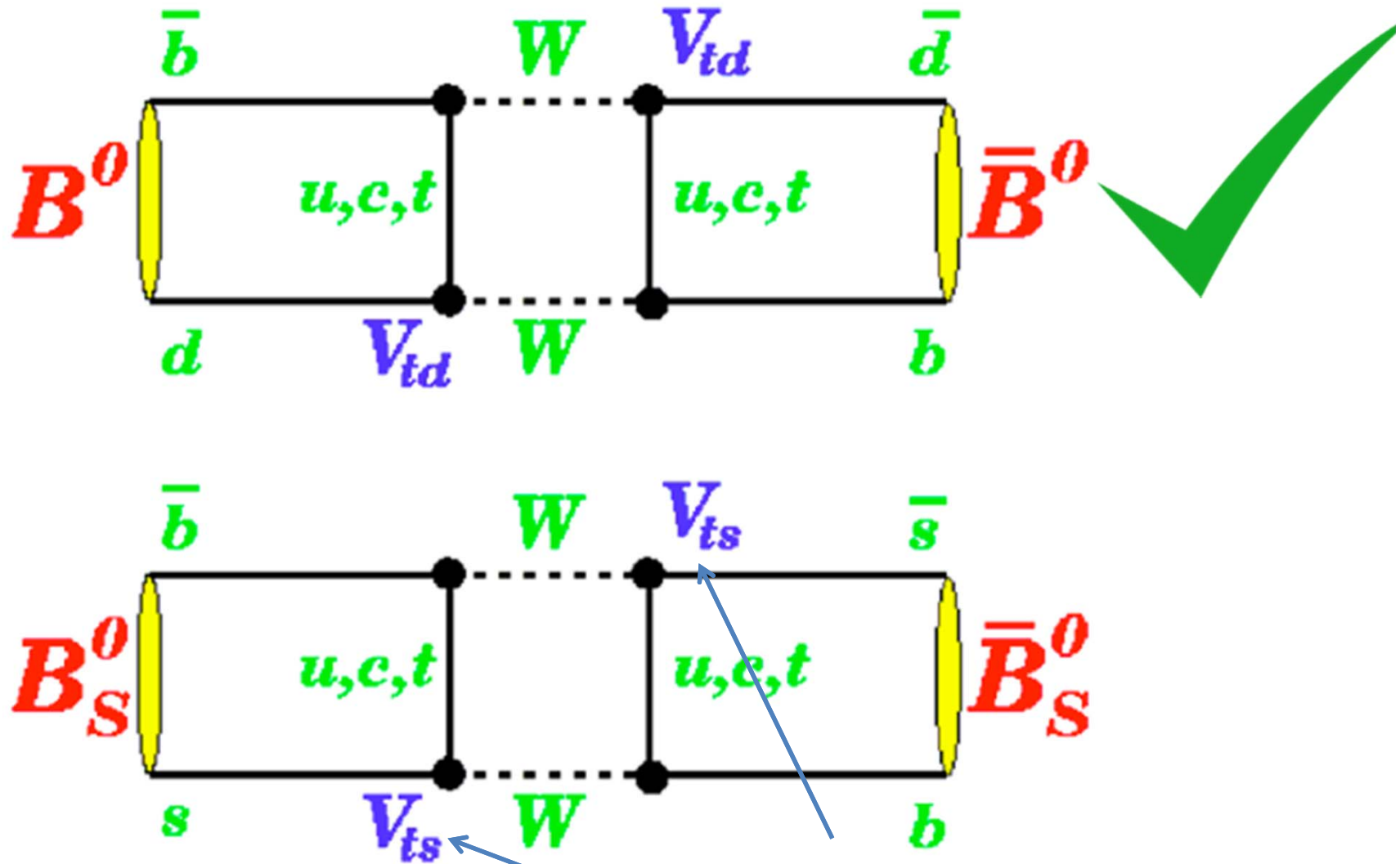
$\rho_d$



But a 10-20% NP amplitude in  $B_d$  mixing is perfectly compatible with all current data.



# Boxes

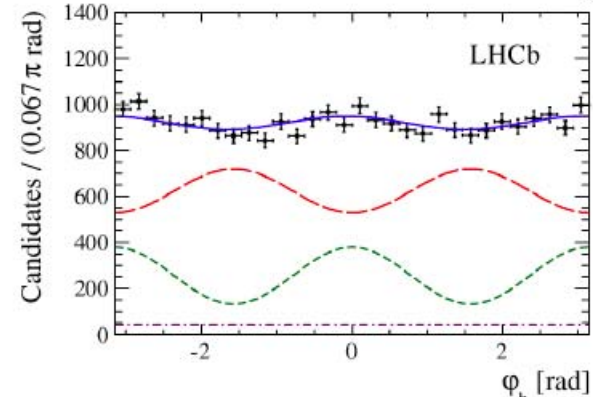
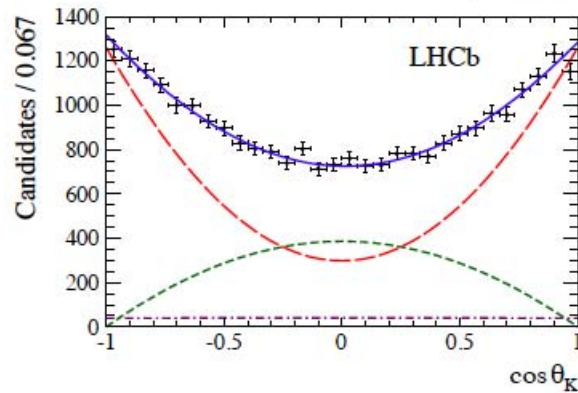
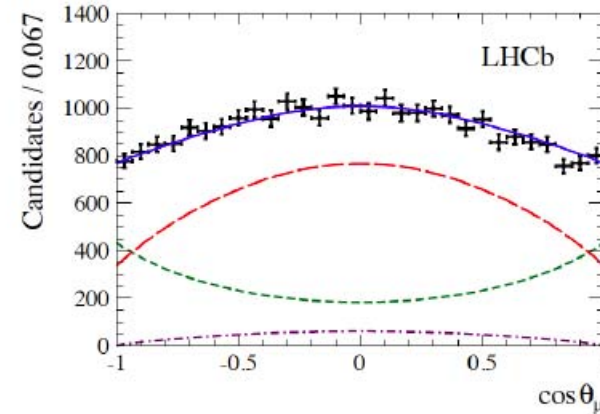
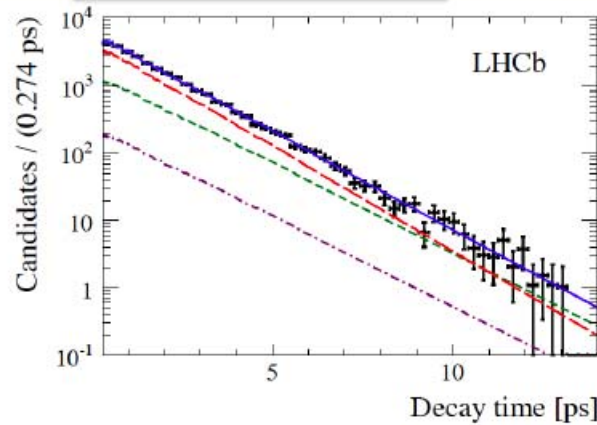


Although B factories can run on the Upsilon(5S), LHCb dominates here

No phase expected from SM but possible from NP particles

# $B_s \rightarrow J/\psi \phi$ angular and decay time projections

PRD 87 (2013)112010, arXiv:1304.2600



$$\Phi_s = +0.07 \pm 0.09 \pm 0.01 \text{ rad}$$

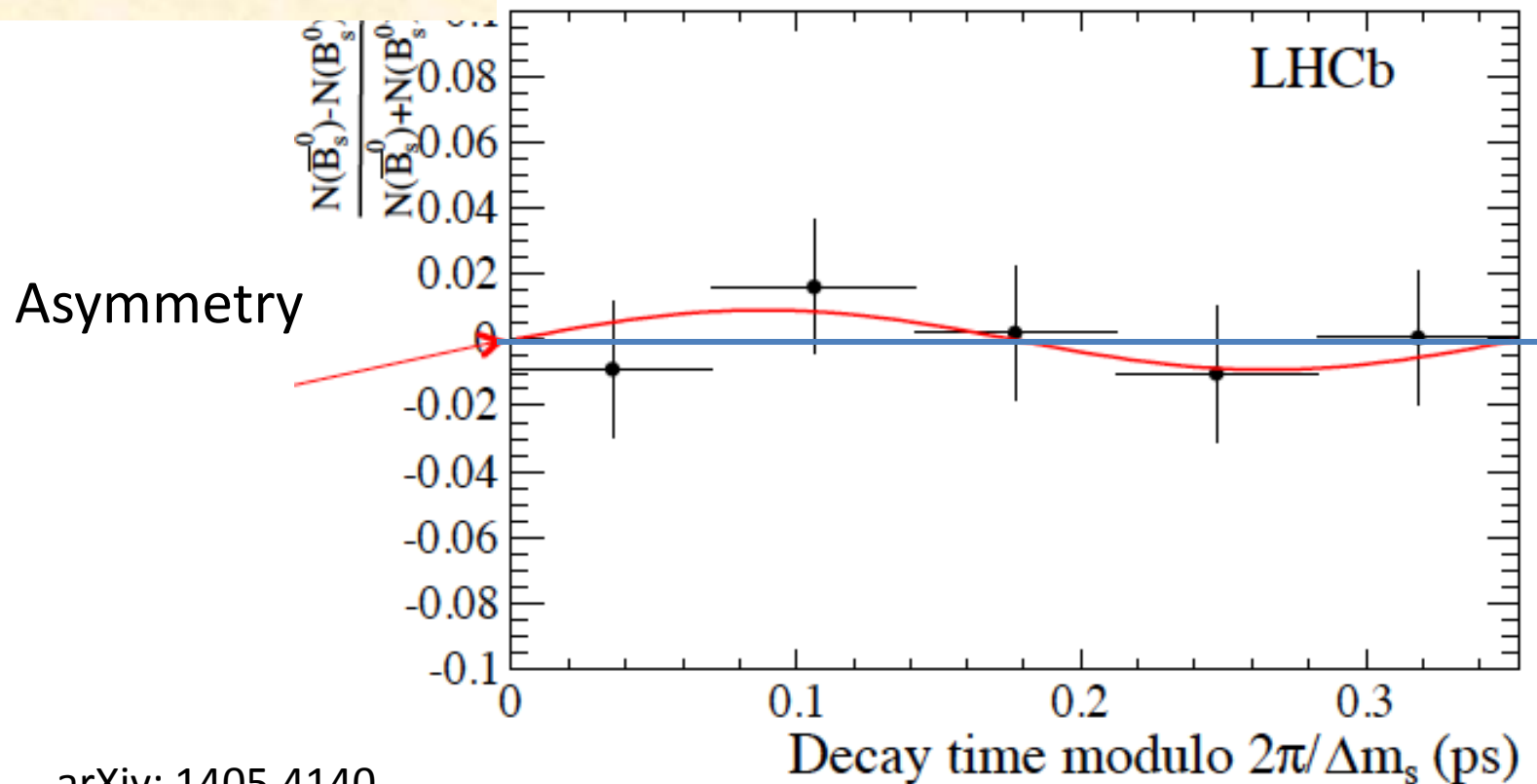
$B_s \rightarrow J/\psi \phi$  is a pseudo-scalar to  $V V$  decay (mixture of two CP eigenstates). This requires multi-dimensional angular analysis.

However,  $B_s \rightarrow J/\psi f_0(980)$  is a pure CP eigenstate since the  $f_0(980)$  is a scalar.

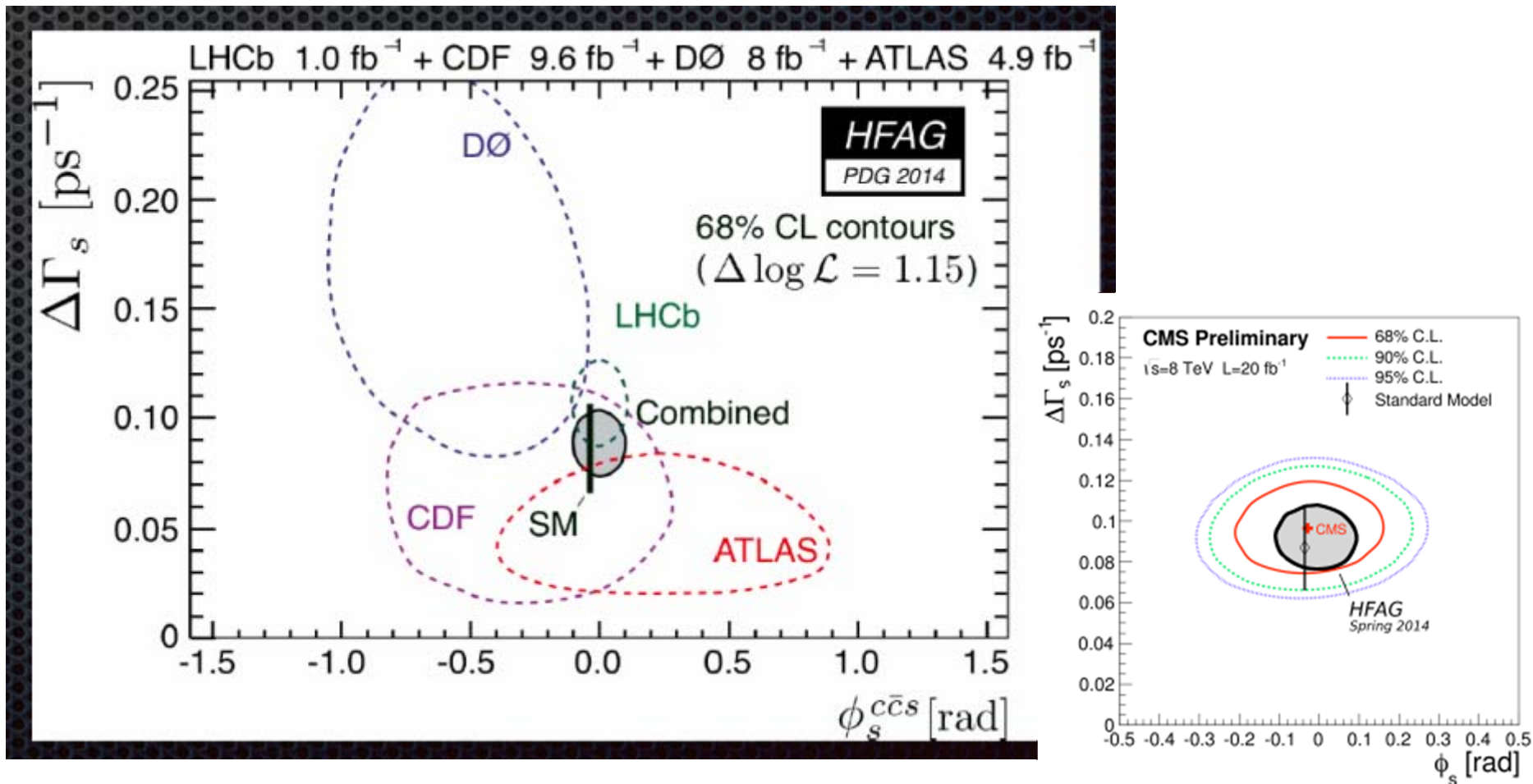
Stone & Zhang pointed out that this mode provides more statistics and a more straightforward analysis. Phys. Rev. D79 (2009) 074024.

$\Phi_s = (70 \pm 68 \pm 8)$   
mrad

Red curve: expectation for  $\Phi_s = 70$  mrad



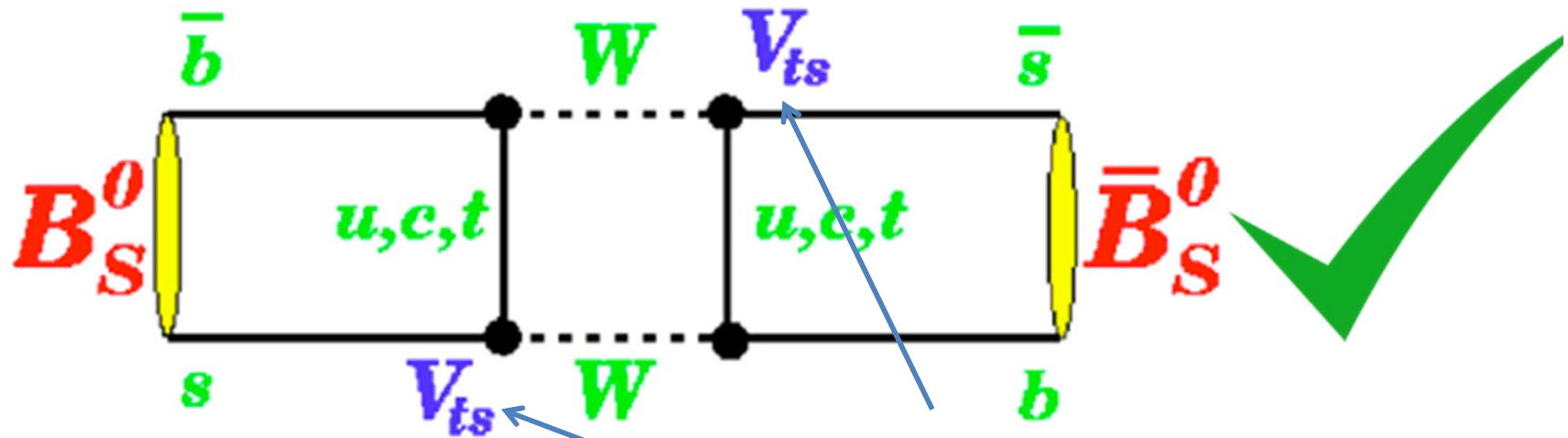
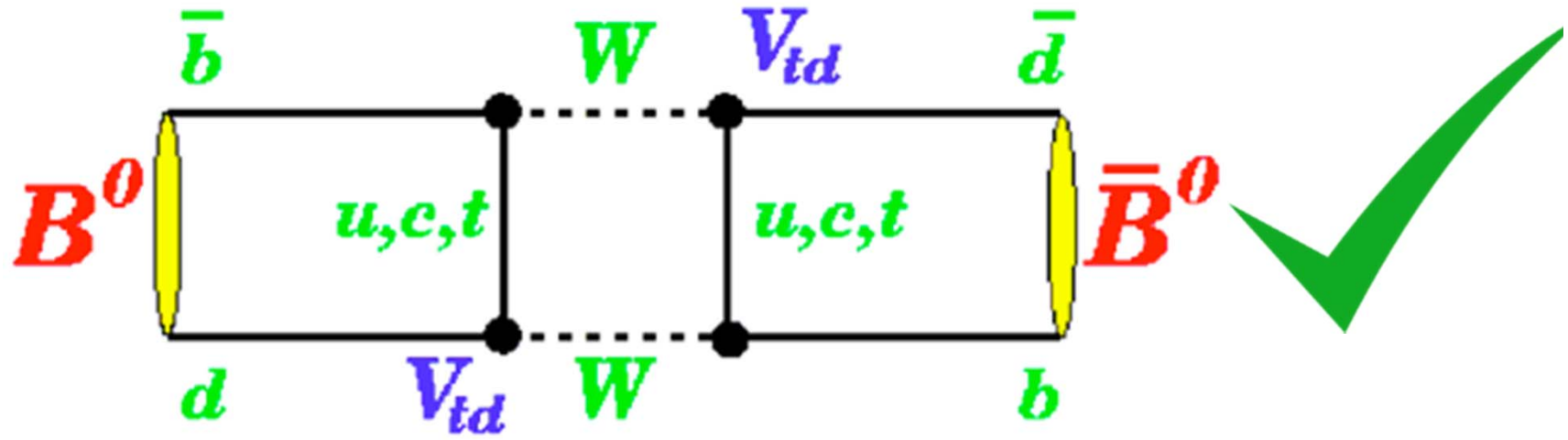
# Results on the phase of $B_s$ -anti $B_s$ mixing (i.e. phase of $V_{ts}$ )



