

2 **Studies on τ decays at Belle II**

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7 **ABSTRACT:** Tau leptons offer a very clean environment to study the process of hadronization and
8 are powerful tools to probe physics beyond the Standard Model (SM), which might be enhanced due
9 to mass-dependent couplings. The Belle II experiment, installed at SuperKEKB asymmetric energy
10 electron-positron collider, will collect the world's largest sample of tau pair events $e^+e^- \rightarrow \tau^+\tau^-$.
11 Direct searches for new invisible mediators, charged lepton flavor violation in τ decays, and tests of
12 the SM via precision measurements of τ lepton properties and couplings are reported in the following
13 article. The results presented here rely on the data collected by Belle II during 2019-2021.

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16 **Contents**

17	1 Introduction	1
18	2 Leptons as discovery tools: the experimental challenges	1
19	3 Experimental facility: Belle II	2
20	3.1 Typical τ signatures in e^+e^- collisions	2
21	4 Searches for a new invisible boson α in τ decays	2
22	5 Direct searches for LFV $\tau \rightarrow \ell\phi$ decays	3
23	6 Measurement of the τ lepton mass	4
24	7 Prospects on lepton flavor universality tests	4

25 **1 Introduction**

26 As the only leptons massive enough to decay into hadrons, τ s not only allow to investigate the
27 hadronization mechanism via their hadronic final states, but might preferentially couple to non-SM
28 physics, through mass-dependent couplings. They are challenging from the experimental point of
29 view, since they can not be detected as long-lived particles, but instead reconstructed from their
30 final-state products, which involve undetectable neutrinos. Furthermore, they allow searching for
31 charged lepton flavor violation (LFV), which would provide an indisputable proof for beyond SM
32 physics. Processes involving LFV can occur in the SM via weak interaction charged currents,
33 due to neutrino oscillations, and are predicted at the level of 10^{-50} , which is beyond the reach
34 of current and future experiments. Belle II has a unique capability to probe both new invisible
35 mediators and LFV in τ decays. Moreover, it can look for indirect signs of non-SM physics through
36 high precision measurements of SM fundamental parameters. We report searches for new invisible
37 particles, τ LFV decays and the measurement of the τ lepton mass using the data collected by
38 the Belle II detector [1] at the SuperKEKB asymmetric energy e^+e^- collider [2]. SuperKEKB
39 mainly operates at a centre-of-mass energy (c.m.) of 10.58 GeV and adopts a nano-beam scheme to
40 reach unprecedented instantaneous luminosity. At the time of this conference, the accelerator had
41 achieved the peak luminosity world's record of $4.7 \times 10^{34} \text{ cm}^{-2}/\text{s}$ and Belle II has so far collected
42 424 fb^{-1} of data, including unique energy scan samples. It is currently in its first long shutdown.

43 **2 Leptons as discovery tools: the experimental challenges**

44 Leptonic production of tau pair processes $e^+e^- \rightarrow \tau^+\tau^-$ provide a very clean physics environment
45 and can rely on precise QED predictions to look for physics beyond the SM. The way is two-fold:

46 one could look for deviations from SM predictions in high precision measurements of very cleaned
 47 and precisely computed observables; a second possibility is instead to search for processes that
 48 would be either forbidden or highly suppressed in the SM and whose observation is *per se* a hint of
 49 new physics. The first class of measurements are mainly systematically limited and to improve the
 50 current results and attain the world's best precision, an excellent understanding of the experiment
 51 performance at the fraction of permille level is required. On the other hand, measurements of rare
 52 or forbidden processes imply to achieve unprecedented luminosity to collect sufficiently large data
 53 sets and devise new analysis techniques to boost signal efficiencies in order to reach sensitivities
 54 below the 10^{-8} level.

55 **3 Experimental facility: Belle II**

56 The Belle II detector is a multipurpose spectrometer surrounding the interaction point and providing
 57 coverage of more than 90% of the solid angle. The details of the Belle II detector can be found
 58 elsewhere [1]. Belle II ensures a very high reconstruction efficiency for neutral particles and
 59 excellent resolutions despite the harsh beam background environment, which are crucial when
 60 dealing with recoiling system and missing-energy final states. Additionally, it is equipped with
 61 dedicated low-multiplicity trigger lines at hardware level, mainly based on calorimetric information,
 62 which were not available at Belle. Profiting from the well known initial state of e^+e^- collisions,
 63 and its near-hermetic coverage, Belle II has a unique capability to probe signatures involving
 64 invisible final states and particles escaping detection. Moreover, the production cross-section for
 65 $e^+e^- \rightarrow \tau^+\tau^-$ events is 0.919 nb at a c.m. energy $\sqrt{s} = 10.58$ GeV, allowing Belle II to make
 66 precision measurements of τ lepton properties.