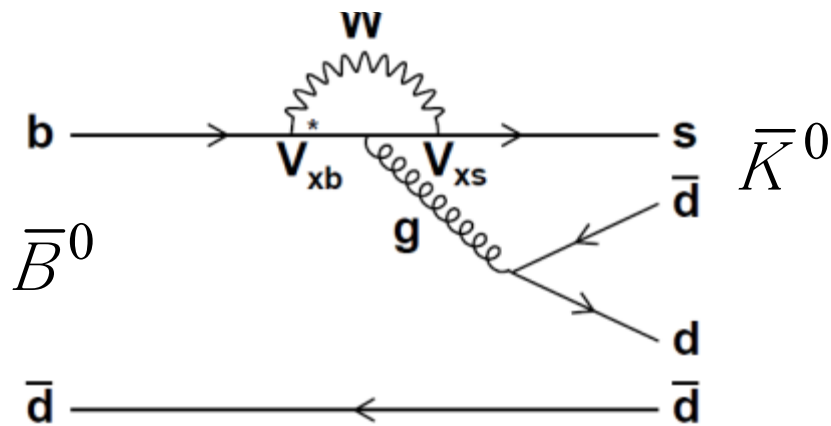
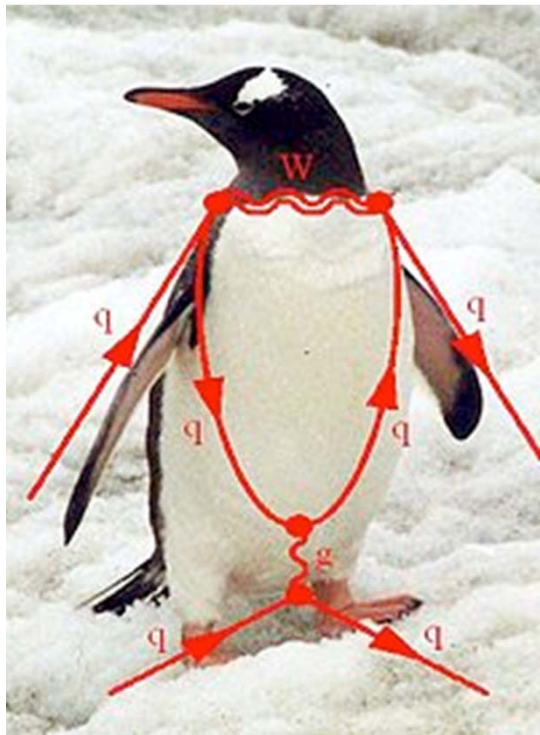
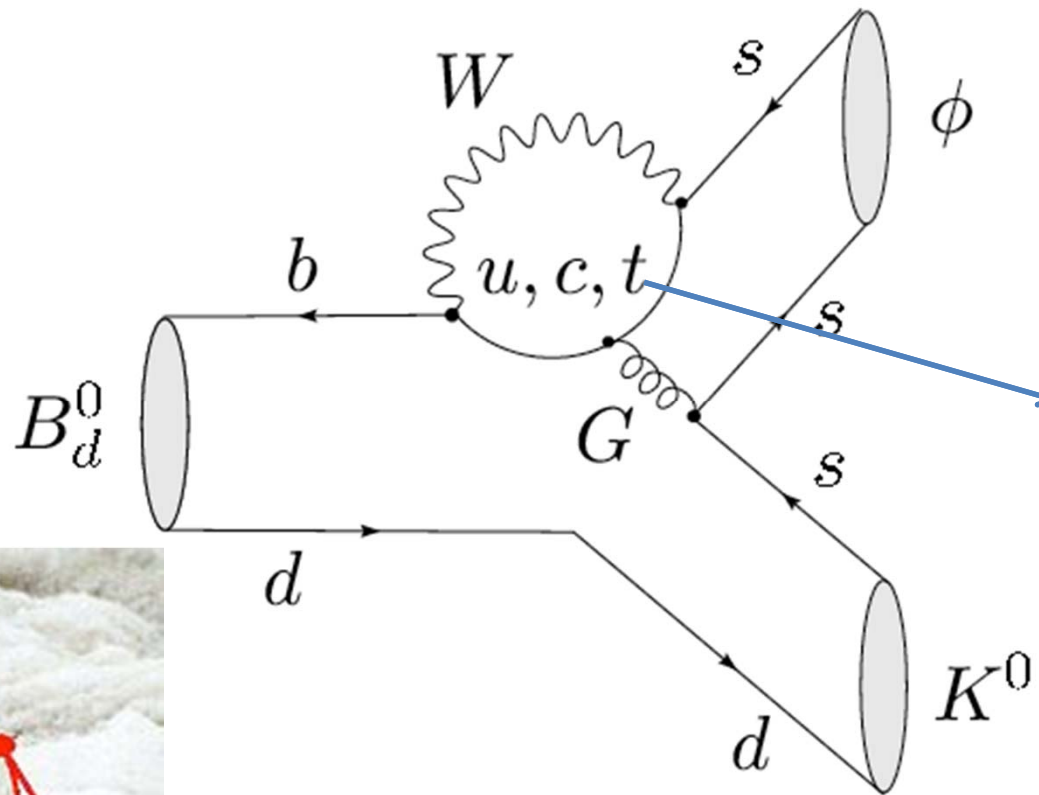
R. Aaij: Does not include arXiv: 1405.4140 using  $J/\psi\pi\pi$  decays, which are dominantly S-wave.

Does not include new CMS result at ICHEP2014 (G. Fedi)

# Boxes



No phase expected from SM but possible from NP particles



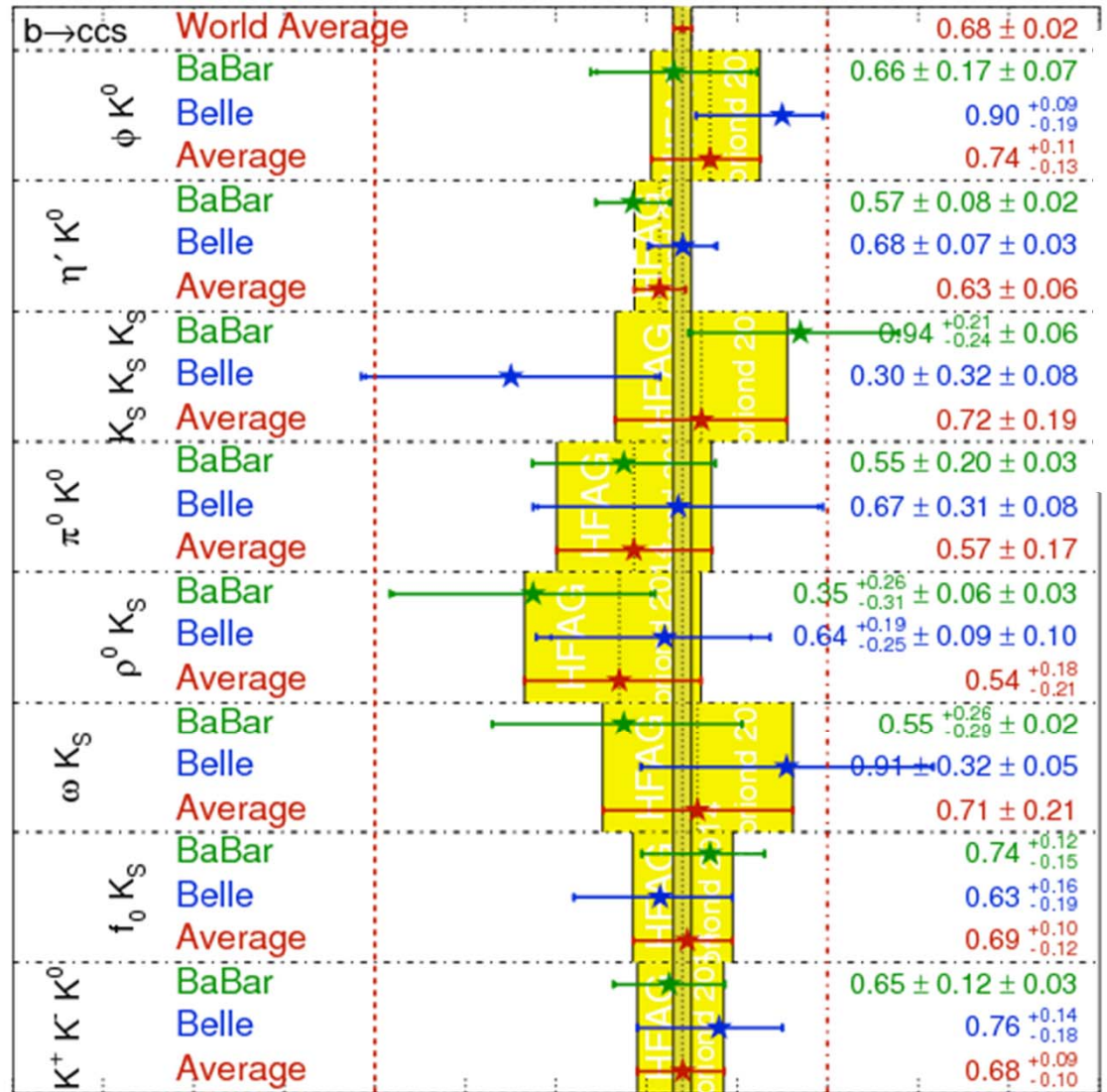
Phase of  $V_{ts}$

ICHEP2014: New Belle results on  $B \rightarrow \omega K_S$ ,  $B \rightarrow \eta' K_S$

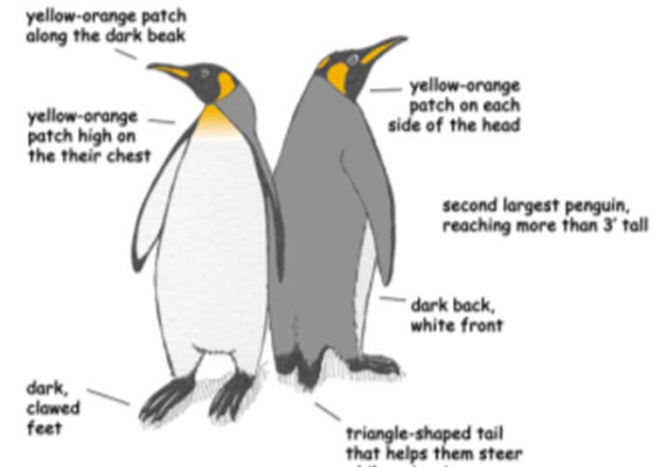
Talk by  
V.Chobanova

$$\sin(2\beta^{\text{eff}}) \equiv \sin(2\phi_1^{\text{eff}})$$

**HFAG**  
Moriond 2014  
PRELIMINARY



## New Physics Phases in Penguin $b \rightarrow s$ decays

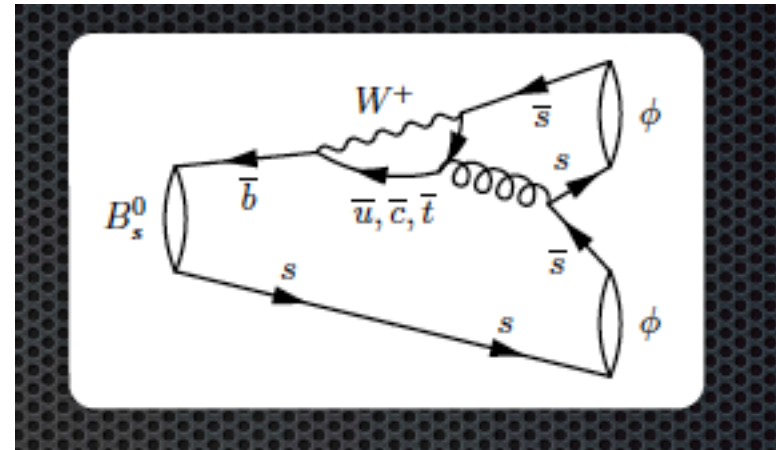


No evidence for NP at current level of sensitivity

LHCb is absent from this game (lower  $K_S$  eff and flavor tagging eff) but contributes in  $B_s$  modes.

M. Needham

But LHCb dominates these  $B_s$  modes  
Again, NP Penguin phase is found to be consistent with zero



## $B_s \rightarrow \phi\phi$ - Time-Dependent Results

Results are found to be:

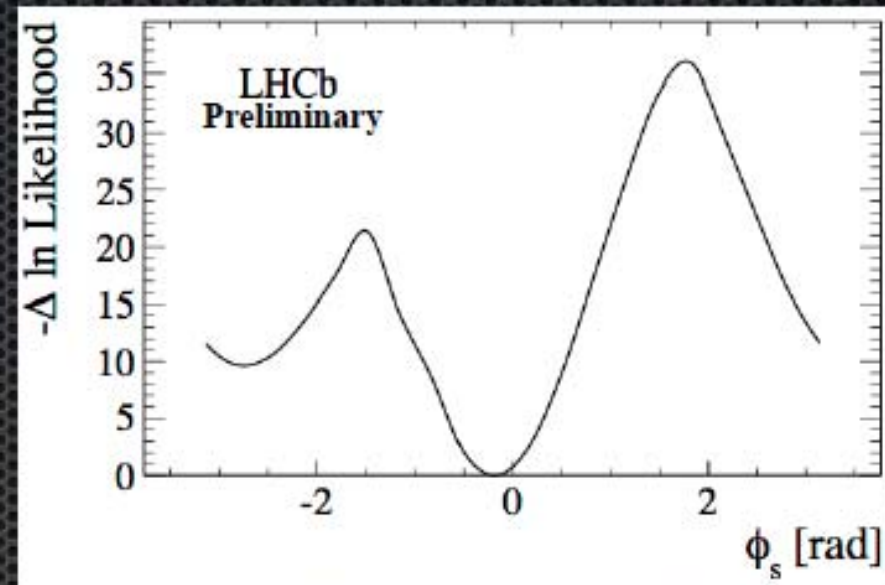
$$\phi_s = -0.17 \pm 0.15 \pm 0.03$$

$$\lambda = 1.04 \pm 0.07 \pm 0.03$$

$$|A_0|^2 = 0.364 \pm 0.012 \pm 0.009$$

$$|A_{\pm}|^2 = 0.305 \pm 0.013 \pm 0.005$$

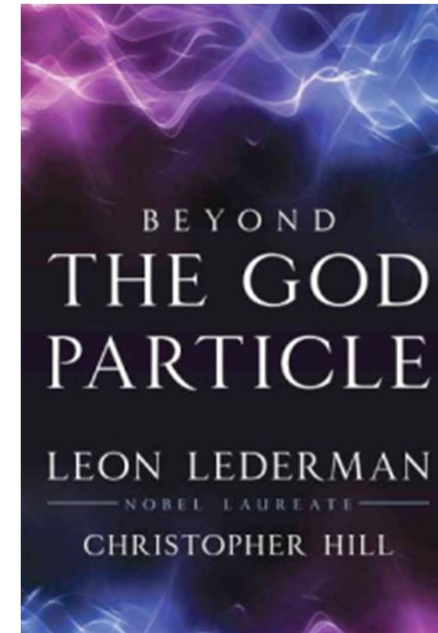
First uncertainty statistical and  
second systematic



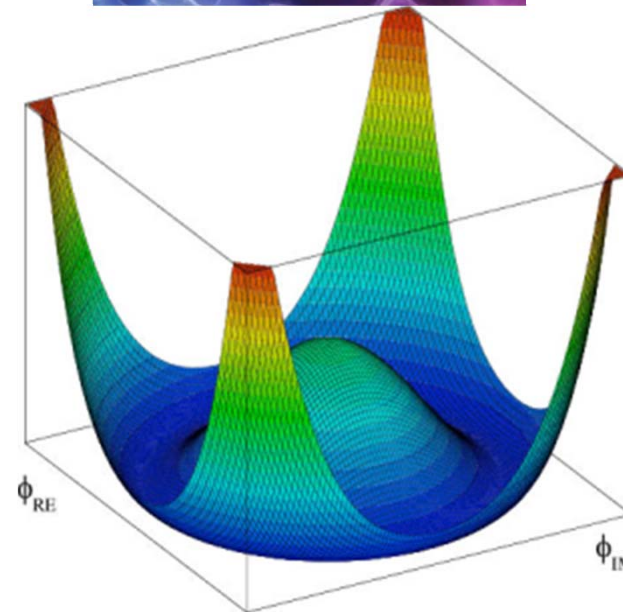


# “Missing Energy” Decays

The BEH boson is now firmly established by experimental results from ATLAS and CMS.



Does the GP (Brout-Englert-Higgs particle) have a “brother” i.e. the charged Higgs ?

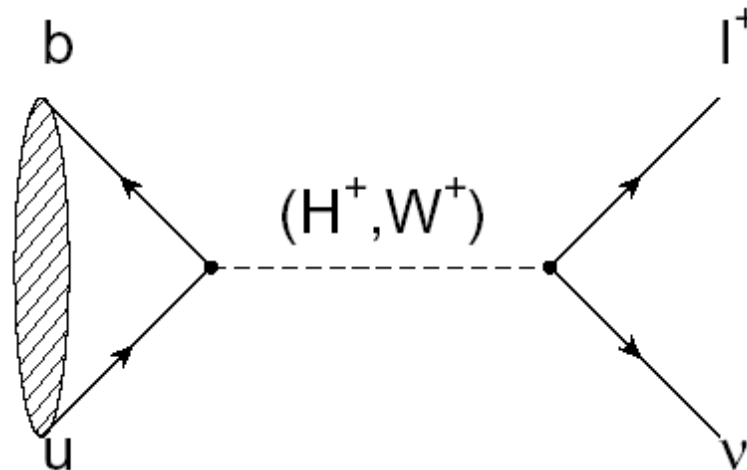


Measurements at B factories and direct searches at hadron colliders take *complementary* approaches to this important question.

$$B^+ \rightarrow \tau^+ \nu_\tau$$

(Decays with *Large Missing Energy*)

Sensitivity to new physics from a charged Higgs



$$\mathcal{B}(B^+ \rightarrow \tau^+ \nu_\tau) = \frac{G_F^2 m_B}{8\pi} m_\tau^2 \left(1 - \frac{m_\tau^2}{m_B^2}\right)^2 f_B^2 |V_{ub}|^2 \tau_B$$

$$\mathcal{B}_{(B \rightarrow \tau \nu)} = \mathcal{B}_{SM} \times \left(1 - \tan^2 \beta \frac{m_{B^\pm}^2}{m_{H^\pm}^2}\right)$$



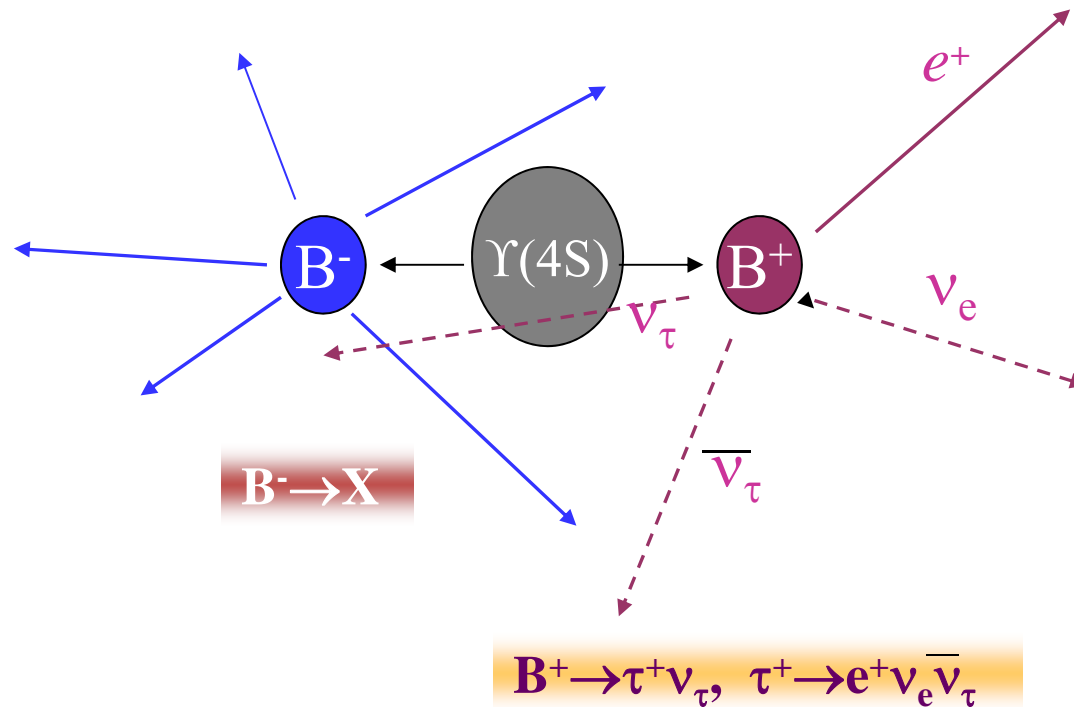
The *B* meson decay constant, determined by the *B* wavefunction at the origin

( $|V_{ub}|$  taken from indep. measurements.)

# Consumer's guide to charged Higgs

- Higgs doublet of type I (couples with equal strength to upper and lower generations)
- Higgs doublet of type II (couples with different strength to u and d-type quarks,  $\tan(\beta) = v_u/v_d$  ( favored NP scenario e.g. MSSM))
- Higgs doublet of type III (not type I or type II; anything goes)

# Why measuring $B \rightarrow \tau \nu$ is non-trivial

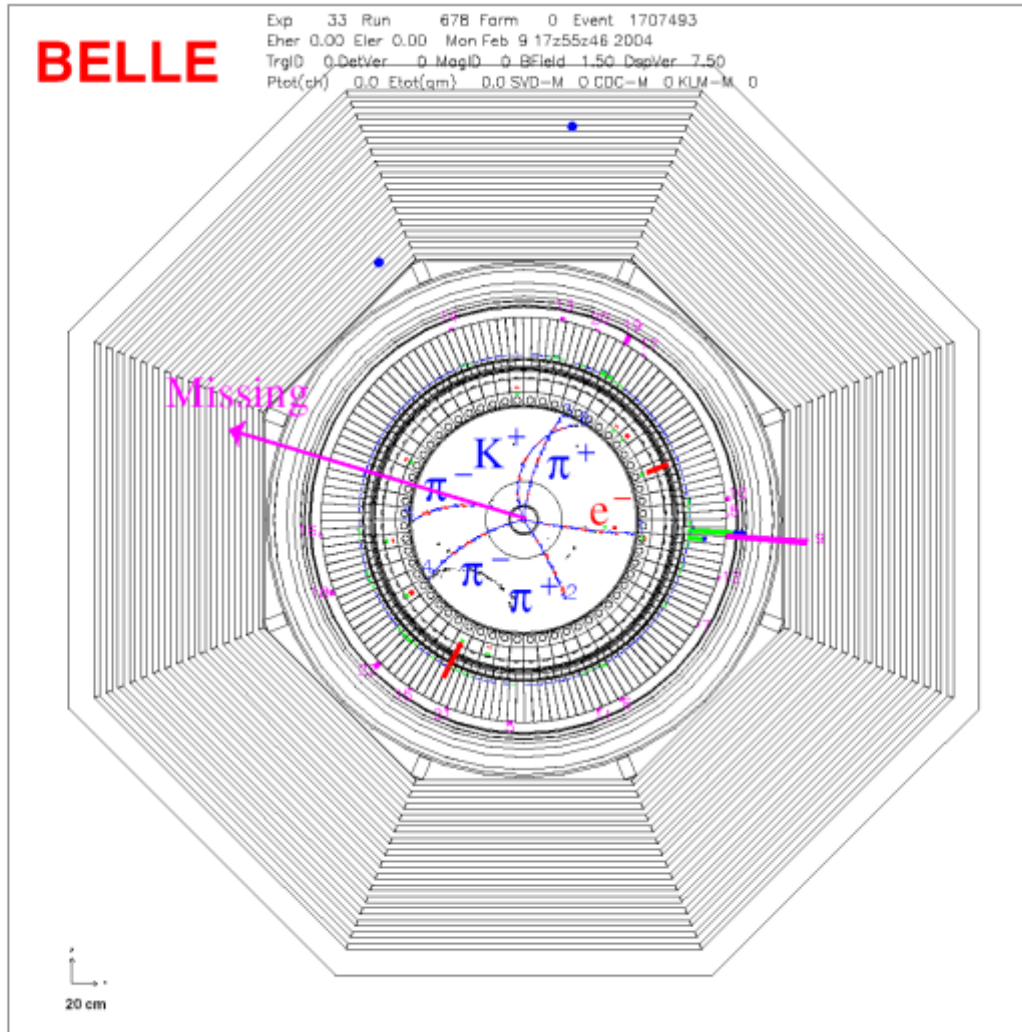
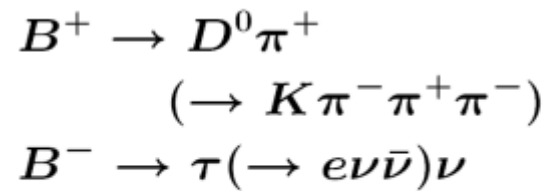


Most of the sensitivity is from tau modes with 1-prong

*The experimental signature is rather difficult:  
B decays to a **single charged track + nothing***

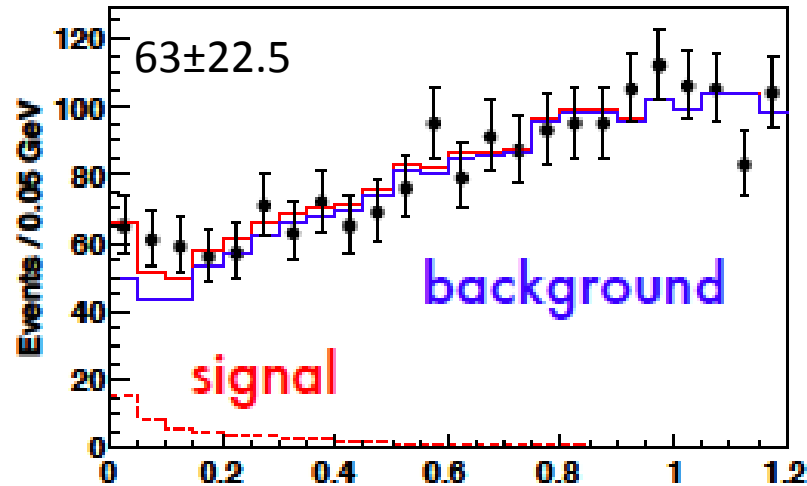
(This may be hard at a hadron collider)

# Example of a Missing Energy Decay ( $B^- \rightarrow \tau^- \nu_\tau$ ) in Data



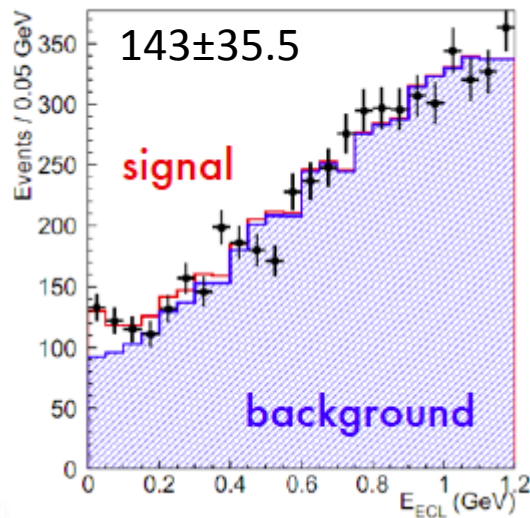
The clean e+e- environment makes this possible

# Example: Belle measurement with full data sample and hadronic tags



Idea: With the “single B meson beam”, we look for a single track from a  $\tau$ , missing energy/momentum and extra calorimeter energy close to zero.

Semileptonic tags



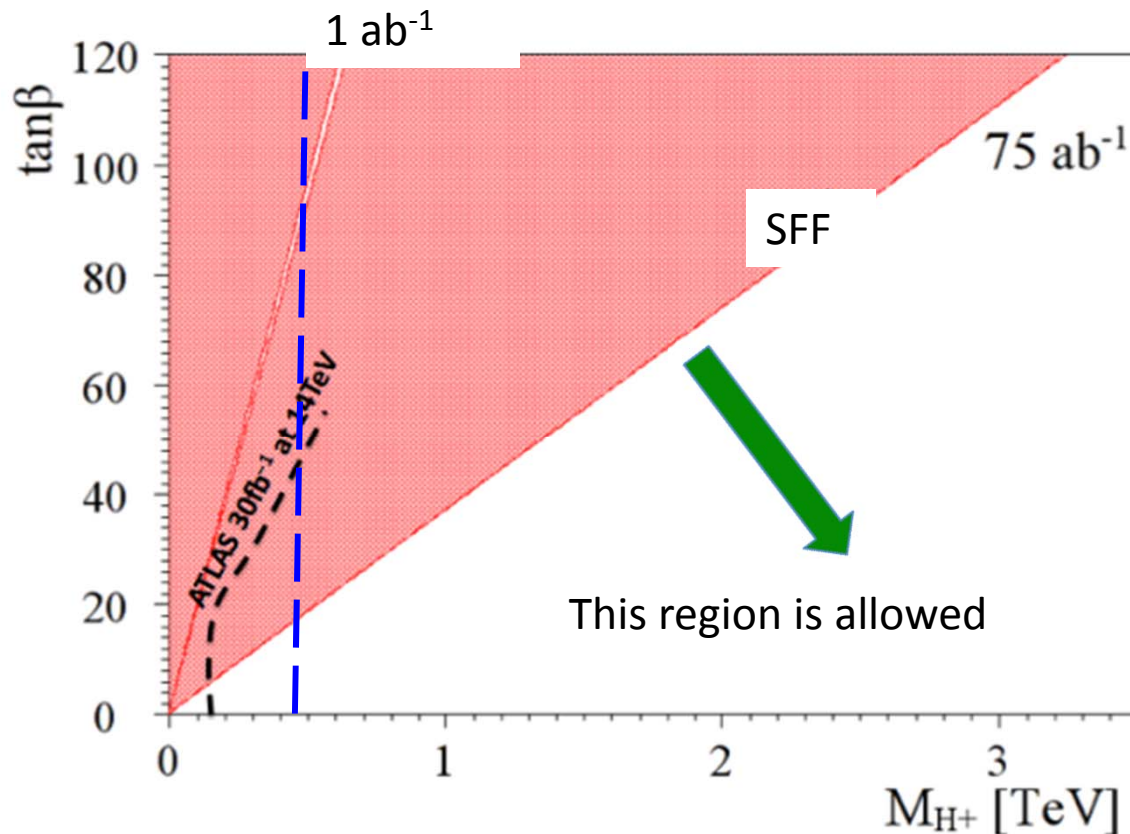
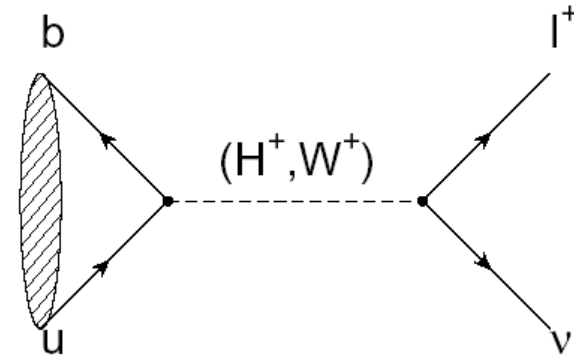
With the full B factory statistics only “evidence”. No single observation from either Belle or BaBar.

The horizontal axis is the “Extra Calorimeter Energy”

# Complementarity of $e^+ e^-$ factories and LHC

(Slide adapted from A. Bevan)

The current combined limit places a stronger constraint than direct searches from LHC exps. for the next few years.

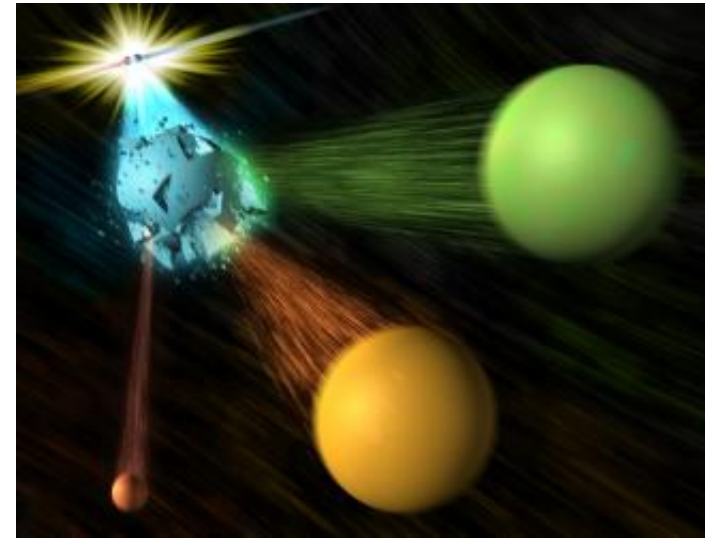
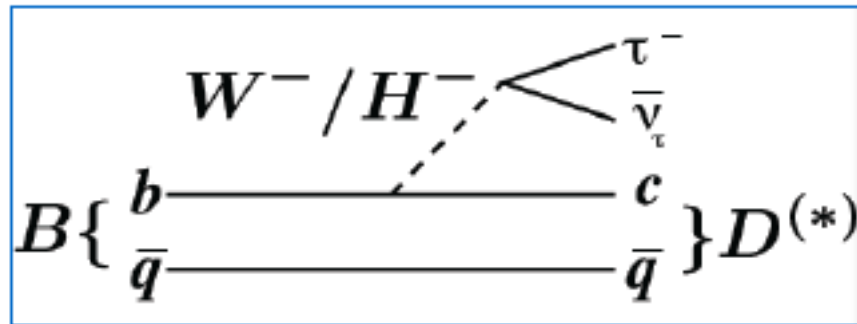


$$r_H = \left( 1 - \frac{m_B^2}{m_H^2} \tan^2 \beta \right)^2$$

Currently **inclusive b to sy** rules out  $m_{H^+}$  below  $\sim 400 \text{ GeV}/c^2$  range (independent of  $\tan\beta$ )

<http://arxiv.org/abs/1208.2788>





$$\mathcal{R}(D^{(*)})_{2\text{HDM}} = \mathcal{R}(D^{(*)})_{\text{SM}} + A_{D^{(*)}} \frac{\tan^2 \beta}{m_{H^+}^2} + B_{D^{(*)}} \frac{\tan^4 \beta}{m_{H^+}^4}$$

	$D\tau\nu$	$D^*\tau\nu$
$A_{D^{(*)}} \text{ (GeV}^2\text{)}$	$-3.25 \pm 0.32$	$-0.230 \pm 0.029$
$B_{D^{(*)}} \text{ (GeV}^4\text{)}$	$16.9 \pm 2.0$	$0.643 \pm 0.085$

$$R(D^{(*)}) = \frac{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \tau^- \bar{\nu}_\tau)}{\mathcal{B}(\bar{B} \rightarrow D^{(*)} \ell^- \bar{\nu}_\ell)} \begin{matrix} \longrightarrow & \text{Signal} \\ \longrightarrow & \text{Normalization } (\ell = e \text{ or } \mu) \end{matrix}$$

Example from a recent BaBar paper

Signals in  $B \rightarrow D^{(*)} \tau \nu$  ( $489 \pm 63$ ,  $888 \pm 63$ )

Missing mass variable:

$$m_{\text{miss}}^2 = p_{\text{miss}}^2 = (p[e^+e^-] - p_{\text{tag}} - p_{D^{(*)}} - p_l)^2$$

$P_l^*$  = momentum of lepton in B rest frame

*Production of B meson pairs at threshold is critical to the separation of backgrounds from the missing energy/momentum signal.*

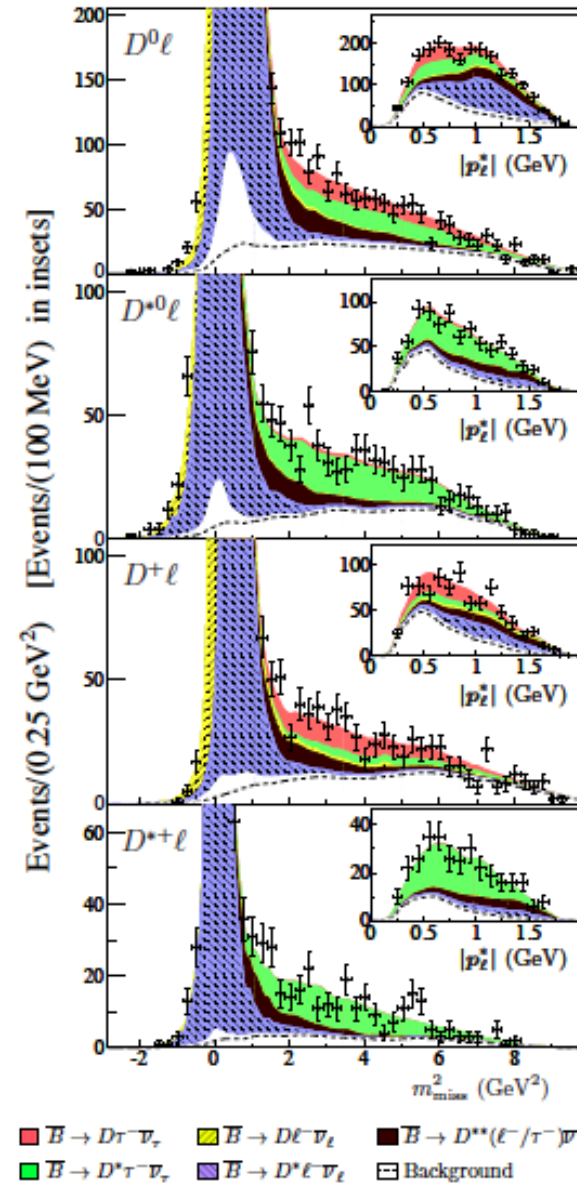


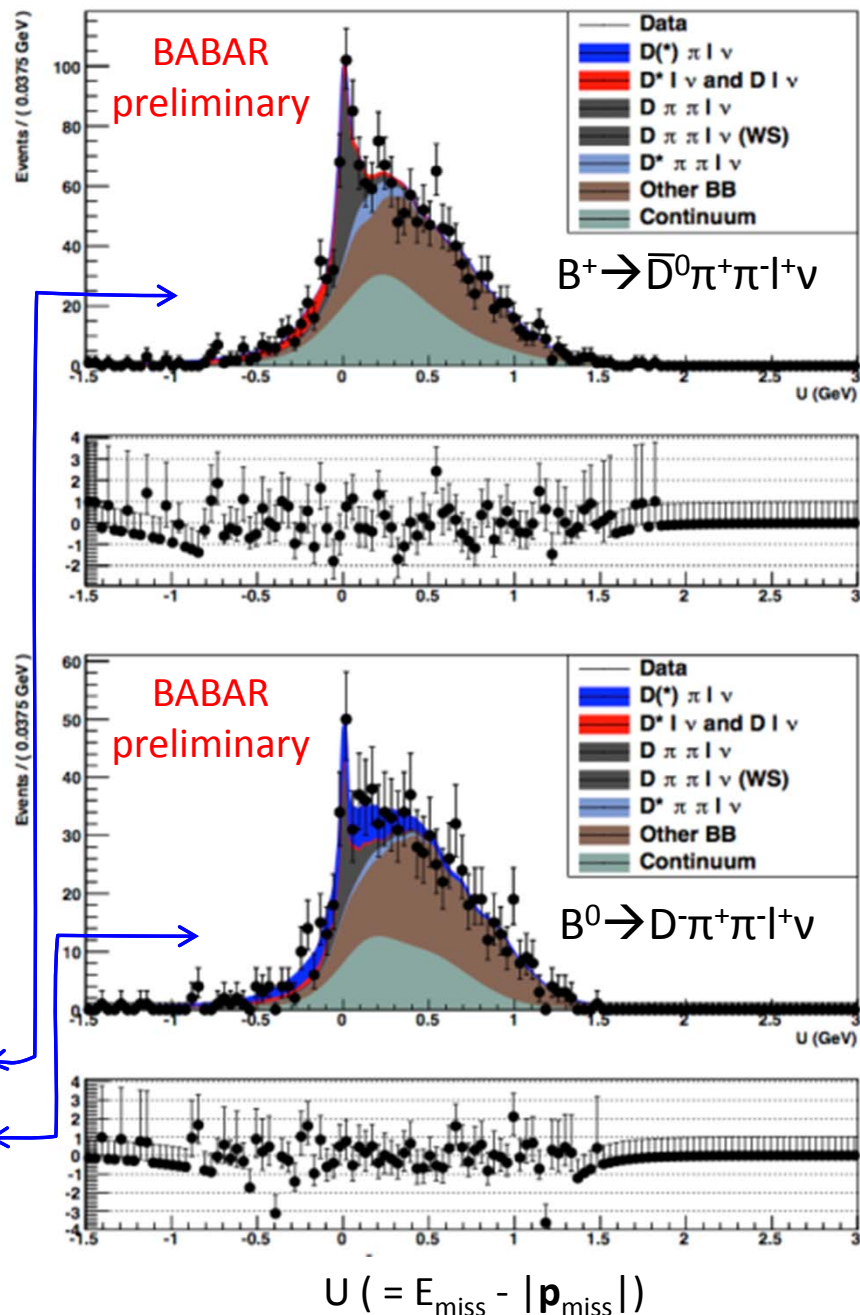
FIG. 1. (Color online) Comparison of the data and the fit projections for the four  $D^{(*)} \ell$  samples. The insets show the  $|p_l^*|$  projections for  $m_{\text{miss}}^2 > 1 \text{ GeV}^2$ , which excludes most of the normalization modes. In the background component, the region above the dashed line corresponds to charge cross-feed, and the region below corresponds to continuum and  $B\bar{B}$ .