67 **3.1 Typical τ signatures in e^+e^- collisions**

68 In $e^+e^- \rightarrow \tau^+\tau^-$ processes, tau candidates are produced back-to-back in c.m. system. Their decay
 69 products are well separated into two opposite hemispheres, defined by the plane perpendicular to
 70 the thrust axis $\hat{\mathbf{n}}_T$, which is the vector maximizing the quantity

$$T = \max_{\hat{\mathbf{n}}_T} \left(\frac{\sum_i |\mathbf{p}_i \cdot \hat{\mathbf{n}}_T|}{\sum_i |\mathbf{p}_i|} \right), \quad (3.1)$$

71 where \mathbf{p}_i is the momentum of the final state particle i , including both charged and neutral particles.
 72 According to the number of charged particles in each hemisphere, and consistently with charge
 73 conservation in τ decays, two main topologies can be selected: the 3×1 -prong decays, as schemat-
 74 ically shown in the left drawing of Figure 1, with three charged particles on one side and only one
 75 in the opposite hemisphere; or the 1×1 -prong decays. Requiring a fixed number of tracks matching
 76 one of these topology classes is a powerful way to suppress the main background from continuum
 77 $e^+e^- \rightarrow q\bar{q}$ processes and enhance signal purity when reconstructing $e^+e^- \rightarrow \tau^+\tau^-$ events.

78 **4 Searches for a new invisible boson α in τ decays**

79 Decays of τ leptons to new LFV bosons are postulated in many models. The process searched for
 80 in this study is $e^+e^- \rightarrow \tau^+(\rightarrow \ell\alpha)\tau^-(\rightarrow \pi^+\pi^-\pi^-\nu)$ and its charged conjugated, where the first τ

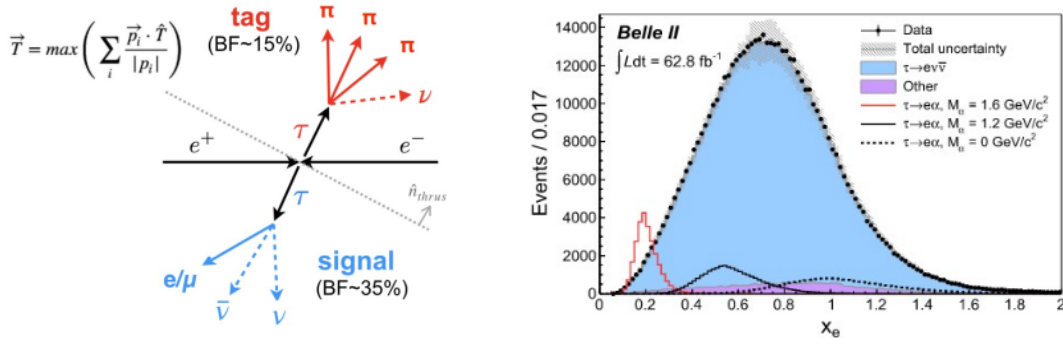


Figure 1. On the left, a typical 3x1-prong decay in $e^+e^- \rightarrow \tau^+\tau^-$ events, where one tau decays into three charged particles and the other one into one charged particle, is depicted. On the right, the distribution of the normalized lepton energy x_ℓ for the electron channel in the search for $\tau \rightarrow \ell\alpha$ is shown. Data are the black dots and the simulation is the stack filled histograms.

81 is defined as the signal and the second one as the tag. The signal τ is searched for in its decay to a
 82 new invisible boson α , accompanied by a lepton $\ell = e, \mu$, therefore 3×1 -prong events are selected
 83 among the reconstructed ones. The signal τ rest-frame is approximated using as energy half the
 84 collision energy $\sqrt{s}/2$ and as momentum direction the opposite to the one of the reconstructed tag
 85 τ . We exploit the kinematic features of the signal process as a two-body decay to discriminate
 86 from the background. Variable shape differences are more prominent in the c.m. frame, where we
 87 look for a narrow peak in the distribution of the normalized lepton energy, reported in the right
 88 plot of Figure 1, over a smooth contribution coming from the irreducible background of $\tau \rightarrow \ell\nu\bar{\nu}_\ell$
 89 processes. In absence of any signal excess in 63 fb^{-1} data, 95% CL upper limits are computed on the
 90 ratio of branching fractions $\mathcal{B}(\tau \rightarrow \ell\alpha)$ normalized to $\mathcal{B}(\tau \rightarrow \ell\nu\bar{\nu}_\ell)$ [3]. This analysis provides
 91 limits between 2-14 times more stringent than the previous one set by ARGUS [4].