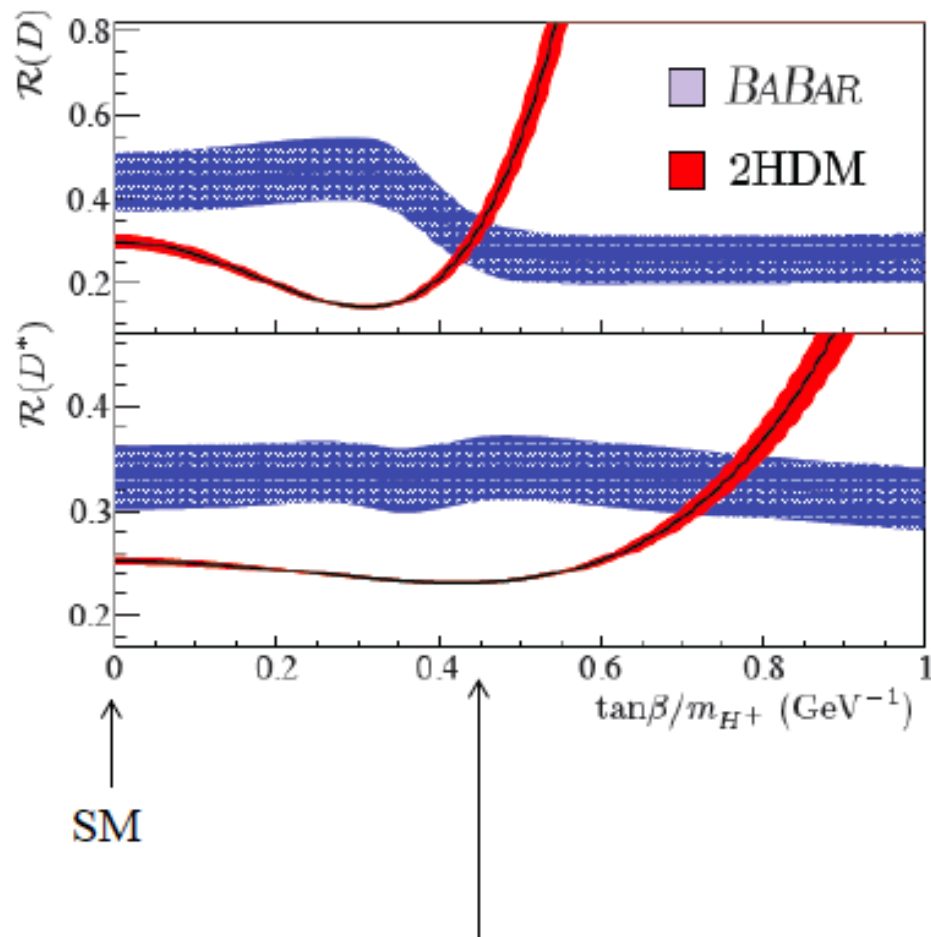
# BABAR@ICHEP2014: multi-body semi-leptonic $b \rightarrow c$ decays

- Background for  $B \rightarrow D^{(*)} \tau \nu$  analysis
- Fully reconstruct tag B and semileptonic signal decays,  
 $B \rightarrow D^{(*)} \pi^+ \pi^- l \nu$
- Averaging  $B^+$  and  $B^0$ , significance is  $5.1\sigma$  for  $D\pi^+\pi^-l\nu$ ,  $3.5\sigma$  for  $D^*\pi^+\pi^-l\nu$
- $B \rightarrow D^{(*)} \pi \pi l \nu$  decays (including  $\pi^0$ ) fill  $\sim 1/2$  of gap  $B(b \rightarrow c l \nu) - \sum B(B \rightarrow X_c l \nu)$
- Parallel contribution by Thomas Lück

Mode	Signal	BF $\times 10^4$ (quoted for $B^+$ )	Signif
$D^0 \pi^+ \pi^-$	$189 \pm 39$	$21.0 \pm 4.5 \pm 2.8 \pm 1.2$	$4.4\sigma$
$D^+ \pi^+ \pi^-$	$57 \pm 20$	$11.4 \pm 4.1 \pm 2.6 \pm 0.6$	$2.6\sigma$
$D^{*0} \pi^+ \pi^-$	$75 \pm 36$	$9.2 \pm 4.4 \pm 2.0 \pm 0.3$	$1.9\sigma$
$D^{*+} \pi^+ \pi^-$	$58 \pm 19$	$13.2 \pm 4.5 \pm 2.4 \pm 0.3$	$2.9\sigma$



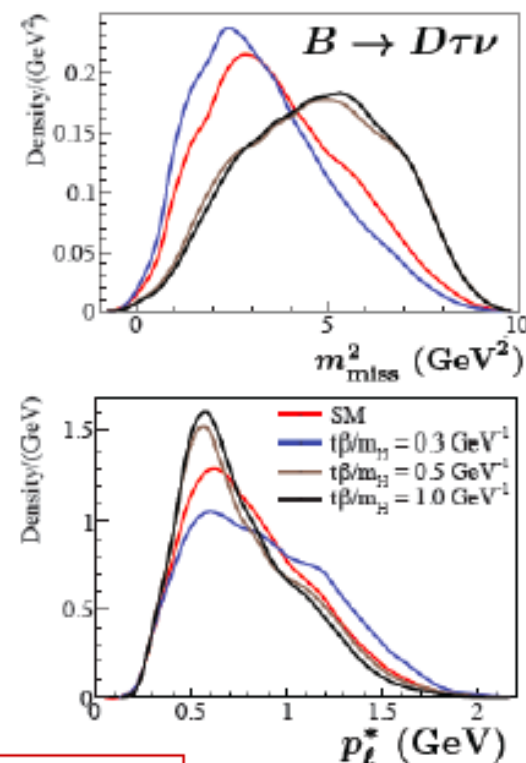
# Limits on type-II 2HDM



$$\tan\beta/m_{H^+} = 0.44 \pm 0.02 \text{ GeV}^{-1}$$

$$\tan\beta/m_{H^+} = 0.75 \pm 0.04 \text{ GeV}^{-1}$$

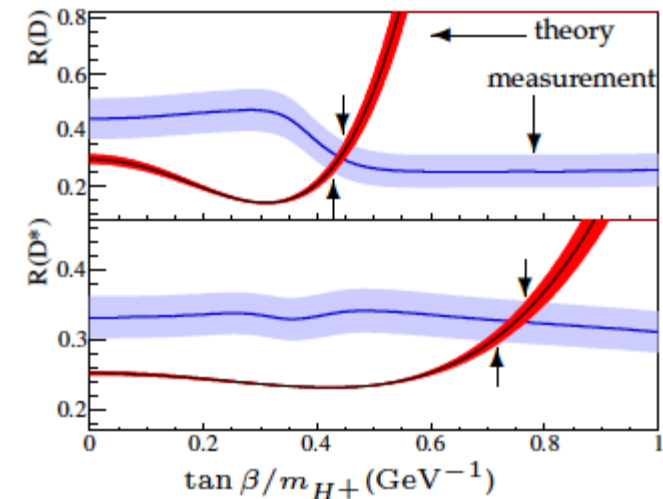
2HDM modifies fit-variable distribution and hence the efficiency



Best point is  $\tan\beta/m_{H^+} = 0.45 \text{ GeV}^{-1}$ , excluded at 99.8% CL ( $3.1 \sigma$ ).  
All other values (with  $m_{H^+} > 15 \text{ GeV}$ ) are worse.

BaBar collaboration, Phys. Rev. Lett. 109, 101802 (2012)

“However, the combination of  $R(D)$  and  $R(D^*)$  excludes the type II 2HDM charged Higgs boson with a 99.8% confidence level for any value of  $\tan(\beta)/m_{H^+}$ ”

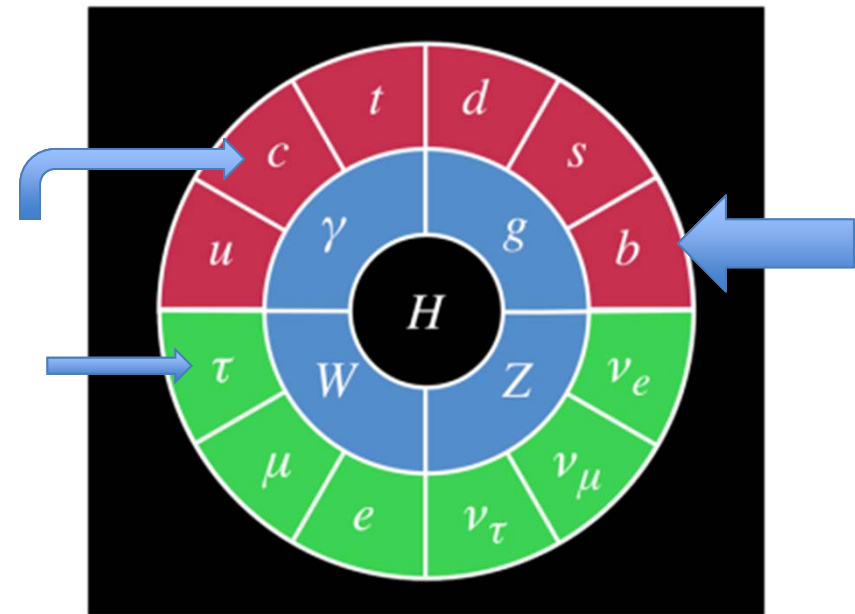


In other words, found NP but have *killed* the 2HDM NP model.

A few points:

- (0) Still waiting for the Belle update (PRD 82,.072005 (2010))
- (1) If the R values fluctuate back down in the future, this *will allow NP from type II charged Higgs to be viable.*
- (2) It is *obvious* that we need two orders of magnitude of data to solve these issues related to the charged Higgs.

2014 is the 50<sup>th</sup> anniversary of the discovery of CP violation in the kaon sector [see <http://pprc.qmul.ac.uk/research/50-years-cp-violation>



## The Next Generation

### *Belle II and the LHCb upgrade*

US P5 report (p. v): “Explore the unknown: new particles, interactions, and physical principles”

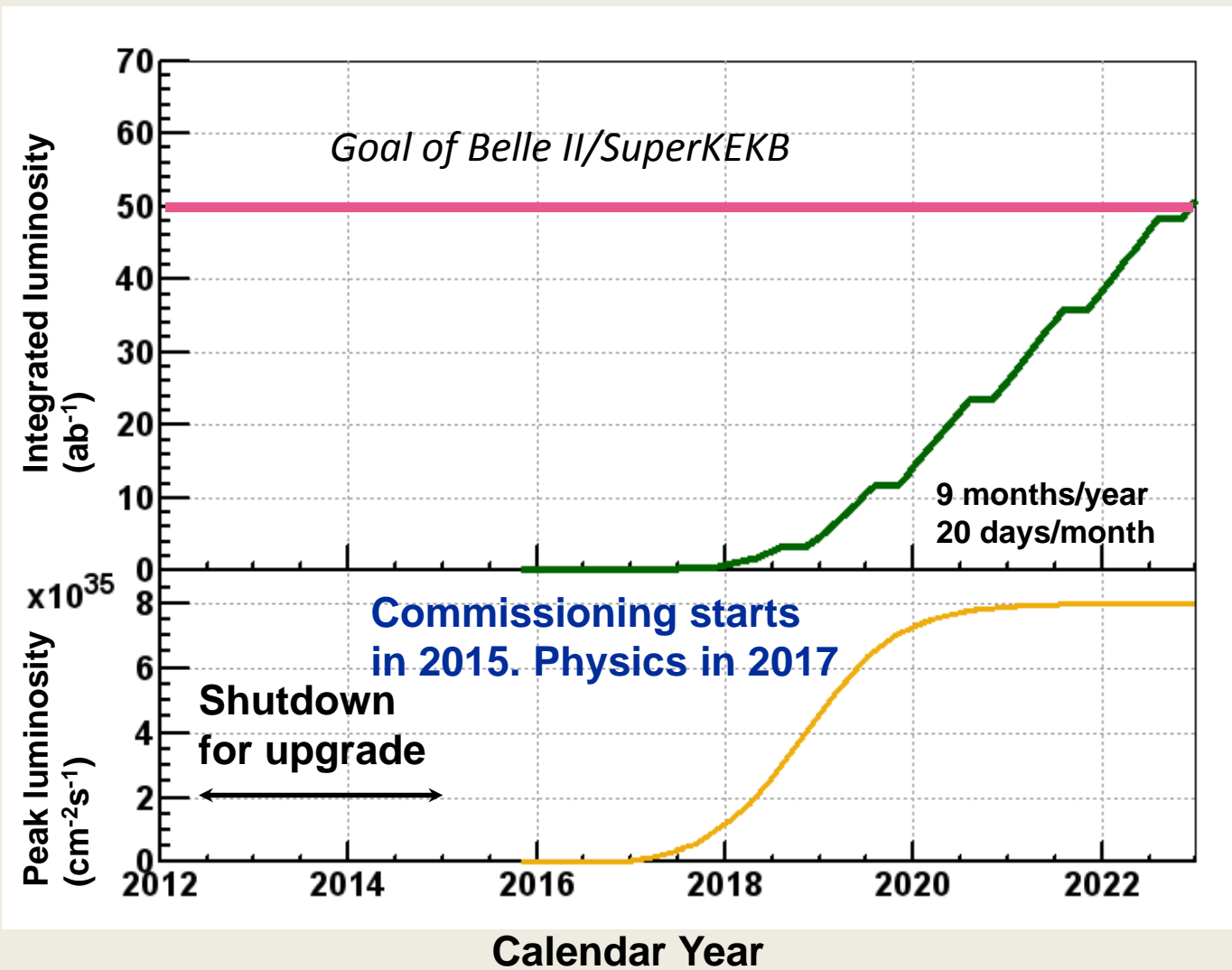
# Physics Reach of Belle II and the LHCb upgrade

Observable	Expected th. accuracy	Expected exp. uncertainty	Facility
CKM matrix			
$ V_{us}  [K \rightarrow \pi \ell \nu]$	**	0.1%	<i>K</i> -factory
$ V_{cb}  [B \rightarrow X_c \ell \nu]$	**	1%	Belle II
$ V_{ub}  [B_d \rightarrow \pi \ell \nu]$	*	4%	Belle II
$\sin(2\phi_1) [c\bar{c}K_S^0]$	***	$8 \cdot 10^{-3}$	Belle II/LHCb
$\phi_2$		$1.5^\circ$	Belle II
$\phi_3$	***	$3^\circ$	LHCb
CPV			
$S(B_s \rightarrow \psi\phi)$	**	0.01	LHCb
$S(B_s \rightarrow \phi\phi)$	**	0.05	LHCb
$S(B_d \rightarrow \phi K)$	***	0.05	Belle II/LHCb
$S(B_d \rightarrow \eta' K)$	***	0.02	Belle II
$S(B_d \rightarrow K^*(\rightarrow K_S^0 \pi^0) \gamma)$	***	0.03	Belle II
$S(B_s \rightarrow \phi \gamma)$	***	0.05	LHCb
$S(B_d \rightarrow \rho \gamma)$		0.15	Belle II
$A_{SL}^d$	***	0.001	LHCb
$A_{SL}^s$	***	0.001	LHCb
$A_{CP}(B_d \rightarrow s \gamma)$	*	0.005	Belle II
rare decays			
$\mathcal{B}(B \rightarrow \tau \nu)$	**	3%	Belle II
$\mathcal{B}(B \rightarrow D \tau \nu)$		3%	Belle II
$\mathcal{B}(B_d \rightarrow \mu \nu)$	**	6%	Belle II
$\mathcal{B}(B_s \rightarrow \mu \mu)$	***	10%	LHCb
zero of $A_{FB}(B \rightarrow K^* \mu \mu)$	**	0.05	LHCb
$\mathcal{B}(B \rightarrow K^{(*)} \nu \nu)$	***	30%	Belle II
$\mathcal{B}(B \rightarrow s \gamma)$		4%	Belle II
$\mathcal{B}(B_s \rightarrow \gamma \gamma)$		$0.25 \cdot 10^{-6}$	Belle II (with $5 \text{ ab}^{-1}$ )
$\mathcal{B}(K \rightarrow \pi \nu \nu)$	**	10%	<i>K</i> -factory
$\mathcal{B}(K \rightarrow e \pi \nu) / \mathcal{B}(K \rightarrow \mu \pi \nu)$	***	0.1%	<i>K</i> -factory
charm and $\tau$			
$\mathcal{B}(\tau \rightarrow \mu \gamma)$	***	$3 \cdot 10^{-9}$	Belle II
$ q/p _D$	***	0.03	Belle II
$\arg(q/p)_D$	***	$1.5^\circ$	Belle II



# SuperKEKB luminosity projection

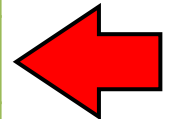
Belle/KEKB recorded  $\sim 1000 \text{ fb}^{-1}$ . Now change units to  $\text{ab}^{-1}$





# Compare the Parameters for KEKB and SuperKEKB

	KEKB Design	KEKB Achieved : with crab	SuperKEKB Nano-Beam
Energy (GeV) (LER/HER)	3.5/8.0	3.5/8.0	4.0/7.0
$\beta_y^*$ (mm)	10/10	5.9/5.9	0.27/0.30
$\beta_x^*$ (mm)	330/330	1200/1200	32/25
$\epsilon_x$ (nm)	18/18	18/24	3.2/5.3
$\epsilon_y / \epsilon_x$ (%)	1	0.85/0.64	0.27/0.24
$\sigma_y$ (mm)	1.9	0.94	0.048/0.062
$\sigma_y$	0.052	0.129/0.090	0.09/0.081
$\sigma_z$ (mm)	4	6 - 7	6/5
$I_{\text{beam}}$ (A)	2.6/1.1	1.64/1.19	3.6/2.6
$N_{\text{bunches}}$	5000	1584	2500
Luminosity ( $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ )	1	2.11	80



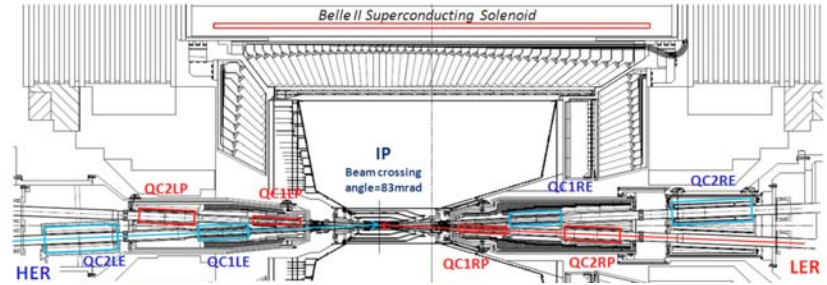
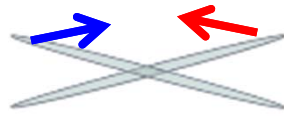
Nano-beams are the key (vertical spot size is  $\sim 50\text{nm}$  !!)

This is not a typo

K-I. Kanazawa@ICHEP2014



Colliding bunches



New superconducting final focusing magnets near the IP

$e^+$  3.6A

$e^-$  2.6A

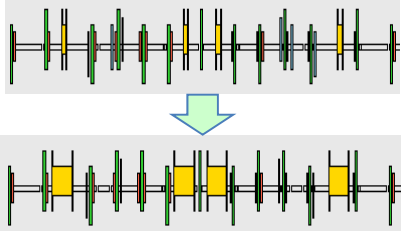
# KEKB to SuperKEKB

- ◆ Nano-Beam scheme  
extremely small  $\beta_y^*$   
low emittance
- ◆ Beam current X 2

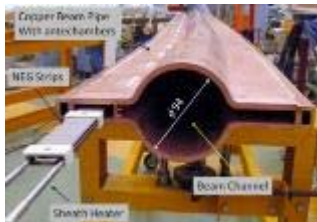
$$L = \frac{\gamma_{\pm}}{2e r_e} \left( 1 + \frac{\sigma_y^*}{\sigma_x^*} \left( \frac{I_{\pm} \xi_{\pm y}}{\beta_y^*} \right) \left( \frac{R_L}{R_y} \right) \right)$$

40 times higher luminosity  
 $2.1 \times 10^{34} \rightarrow 8 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$

Redesign the lattice to reduce the emittance (replace short dipoles with longer ones, increase wiggler cycles) (*all magnets installed 8/2014*)



Replace beam pipes with TiN-coated beam pipes with antechambers (*85% installed*)



Reinforce RF systems for higher beam currents

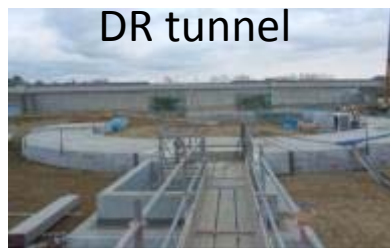
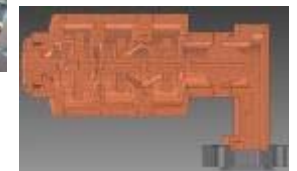


Improve monitors and control system

Injector Linac upgrade

Upgrade positron capture section

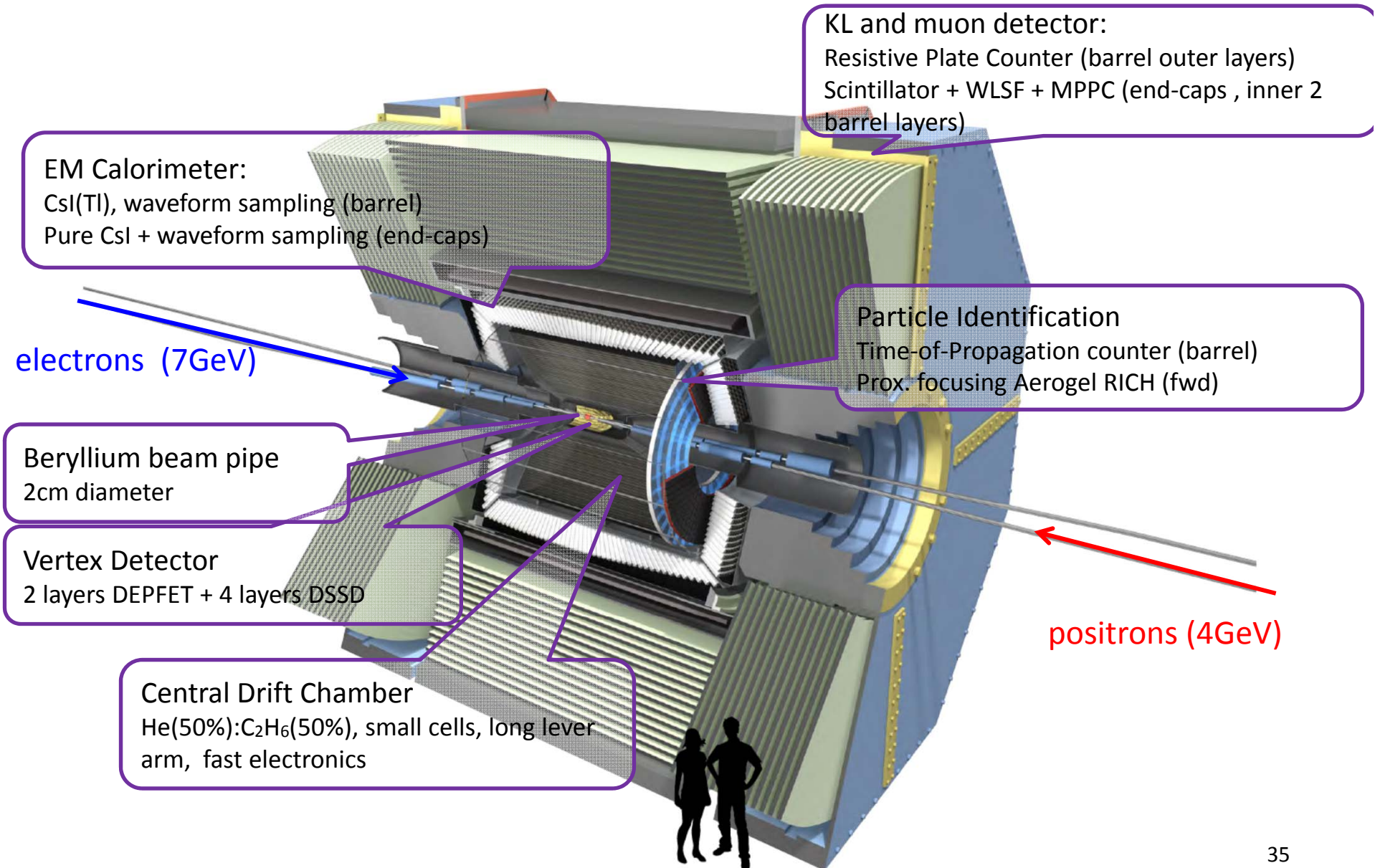
Low emittance RF electron gun



DR tunnel

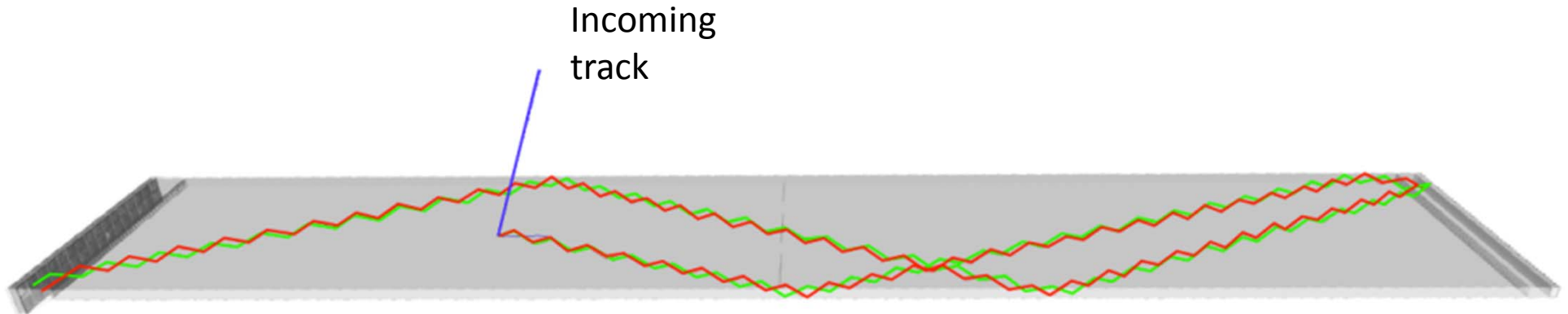
New  $e^+$  Damping Ring constructed

# Belle II Detector



# Barrel PID

A GEANT4 event display of a 2 GeV pion and kaon interacting in a TOP quartz bar. (Japan, US, Slovenia, Italy)



## Vertexing/Inner Tracking

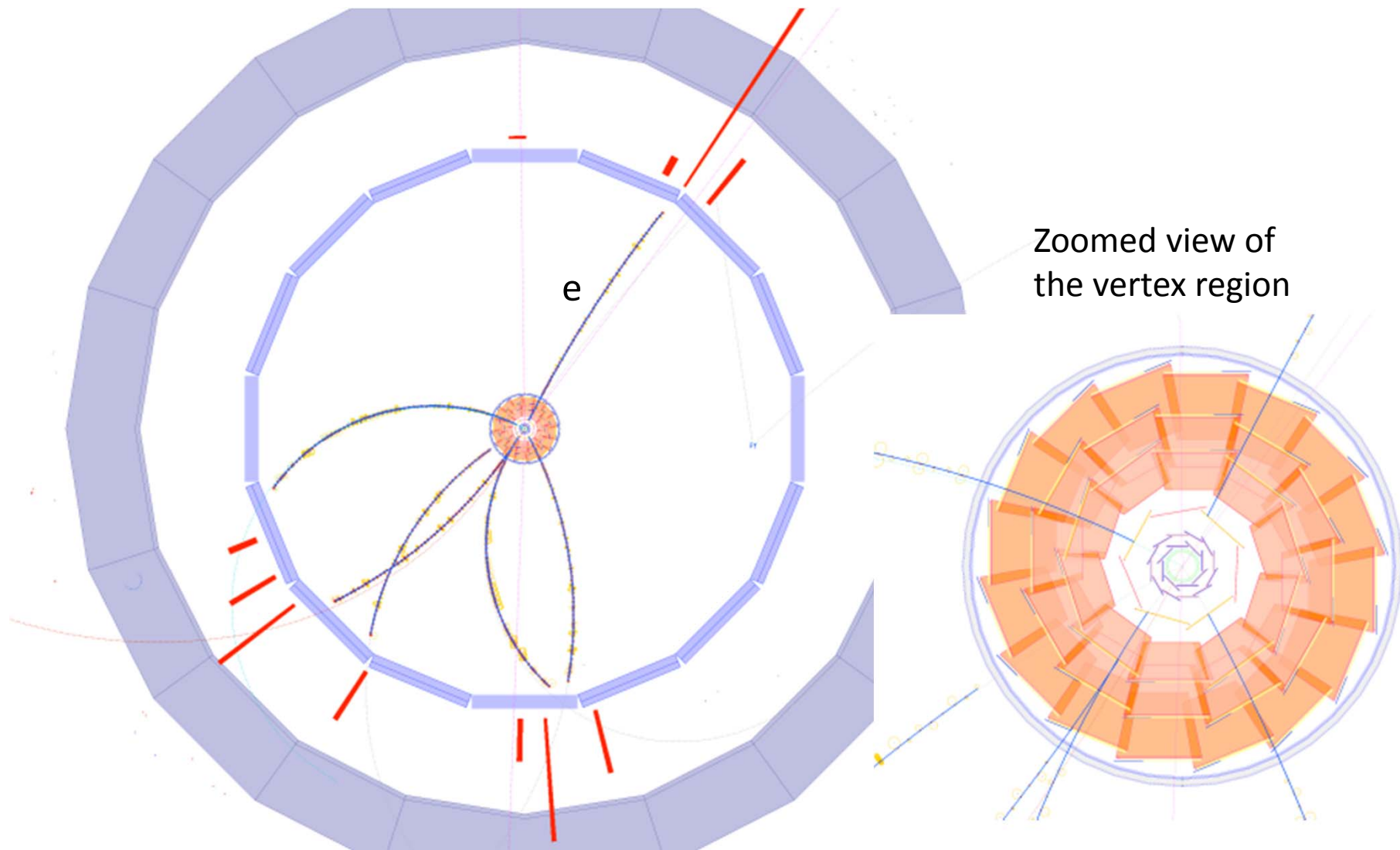


- Beampipe  $r=10$  mm
- DEPFET pixels (Germany, Czech Republic...)
  - Layer 1  $r=14$  mm
  - Layer 2  $r=22$  mm
- DSSD (double sided silicon detectors) FWD/BWD
  - Layer 3  $r=38$  mm (Australia) Italy
  - Layer 4  $r=80$  mm (India)
  - Layer 5  $r=115$  mm (Austria)
  - Layer 6  $r=140$  mm (Japan)

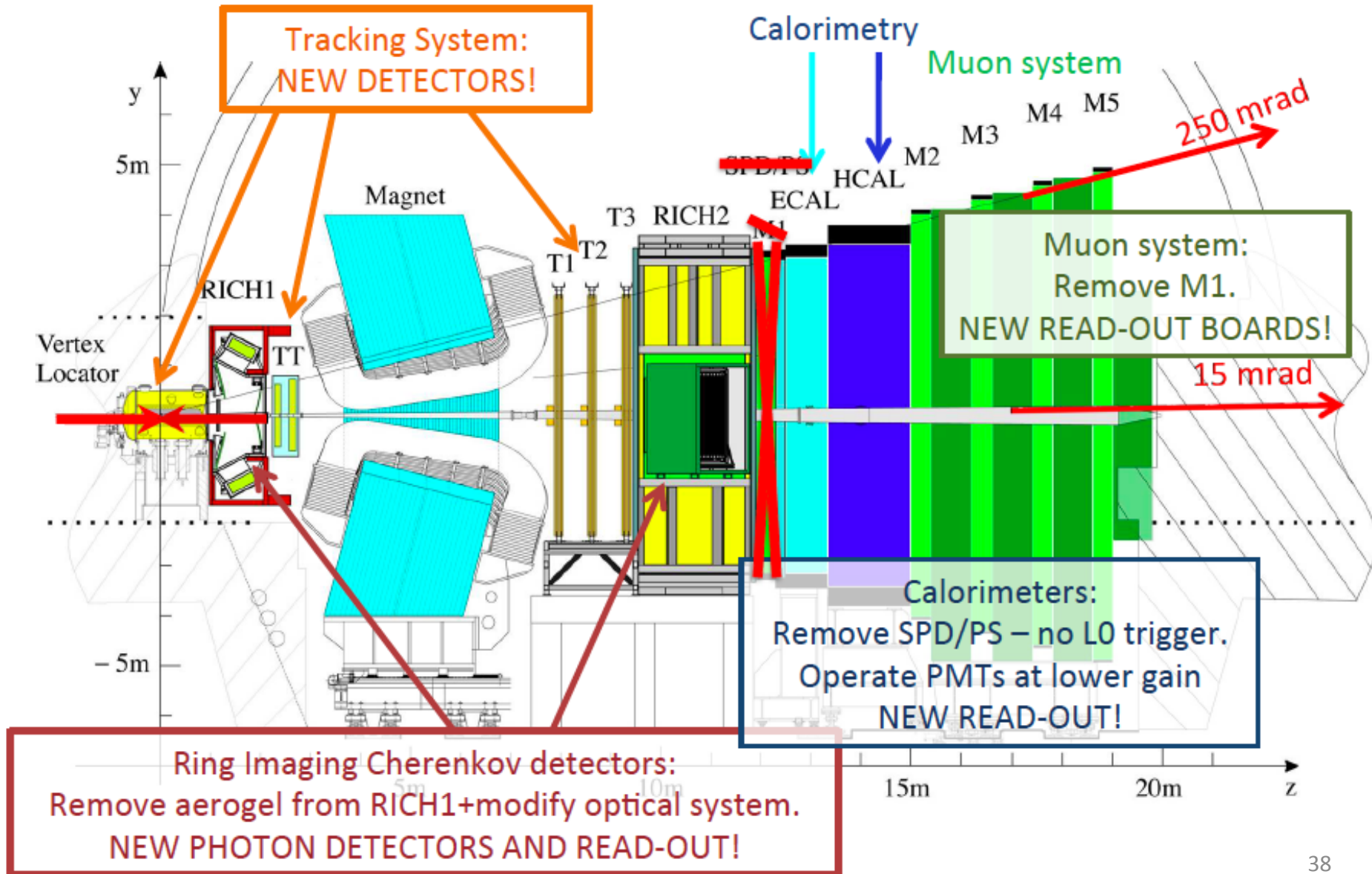
+Poland, Korea

# “Missing Energy Decay” in a Belle II GEANT4 MC

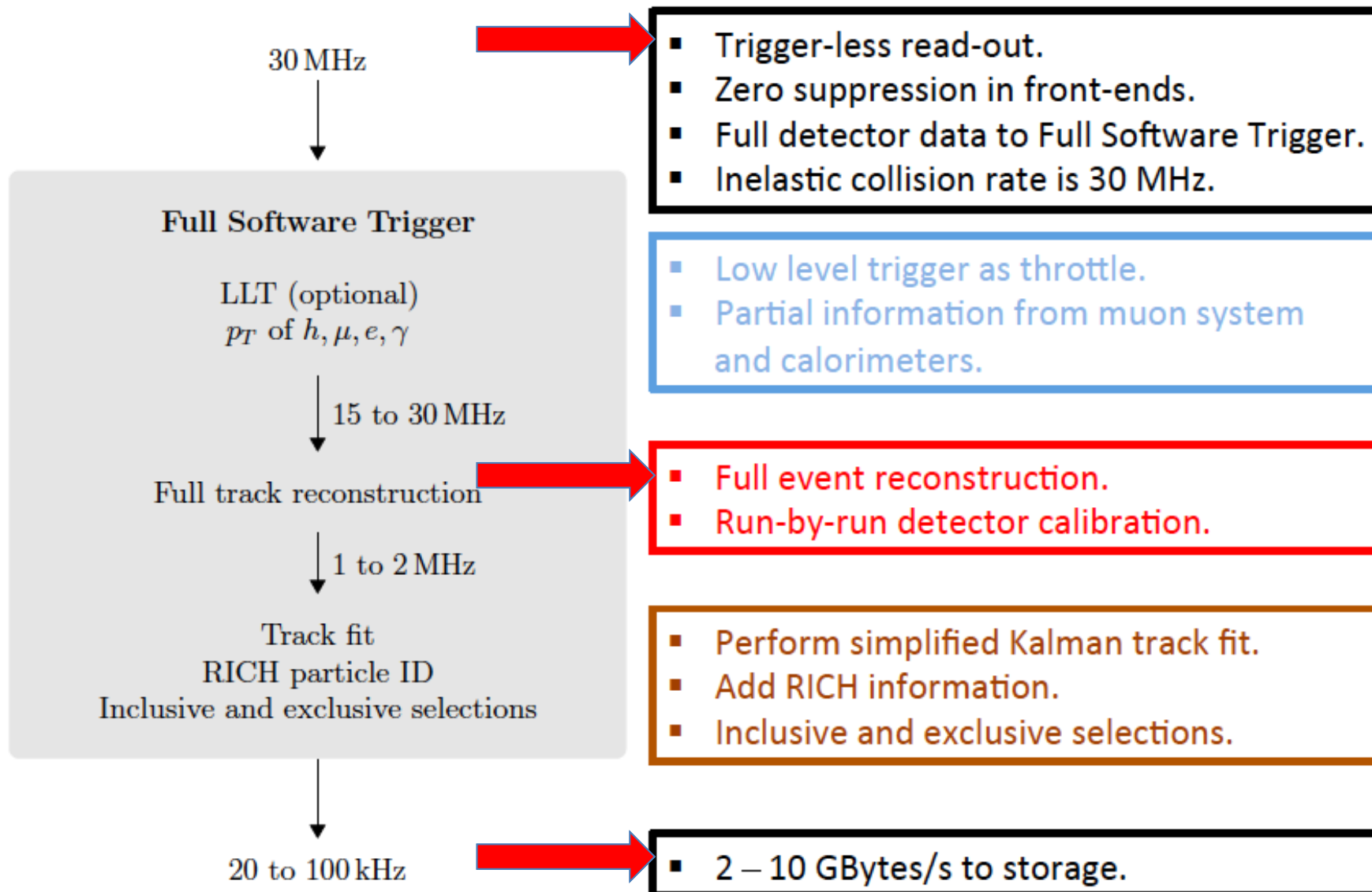
$B \rightarrow \tau \nu$ ,  $\tau \rightarrow e \nu \nu$      $B \rightarrow D \pi$ ,  $D \rightarrow K \pi \pi \pi$



# Upgraded LHCb detector

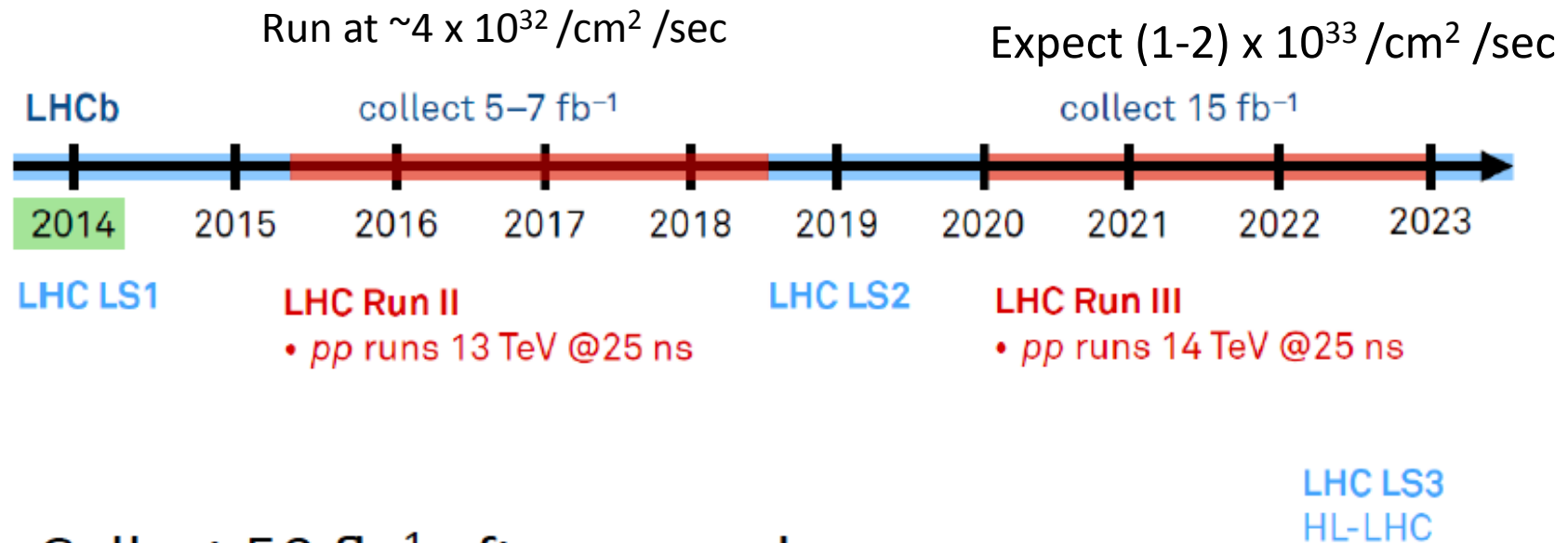


# Upgrade Trigger



# LHCb upgrade schedule

## Schedule / timeline



- Collect  $50 \text{ fb}^{-1}$  after upgrade.
- Continue taking data during HL-LHC.



# Conclusion/Next Generation.

- 2014 is the 50<sup>th</sup> anniversary of the discovery of CP violation in the kaon sector.
- The e<sup>+</sup> e<sup>-</sup> B factories confirmed that the KM phase is responsible for most of the observed CPV.
- LHCb has ruled out large CPV phases from NP in the B<sub>s</sub> sector.
- Nevertheless, 10-20% NP effects are consistent with all current data.
- “Missing energy decays” provide important constraints on the charged Higgs.

*Belle II roll-in in early 2017 with first physics runs and the LHCb upgrade in ~2020. These facilities will inaugurate a new era of flavor physics and the study of CP violation.*

# Backup slides



# New Reference *for the Next Generation*

## **The Physics of the B Factories**

<http://arxiv.org/abs/1406.6311>

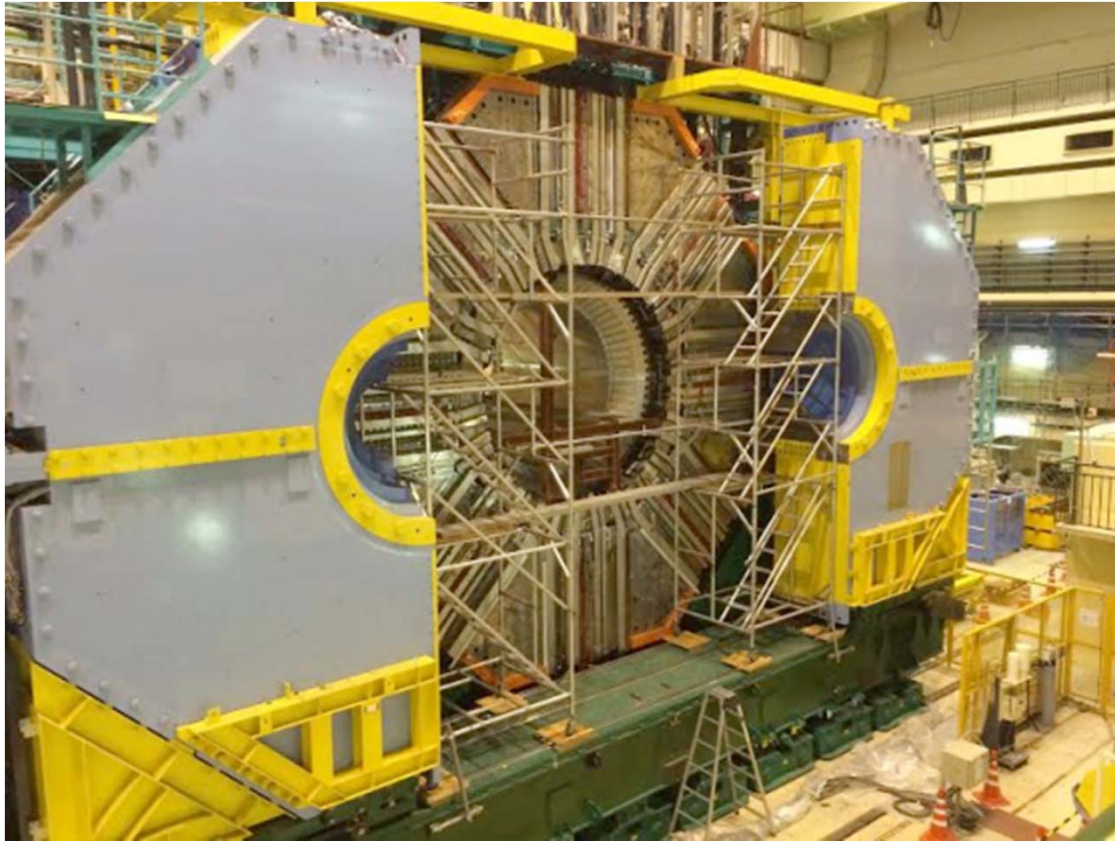
This work is on the Physics of the B Factories. Part A of this book contains a brief description of the SLAC and KEK B Factories as well as their detectors, BaBar and Belle, and data taking related issues. Part B discusses tools and methods used by the experiments in order to obtain results. The results themselves can be found in Part C.

Comments: 928 pages

Subjects: High Energy Physics - Experiment (hep-ex); High Energy Physics - Phenomenology (hep-ph)

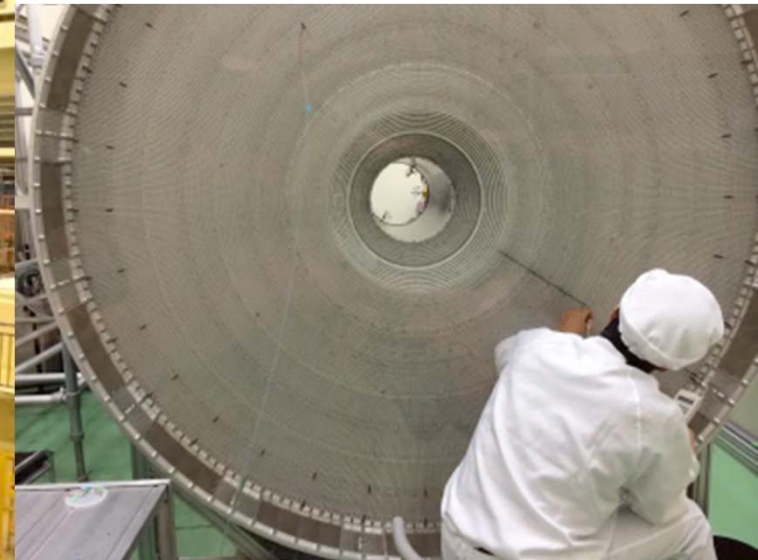
Report number: SLAC-PUB-15968, KEK Preprint 2014-3

## The scene at KEK in June 2014



*Tsukuba Hall*

Belle II detector with barrel KLM upgrade and forward muon endcap upgrade now installed. Next is the backward KLM upgrade.



*Fuji Hall*

Central Drift Chamber

More backup

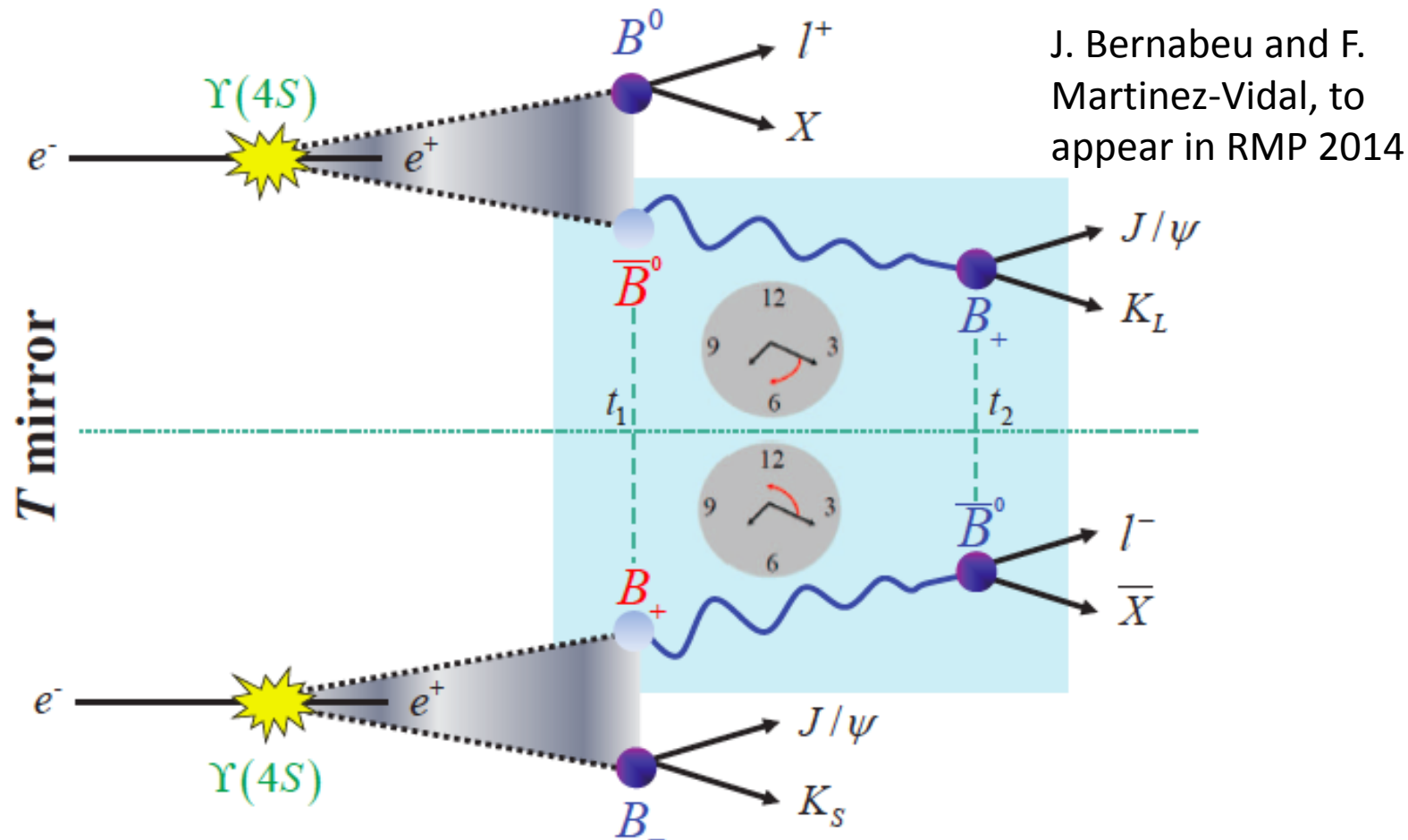
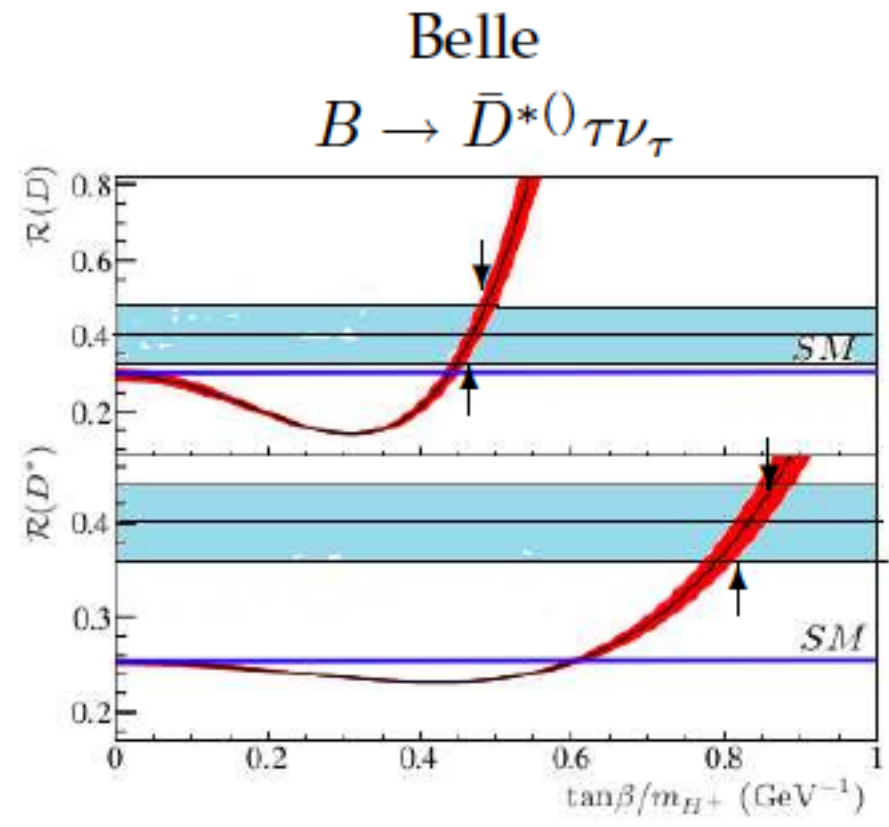
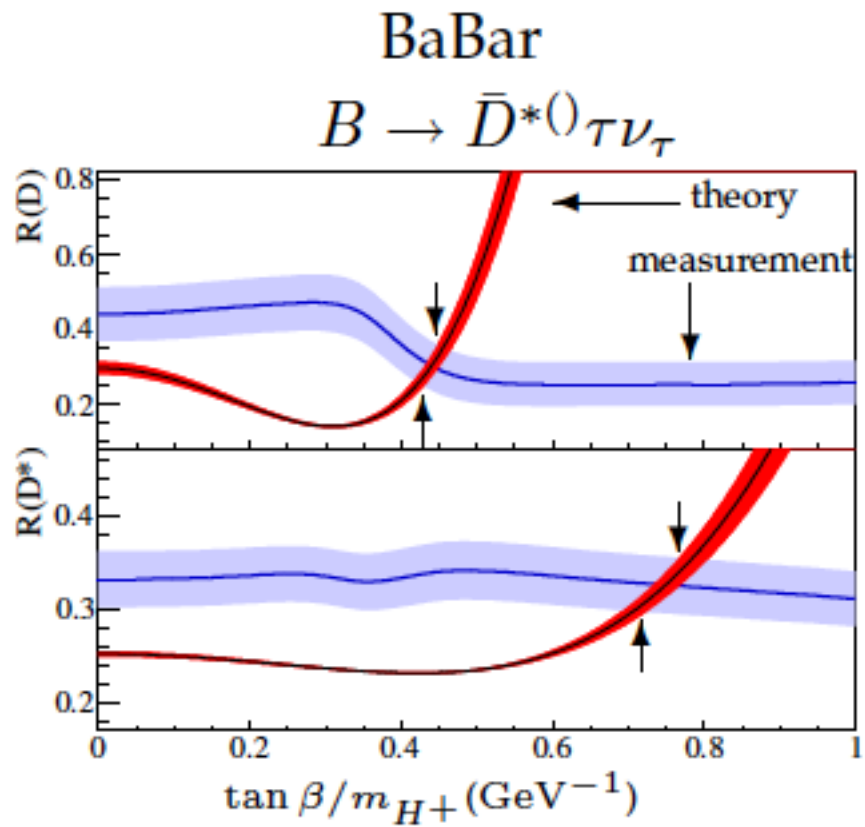


FIG. 11. Foundations of the time-reversal experiment. Electron-positron collisions at the asymmetric  $B$  factory produce  $\Upsilon(4S)$  resonances, each of which decays through strong interaction in an entangled pair of  $B$  mesons. When one  $B$  meson decays at  $t_1$ , the identity of the other is “tagged” without measuring it specifically. In the top panel, the  $B$  meson observed to decay to the final state  $l^+ X$  at  $t_1$  transfers information to the (still living) partner meson and dictates that it is in a  $\bar{B}^0$  state. This surviving meson tagged as  $\bar{B}^0$  is observed later at  $t_2$ , encapsulating a time ordering, to decay into a final state  $J/\psi K_L^0$  that filters the  $B$  meson to be in a  $B_+$  state, a linear combination of  $B^0$  and  $\bar{B}^0$  states. This case corresponds to a transition  $\bar{B}^0 \rightarrow B_+$ . To study time reversal we have to compare the rate at which this transition occurs to the rate of the time-reversed transition,  $B_+ \rightarrow \bar{B}^0$  (bottom panel). Adapted from<sup>2</sup>.

# Comparison of BaBar and Belle (not updated)



# New CMS Bs mixing result





# Signal model

We use the same notations as LHCb [arXiv:1304.2600]:

$$\frac{d^4\Gamma(B_s(t))}{d\Theta dt} = X(\Theta, \alpha, t) = \sum_{i=1}^{10} O_i(\alpha, t) \cdot g_i(\Theta),$$

$$O_i(\alpha, t) = N_i e^{-\Gamma_s t} \left[ a_i \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_i \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) + c_i \cos(\Delta m_s t) + d_i \sin(\Delta m_s t) \right]$$

$i$	$g_i(\theta_T, \psi_T, \phi_T)$	$N_i$	$a_i$	$b_i$	$c_i$	$d_i$
1	$2 \cos^2 \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_0(0) ^2$	1	$D$	$C$	$-S$
2	$\sin^2 \psi_T (1 - \sin^2 \theta_T \sin^2 \phi_T)$	$ A_{\parallel}(0) ^2$	1	$D$	$C$	$-S$
3	$\sin^2 \psi_T \sin^2 \theta_T$	$ A_{\perp}(0) ^2$	1	$-D$	$C$	$S$
4	$-\sin^2 \psi_T \sin 2\theta_T \sin \phi_T$	$ A_{\parallel}(0)A_{\perp}(0) $	$C \sin(\delta_{\perp} - \delta_{\parallel})$	$S \cos(\delta_{\perp} - \delta_{\parallel})$	$\sin(\delta_{\perp} - \delta_{\parallel})$	$D \cos(\delta_{\perp} - \delta_{\parallel})$
5	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin^2 \theta_T \sin 2\phi_T$	$ A_0(0)A_{\parallel}(0) $	$\cos(\delta_{\parallel} - \delta_0)$	$D \cos(\delta_{\parallel} - \delta_0)$	$C \cos(\delta_{\parallel} - \delta_0)$	$-S \cos(\delta_{\parallel} - \delta_0)$
6	$\frac{1}{\sqrt{2}} \sin 2\psi_T \sin 2\theta_T \sin \phi_T$	$ A_0(0)A_{\perp}(0) $	$C \sin(\delta_{\perp} - \delta_0)$	$S \cos(\delta_{\perp} - \delta_0)$	$\sin(\delta_{\perp} - \delta_0)$	$D \cos(\delta_{\perp} - \delta_0)$
7	$\frac{2}{3}(1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_S(0) ^2$	1	$-D$	$C$	$S$
8	$\frac{1}{3}\sqrt{6} \sin \psi_T \sin^2 \theta_T \sin 2\phi_T$	$ A_S(0)A_{\parallel}(0) $	$C \cos(\delta_{\parallel} - \delta_S)$	$S \sin(\delta_{\parallel} - \delta_S)$	$\cos(\delta_{\parallel} - \delta_S)$	$D \sin(\delta_{\parallel} - \delta_S)$
9	$\frac{1}{3}\sqrt{6} \sin \psi_T \sin 2\theta_T \cos \phi_T$	$ A_S(0)A_{\perp}(0) $	$\sin(\delta_{\perp} - \delta_S)$	$-D \sin(\delta_{\perp} - \delta_S)$	$C \sin(\delta_{\perp} - \delta_S)$	$S \sin(\delta_{\perp} - \delta_S)$
10	$\frac{4}{3}\sqrt{3} \cos \psi_T (1 - \sin^2 \theta_T \cos^2 \phi_T)$	$ A_S(0)A_0(0) $	$C \cos(\delta_0 - \delta_S)$	$S \sin(\delta_0 - \delta_S)$	$\cos(\delta_0 - \delta_S)$	$D \sin(\delta_0 - \delta_S)$

$$C = \frac{1 - |\lambda|^2}{1 + |\lambda|^2},$$

$$S = -\frac{2|\lambda| \sin \phi_s}{1 + |\lambda|^2},$$

$$D = -\frac{2|\lambda| \cos \phi_s}{1 + |\lambda|^2}$$



$|\lambda|$  includes possible contribution from CP violation in direct decay, we assume  $|\lambda| = 1$  and we assign a systematic.

$\Delta\Gamma_s > 0$ : we use previous LHCb results.  $\alpha$  physics parameters ( $\Delta\Gamma_s, \phi_s, c_T, |A_0|^2, |A_S|^2, |A_{\perp}|^2, \delta_{\parallel\pm}, \delta_{S\perp}, \delta_{\perp}$ )



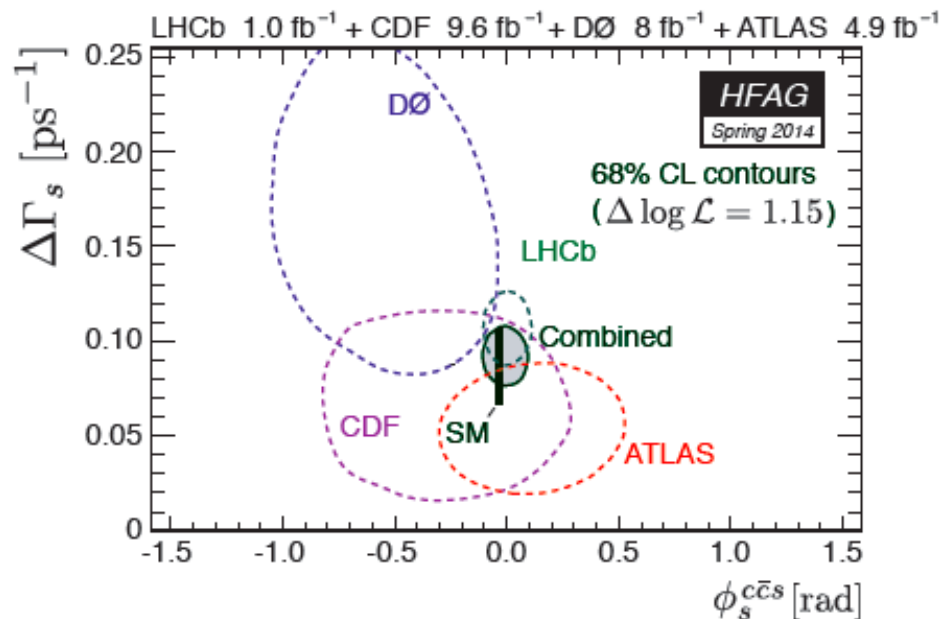
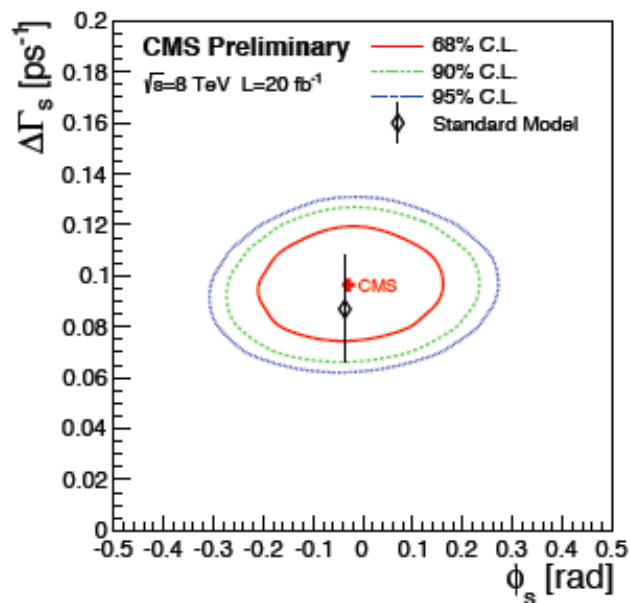


## Results summary

- Analysing the 2012 CMS data ( $20.0 \text{ fb}^{-1}$ ), we selected 49k  $B_s$  signal events. We obtain:

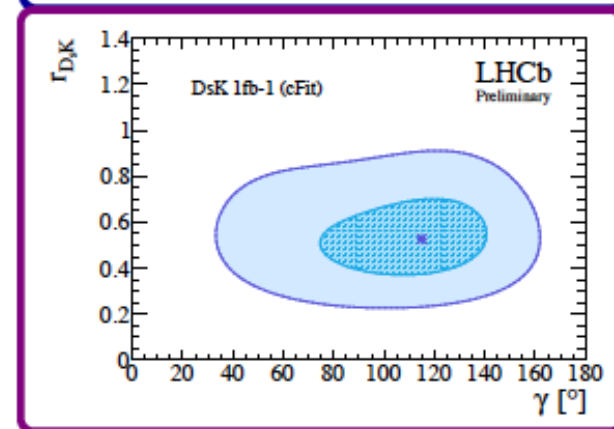
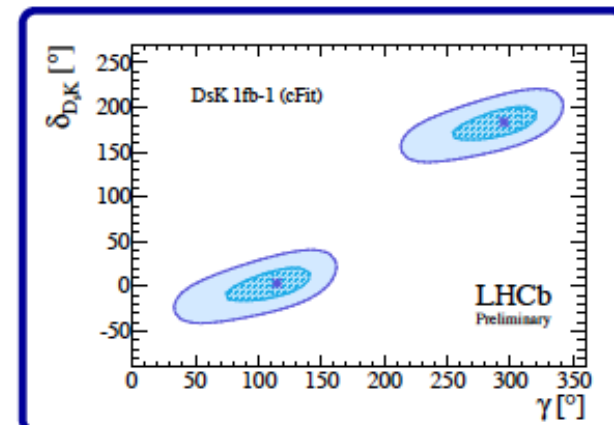
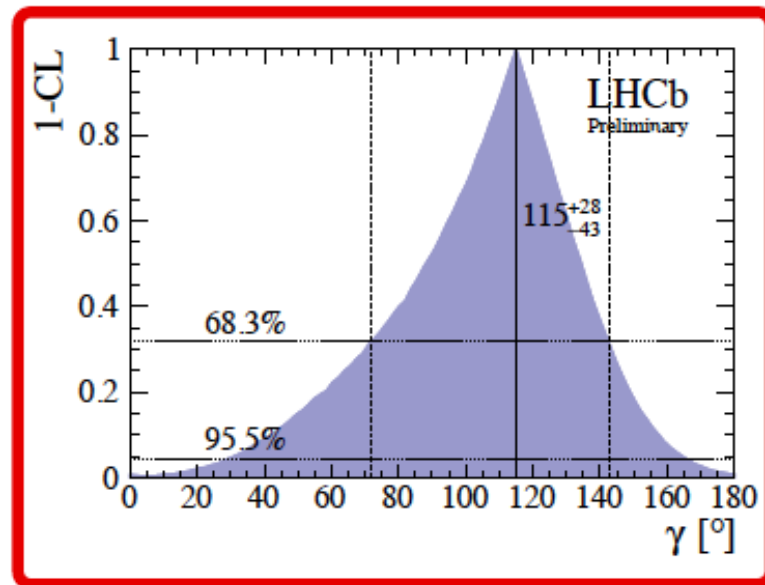
$$\phi_s = -0.03 \pm 0.11 \text{ (stat.)} \pm 0.03 \text{ (syst.) rad}$$
$$\Delta\Gamma_s = 0.096 \pm 0.014 \text{ (stat.)} \pm 0.007 \text{ (syst.) ps}^{-1}$$

- Contour plot (stat. only), constraining  $\Delta\Gamma_s > 0$ :



# New LHCb $\gamma$ results

- $\gamma$  angle extraction based on cFit results.
- Statistical and systematic uncertainties and correlations are taken into account.
- Constraint:  $|C|^2 + |S_f|^2 + |A_f^{\Delta\Gamma}|^2 = 1$



$$\begin{aligned} \gamma &= (115^{+28}_{-43})^\circ \\ \delta_{D_s K} &= (3^{+19}_{-20})^\circ \\ r_{D_s K} &= 0.53^{+0.17}_{-0.16} \end{aligned}$$

➤ Many modes discovered and studied mainly based on  $1\text{fb}^{-1}$  and to be updated with full datasets ( $3\text{fb}^{-1}$ )

➤ gamma combination with part of current measurements

1  $\text{fb}^{-1}$  GLW/ADS on  $B^\pm \rightarrow Dh^\pm, D \rightarrow hh$

Phys. Lett. B 712 (2012) 203

1  $\text{fb}^{-1}$  ADS on  $B^\pm \rightarrow Dh^\pm, D \rightarrow K\pi\pi$

Phys. Lett. B 723 (2013) 44

1  $\text{fb}^{-1}$ + 2  $\text{fb}^{-1}$  GGSZ (MI) on  $B^\pm \rightarrow DK^\pm, D \rightarrow Kshh$

Phys. Lett. B 718 (2012) 43  
LHCb-CONF-2013-004

CLEO inputs on D system nuisance parameters

CLEO Collaboration, Phys. Rev. D80 (2009) 031105

(Interesting proposals to use D mixing to further constrain, i.e.  $K3\pi$ )

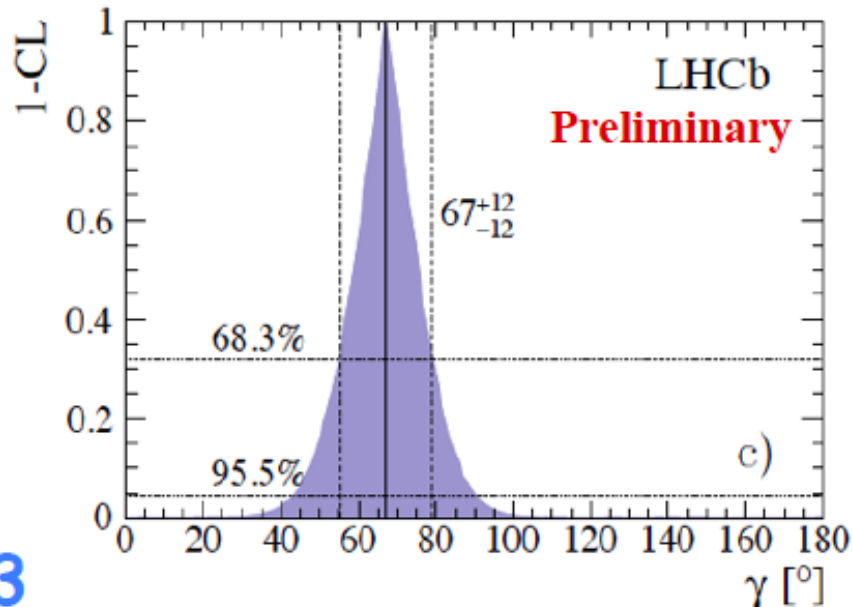
S. Harnew, J. Rademacker  
Phys. Lett. B 728 (2014) 296

Additional D mixing constrain from LHCb

Phys. Rev. Lett. 110 (2012) 101802

Direct CP violation in  $D \rightarrow KK$  and  $D \rightarrow \pi\pi$

HFAG



$$\gamma = (67 \pm 12)^\circ \quad \gamma \in [55.1, 79.1]^\circ \text{ at 68\% CL}$$

$$\gamma \in [43.9, 89.5]^\circ \text{ at 95\% CL}$$

$$\delta_B^K = (114.3^{+12}_{-13})^\circ$$

$$r_B^K = 0.0923^{+0.0078}_{-0.0080}$$

$$\delta_B^K \in [101.3, 126.3]^\circ \text{ at 68\% CL}$$

$$r_B^K \in [0.0843, 0.1001] \text{ at 68\% CL}$$

$$\delta_B^K \in [88.7, 136.3]^\circ \text{ at 95\% CL}$$

$$r_B^K \in [0.0762, 0.1075] \text{ at 95\% CL}$$

LHCb-CONF-2013-006

Combining with B-factories:

CKMFitter

UTfit

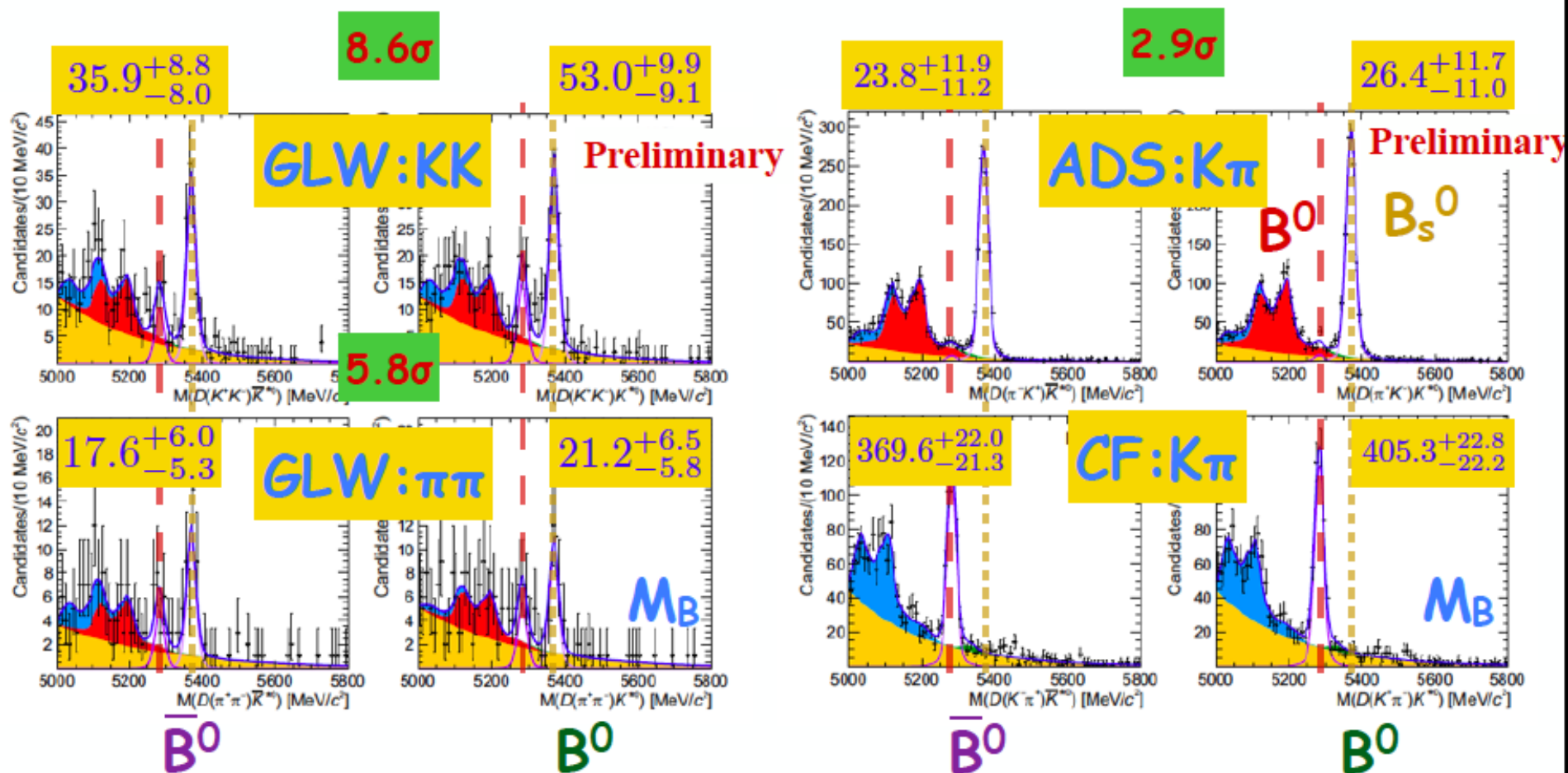
$$\gamma = (70.0^{+7.7}_{-9.0})^\circ$$

$$\gamma = (68.3 \pm 7.5)^\circ$$

# GLW/ADS Analysis of $B^0 \rightarrow DK^{*0}$

LHCb-PAPER-2014-028; 3 fb<sup>-1</sup>

➤ Similar as  $B^+ \rightarrow DK^+$ ,  $\gamma$  could also be measured with self-tagged neutral B decays; Diluted due to resonance shape of  $K^*$



# Results and Interpretation of $B^0 \rightarrow DK^{*0}$

LHCb-PAPER-2014-028; 3 fb<sup>-1</sup>

➤ 12 ratios made according to yields from  $B^0$  and  $B_s^0$

<b>GLW</b>	$\frac{N(\bar{B}^0) - N(B^0)}{N(\bar{B}^0) + N(B^0)}$	$B^0$	$\mathcal{A}_d^{KK} = -0.198^{+0.144}_{-0.145} +0.019_{-0.020}$	$\mathcal{A}_d^{\pi\pi} = -0.092^{+0.217}_{-0.217} +0.019_{-0.019}$
	$2 \times \frac{\Gamma(\bar{B}^0 \rightarrow D_{CP+} \bar{K}^{*0}) + \Gamma(B^0 \rightarrow D_{CP+} K^{*0})}{\Gamma(\bar{B}^0 \rightarrow D^0 \bar{K}^{*0}) + \Gamma(B^0 \rightarrow \bar{D}^0 K^{*0})}$	$B^0$	$\mathcal{R}_d^{KK} = 1.054^{+0.165}_{-0.153} +0.044_{-0.044}$	$\mathcal{R}_d^{\pi\pi} = 1.214^{+0.283}_{-0.252} +0.053_{-0.053}$
		$B_s^0$	$\mathcal{A}_s^{KK} = -0.044^{+0.073}_{-0.073} +0.019_{-0.020}$	$\mathcal{A}_s^{\pi\pi} = 0.064^{+0.130}_{-0.131} +0.018_{-0.019}$
		$B^0/B_s^0$	$\mathcal{R}_{ds}^{KK} = 0.103^{+0.018}_{-0.016} +0.009_{-0.009}$	$\mathcal{R}_{ds}^{\pi\pi} = 0.147^{+0.040}_{-0.036} +0.012_{-0.012}$

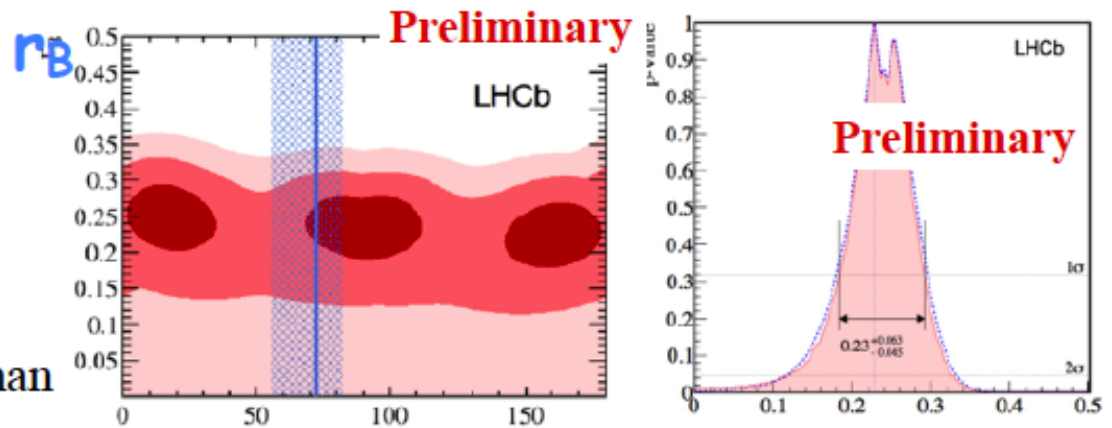
<b>ADS</b>	$\frac{\Gamma(\text{ADS})}{\Gamma(\text{CF})}$	$B^0$	$\mathcal{R}_d^+ = 0.057^{+0.029}_{-0.027} +0.009_{-0.012}$	$\mathcal{R}_d^- = 0.056^{+0.032}_{-0.030} +0.009_{-0.012}$
	$\frac{N(\bar{B}_{(s)}^0) - N(B_{(s)}^0)}{N(\bar{B}_{(s)}^0) + N(B_{(s)}^0)}$	$B^0, B_s^0$	$\mathcal{A}_d^{K\pi} = -0.032^{+0.041}_{-0.041} +0.019_{-0.020}$	$\mathcal{A}_s^{K\pi} = -0.014^{+0.025}_{-0.025} +0.019_{-0.019}$

➤ Contribution to  $\gamma$

➤ Dilution factor ( $0.95 \pm 0.03$ ) from toy simulation which models  $B^0 \rightarrow DK\pi$  resonances

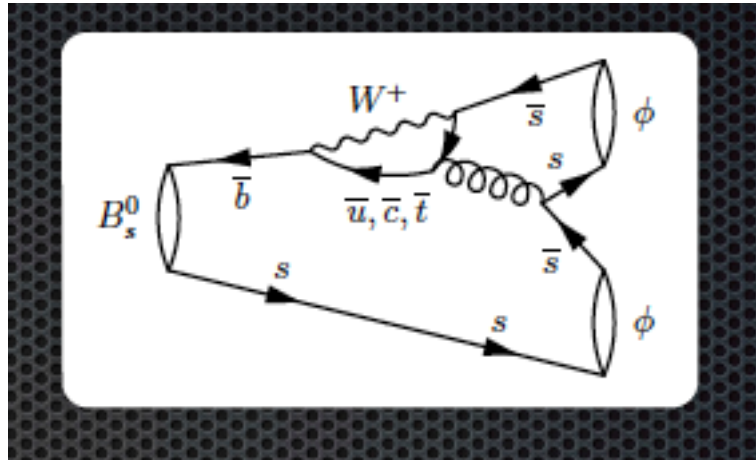
➤ Constrain on  $r_B$  in  $B^0 \rightarrow DK^{*0}$

➤ Compatible and more accurate than previous B-factory measurements

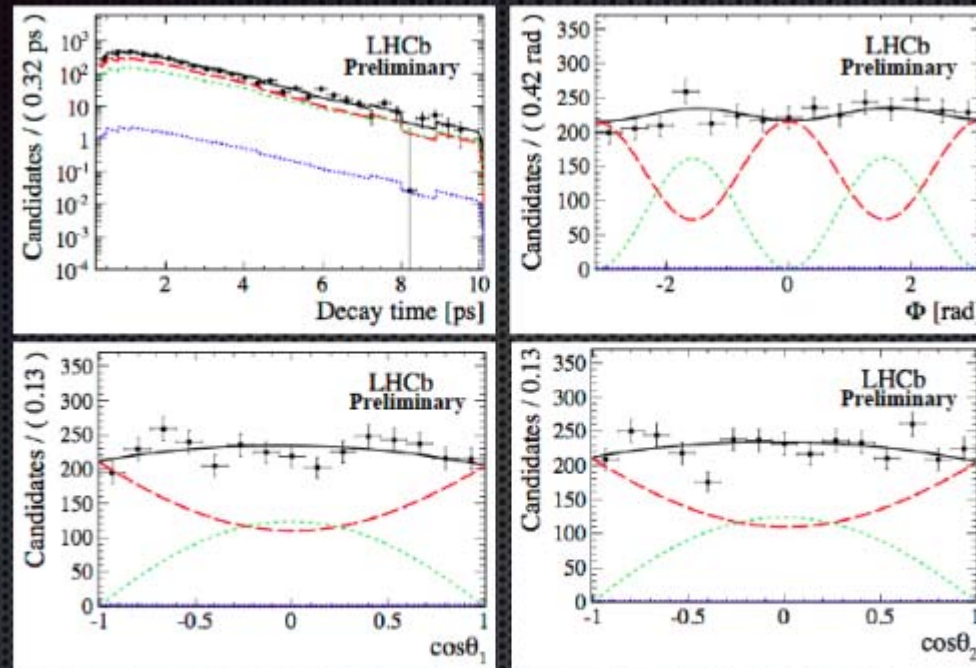


But LHCb dominates on these  $B_s$  modes

M.Needham@ICHEP2014



## $B_s \rightarrow \phi\phi$ - Time-Dependent Results



Projections are s-weighted and include acceptances,  
 Decay time acceptance from  $B_s \rightarrow D_s \pi$  data,  
 Angular acceptance from simulated events.



# Belle results on $\gamma$ at ICHEP2014

# ADS Measurement for $B^0 \rightarrow [K\pi]_D K^{*0}$

$$\mathcal{R}_{DK^{*0}} = r_S^2 + r_D^2 + 2kr_S r_D \cos(\delta_S + \delta_D) \cos\phi_3, \quad r_S^2 \equiv \frac{\Gamma(B^0 \rightarrow D^0 K^+ \pi^-)}{\Gamma(B^0 \rightarrow \bar{D}^0 K^+ \pi^-)} = \frac{\int dp A_{b \rightarrow u}^2(p)}{\int dp A_{b \rightarrow c}^2(p)}$$



PRD80 (2009) 031102

$R_{DK^*} = 0.067^{+0.070}_{-0.054} \pm 0.018$ ,  $R_{DK^*} < 0.244 @ 95\%CL$  for  $K\pi$  mode

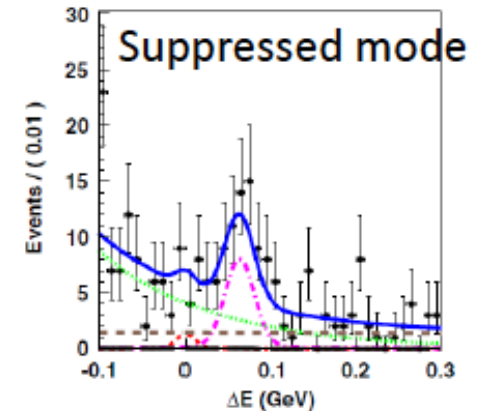
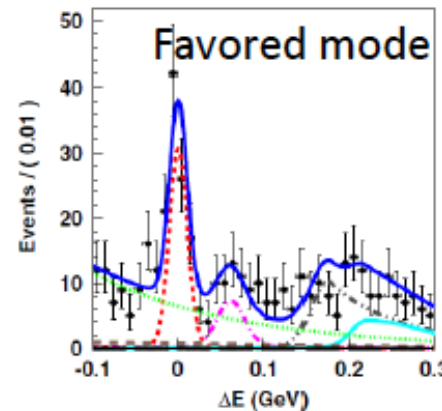
$r_S [0.07, 0.41]$  (95%CL)  $r_S = 0.26$  (most probable value) w/ combined  $K\pi, K\pi\pi^0, K\pi\pi\pi$



PRD86(2012)011101 772M

$R_{DK^*} = 0.045^{+0.056}_{-0.050} \pm 0.028$

$R_{DK^*} < 0.16 @ 95\%CL$  for  $K\pi$  mode



Approximation  $R_{DK^*} \approx r_S^2$  since  $r_D = 0.06$  from J.Phys.G 33 1(2006)

Under this naïve assumption,  $r_S \sim 0.21$

Combining the lower limit of  $r_S$  by BABAR and additional independent measurement by Belle  
It is possible that  $r_S$  bigger than the  $r_B \sim 0.1$

Also refer previous  
Dr. Wenbin Qian's Talk

This motivate the Dalitz analysis for  $B^0 \rightarrow [K_S \pi \pi]_D K^{*0}$

# ADS Measurement for $B^- \rightarrow [K\pi\pi^0]_D K^-$

$$R_{ADS}^{K\pi\pi^0} = r_B^2 + r_D^2 + 2r_B r_D R_{K\pi\pi^0} \cos(\delta_B + \delta_D^{K\pi\pi^0}) \cos\phi_3$$

Inclusive ( $\geq 3$ -body decay) ADS  
PRD 68, 033003 (2003)

$$A_{ADS}^{K\pi\pi^0} = 2r_B r_D R_{K\pi\pi^0} \sin(\delta_B + \delta_D^{K\pi\pi^0}) \sin\phi_3 / R_{ADS}^{K\pi\pi^0}$$

- Coherence factor  $R_{K\pi\pi^0} e^{i\delta_D^{K\pi\pi^0}} = \frac{\int d\vec{m} A_{DCS}(\vec{m}) A_{CF}(\vec{m}) e^{i\delta(\vec{m})}}{\sqrt{\int d\vec{m} A_{DCS}^2(\vec{m}) \int d\vec{m} A_{CF}^2(\vec{m})}}$   $\vec{m} = [m_{K\pi}^2, m_{K\pi^0}^2], 0 \leq R_{K\pi\pi^0} \leq 1$

CLEO-c update ( $R_{K\pi\pi^0}, \delta_D^{K\pi\pi^0}$ ) measurement

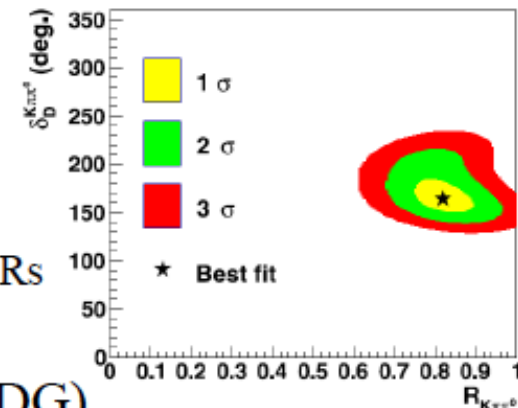
PRD 80, 031105(R)

$$R_{K\pi\pi^0} = 0.84 \pm 0.07, \delta_D^{K\pi\pi^0} = 227^{+14}_{-17}$$

PLB 731(2014)197-203

$$R_{K\pi\pi^0} = 0.82 \pm 0.07, \delta_D^{K\pi\pi^0} = 164^{+20}_{-14}$$

Updates of  
Mixing params  
 $\delta_D^{K\pi}, (r_D^{K\pi})^2, BRs$

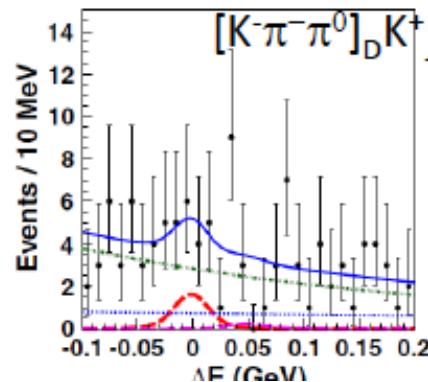
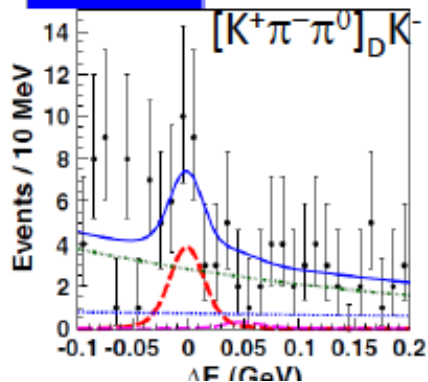


- $(r_B, \delta_B)$  same as  $[K\pi]_D K, r_D = (2.20 \pm 0.10) \times 10^{-3}$  (PDG)



PRD88, 091104(R) (2013) 772M

$$A_{ADS}^{K\pi\pi^0} = 0.41 \pm 0.30(\text{stat}) \pm 0.05(\text{syst})$$

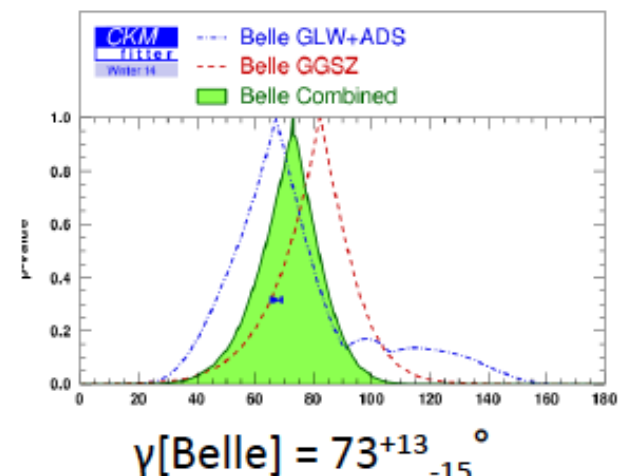
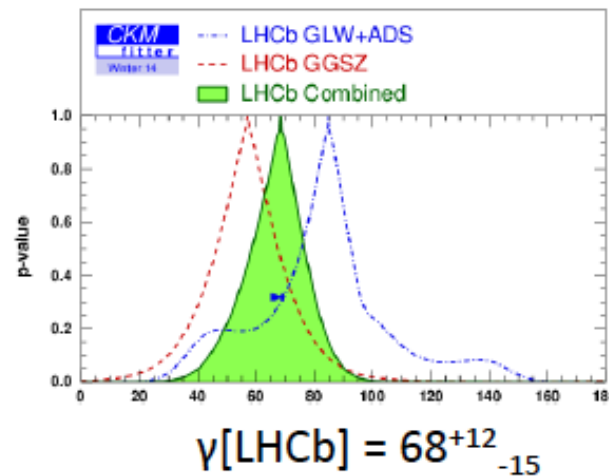
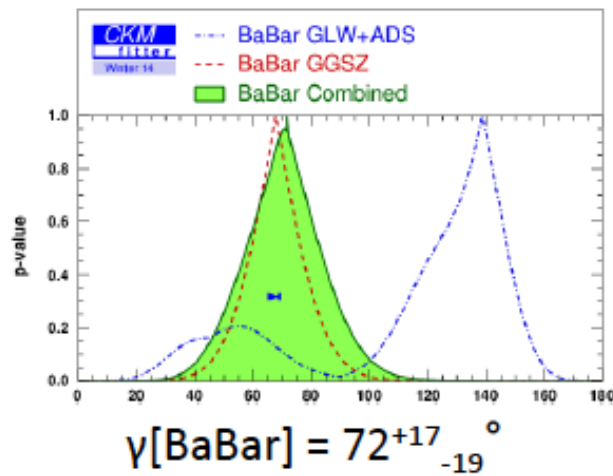


$$R_{ADS}^{K\pi\pi^0} = [1.98 \pm 0.62(\text{stat}) \pm 0.24(\text{syst})] \times 10^{-2}$$

First evidence of  $[K\pi\pi^0]_D K$  at  $3.2\sigma$   
 $B \rightarrow [K\pi\pi^0]_D K$  ADS result should use  
the CLEO-c updated parameters for  $\gamma/\phi_3$

# Summary

- Fruitful outputs: First evidence ADS  $[K\pi]_D K$  and  $[K\pi\pi^0]_D K$ , Unique GLW CP- mode, First model independent Dalitz analysis...
- $B^0 \rightarrow [K\pi]_D K^{*0}$  ADS results may indicate bigger  $r_s$  value than the  $r_B$ . It is possible higher sensitivity of  $\gamma/\phi_3$  than the  $B^- \rightarrow [K_S \pi\pi]_D K^{*-}$
- $B \rightarrow [K\pi\pi^0]_D K$  ADS result should use the CLEO-c updated parameters for  $\gamma/\phi_3$  determination.
- There is still a lot of modes playing an important role for  $\gamma/\phi_3$  determination from Belle



- $B^0 \rightarrow [K_S \pi\pi]_D K^{*0}$  and  $B^- \rightarrow [K_S K\pi]_D K^-$  Dalitz etc. analyses on going!

New BaBar result on CPV in  $B \rightarrow K^* \pi$

# Combined fit to $B^+$ and $B^-$ samples

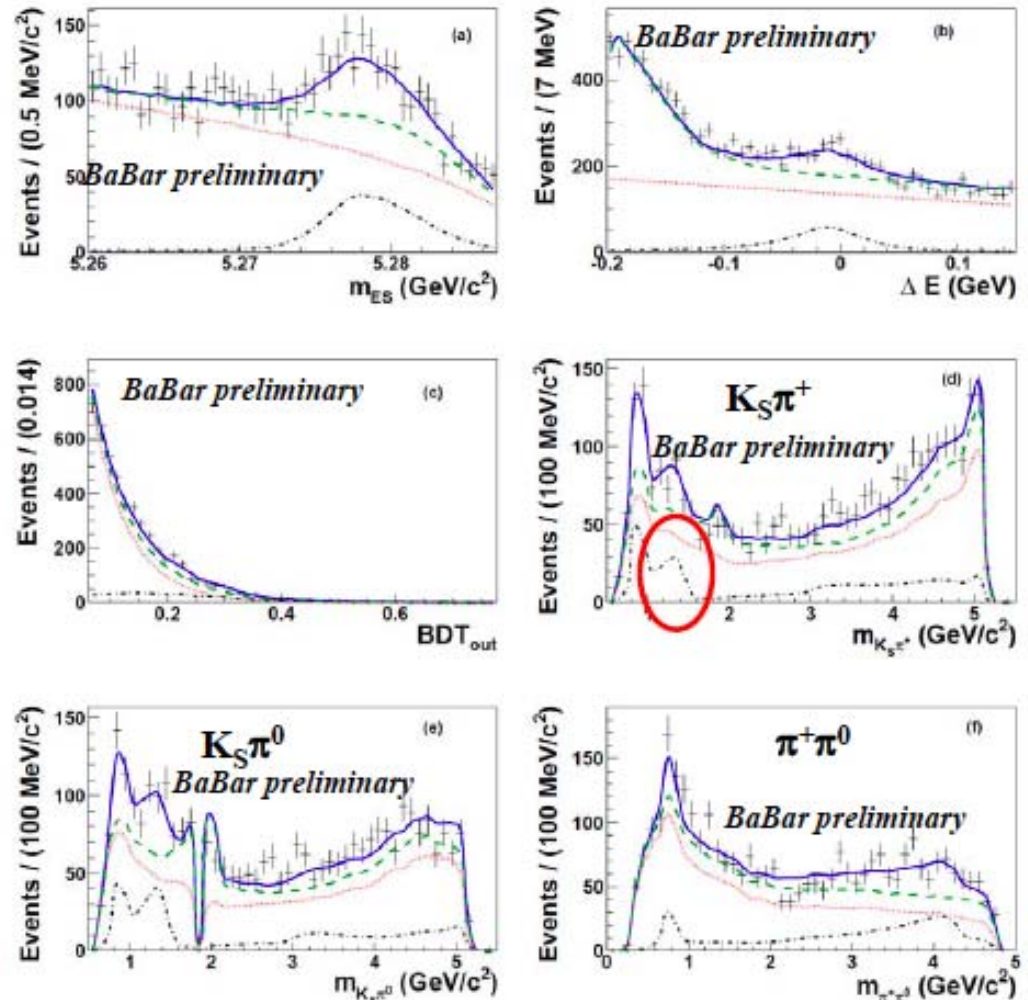
- About 32,000 candidates in the fitted sample

Best fit:  $1,014 \pm 63$  (stat.) signal events

- First measurement of inclusive  $K^0\pi^+\pi^0$  branching fraction:  
 $(45.9 \pm 2.6 \pm 3.0 \pm 8.6) \times 10^{-6}$
- First measurement of the  $K^{*+}_0(1430)\pi^0$  branching fraction:  
 $(17.2 \pm 2.4 \pm 1.5 \pm 1.8) \times 10^{-6}$  ( $5.4\sigma$ )
- All significant branching fractions

Decay channel	$\mathcal{B}$ ( $10^{-6}$ )
$K^0\pi^+\pi^0$	$45.9 \pm 2.6 \pm 3.0 \pm 8.6$
<i>BaBar preliminary</i>	
$K^{*0}(892)\pi^+$	$14.6 \pm 2.4 \pm 1.4 \pm 0.5$
$K^{*+}(892)\pi^0$	$9.2 \pm 1.3 \pm 0.6 \pm 0.5$
$K^{*0}_0(1430)\pi^+$	$50.0 \pm 4.8 \pm 6.1 \pm 4.0$
$K^{*+}_0(1430)\pi^0$	$17.2 \pm 2.4 \pm 1.5 \pm 1.8$
$\rho^+(770)K^0$	$9.4 \pm 1.6 \pm 1.1 \pm 2.6$

→  $\Sigma(\text{resonance BFs}) < \text{inclusive BF}$   
 due to destructive interferences



Signal-enhanced distributions

# Simultaneous fit of the separate $B^+$ and $B^-$ data samples

- Compute the **inclusive & exclusive CP asymmetries** ( $A_{CP}$ )
- **Parameterization** of the coefficients for the 2-body resonant decay modes

$$c_j = (x_j + \Delta x_j) + i(y_j + \Delta y_j),$$

$$\bar{c}_j = (x_j - \Delta x_j) + i(y_j - \Delta y_j),$$

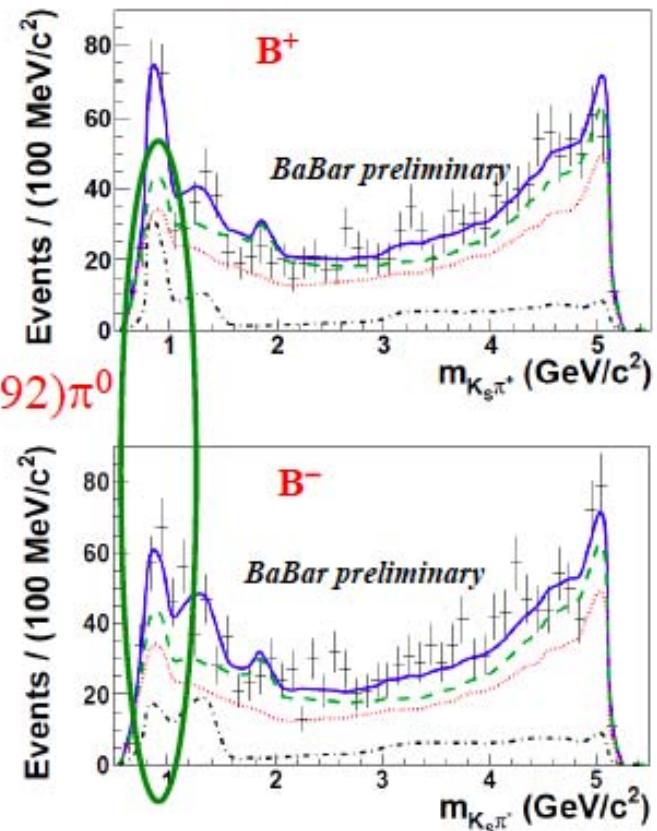
→  $A_{CP}$  computation

$$A_{CP,j} = \frac{|\bar{c}_j|^2 - |c_j|^2}{|\bar{c}_j|^2 + |c_j|^2} = -\frac{2(x_j\Delta x_j + y_j\Delta y_j)}{x_j^2 + \Delta x_j^2 + y_j^2 + \Delta y_j^2}.$$

- Results: **first evidence of direct CPV for  $K^{*+}(892)\pi^0$**

Decay channel	$A_{CP}$
$K^0\pi^+\pi^0$	$0.07 \pm 0.05 \pm 0.03 \pm 0.04$
<i>BaBar preliminary</i>	
$K^{*0}(892)\pi^+$	$-0.12 \pm 0.21 \pm 0.08 \pm 0.11$
$K^{*+}(892)\pi^0$	$-0.52 \pm 0.14 \pm 0.04 \pm 0.04$
$K_0^{*0}(1430)\pi^+$	$0.14 \pm 0.10 \pm 0.04 \pm 0.14$
$K_0^{*+}(1430)\pi^0$	$0.26 \pm 0.12 \pm 0.08 \pm 0.12$
$\rho^+(770)K^0$	$0.21 \pm 0.19 \pm 0.07 \pm 0.30$

(3.4 $\sigma$ )



# Outlook

- First observation of the charmless decay  $B^+ \rightarrow K_S \pi^+ \pi^0$
- First evidence ( $3.4\sigma$ ) of direct CP violation for  $B^+ \rightarrow K^{*+}(892)\pi^0$

- $\Delta A_{\text{CP}}(K^* \pi) = \underbrace{A_{\text{CP}}(K^{*+}\pi^0)}_{\text{This new result}} - \underbrace{A_{\text{CP}}(K^{*+}\pi^-)}_{\text{HFAG average } -0.23 \pm 0.06} = -0.29 \pm 0.16$

→ To be compared with  $\Delta A_{\text{CP}}(K\pi) = 0.122 \pm 0.022$

- Results consistent with previous measurements (when available)
- No deviation from Standard Model expectations
- More statistics needed to study further these decays  
→ Belle-2
- BaBar article to be submitted to PRD soon



# Belle charm mixing and CPV results at ICHEP 2014

- First observation of  $D^0$ - $\bar{D}^0$  mixing in  $e^+e^-$  collision in the measurement of time-dependent ratio of WS to RS decay rates

$$x'^2 = (0.09 \pm 0.22) \times 10^{-3} \quad y' = (4.6 \pm 3.4) \times 10^{-3}$$

⇒ no mixing hypothesis is excluded at  $5.1\sigma$  level

- Updated measurement of  $D^0$ - $\bar{D}^0$  mixing in  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$

$$x = (0.56 \pm 0.19)\% \quad y = (0.30 \pm 0.15)\%$$

⇒ significance of mixing is estimated to be  $2.5\sigma$

⇒ No evidence for CP violation in the decay

- Significantly improved measurement of time-integrated CP violating asymmetry  $A_{CP}$  in  $D^0 \rightarrow \pi^0 \pi^0$  and the result is consistent with no CPV

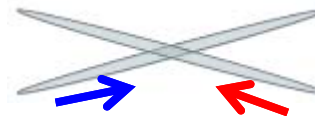
$$A_{CP}(\pi^0 \pi^0) = (-0.03 \pm 0.64 \pm 0.10)\%$$

⇒ updated the existing measurement of CP asymmetry in  $D^0 \rightarrow K_S^0 \pi^0$

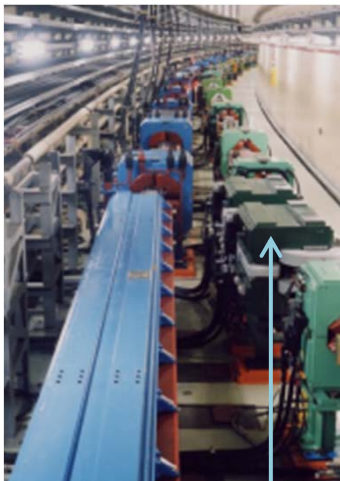
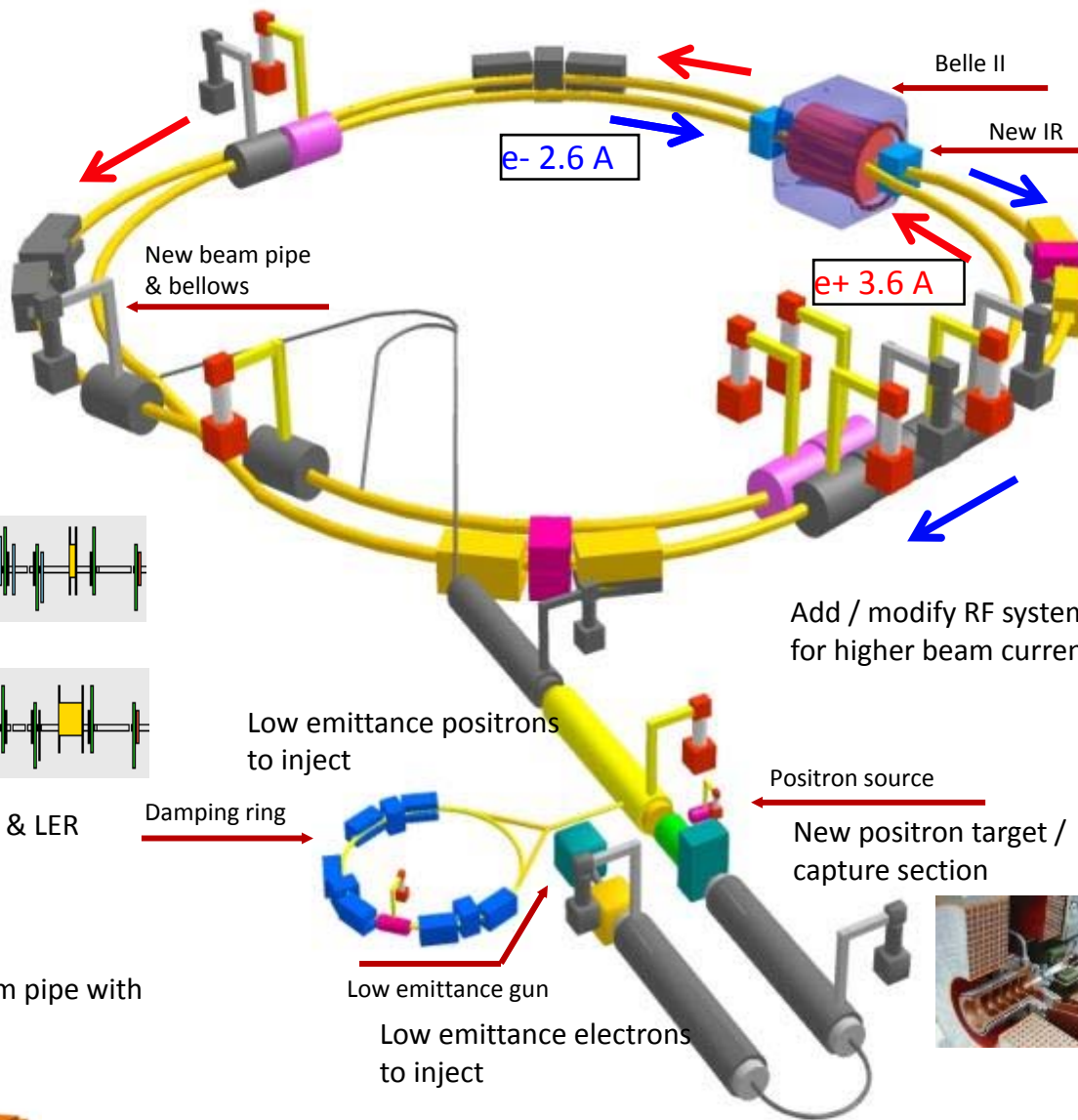
# KEKB to SuperKEKB



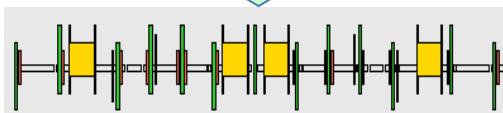
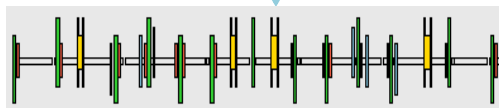
Colliding bunches



New superconducting / permanent final focusing quads near the IP

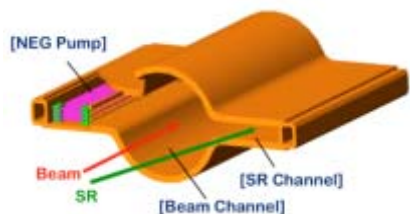


Replace short dipoles with longer ones (LER)



Redesign the lattices of HER & LER to squeeze the emittance

TiN-coated beam pipe with antechambers



Add / modify RF systems for higher beam current



**To obtain x40 higher luminosity**