92 5 Direct searches for LFV $\tau \rightarrow \ell\phi$ decays

93 Possible new mediators may enhance the branching fraction for τ LFV decays $\tau^- \rightarrow \ell^-\phi$ up
 94 to observable levels of $10^{-11} - 10^{-8}$, and accommodate for flavor anomalies observed in lepton
 95 flavor universality tests with B decays [5]. In contrast to previous searches for $\tau^- \rightarrow \ell^-\phi$ decays
 96 performed at Belle [6] on $e^+e^- \rightarrow \tau^+\tau^-$ events, we apply for the first time an *untagged* approach.
 97 Only the signal τ decay to a ϕ meson candidate and a lepton, either muon or electron, is explicitly
 98 reconstructed and the other τ is not required to decay to any specific known final state. Event
 99 kinematics features and signal properties are used in a BDT classifier to suppress the background,
 100 with double the final signal efficiency for the muon mode with respect to previous analyses.
 101 Yields are extracted with a Poisson counting experiment approach from windows peaking at the
 102 known τ mass and at zero in the 2D plane of $(M_\tau, \Delta E_\tau)$, respectively, with ΔE_τ the difference
 103 between the reconstructed energy of the signal τ in the c.m. frame and half the collision energy.
 104 We find no significant excess and set 90% CL upper limits on the branching fractions to be
 105 $\mathcal{B}_{\text{UL}}(\tau \rightarrow e\phi) = 23 \times 10^{-8}$ and $\mathcal{B}_{\text{UL}}(\tau \rightarrow \mu\phi) = 9.7 \times 10^{-8}$ [7].

106 6 Measurement of the τ lepton mass

107 Lepton properties are fundamental parameters of the SM and need to be measured with the highest
 108 precision. Belle II is suitable to access several τ lepton properties. By applying the pseudo-mass
 109 M_{min} technique to reconstructed $e^+e^- \rightarrow \tau^+\tau^-$ events from 190 fb^{-1} data, we provide the world's
 110 most precise measurement of the τ mass M_τ . The measured value is extracted from a fit to
 111 the endpoint of the distribution $M_{min} = \sqrt{M_{3\pi}^2 + 2(\sqrt{s}/2 - E_{3\pi}^*)(E_{3\pi}^* - P_{3\pi}^*)}$, which is computed
 112 from events where the signal τ is reconstructed in its decays to three charged pions and the other τ
 113 decaying into one charged particle. The distributions of the pseudomass in simulation and data is
 114 shown in the left plot of Figure 2. An excellent control of the systematic sources, dominated by the
 115 calibration of the beam energies and the charged-particle momenta scale, is required to reduce the
 116 total systematic uncertainty to $0.11 \text{ MeV}/c^2$, achieving the most precise measurement to date of the
 τ lepton mass of $1777.09 \pm 0.08_{\text{stat}} \pm 0.11_{\text{sys}}$ [8].

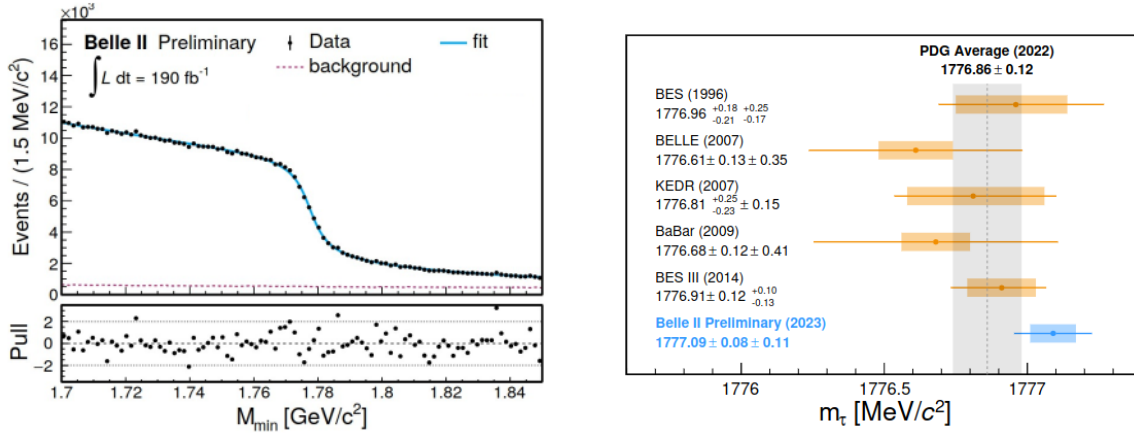


Figure 2. On the left, the spectrum of the reconstructed pseudomass in data (black dots) and the superimposed fit (solid blue line) are shown. The bottom inset plot displays differences between data and fit result divided by the statistical uncertainties. On the right, the summary of the most precise measurements of the tau mass to date, compared to the world average (gray band) and this work result (blue text).

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118 7 Prospects on lepton flavor universality tests

119 Lepton flavor universality (LFU) in SM assumes all three leptons have equal coupling strength to
 120 the charged gauge bosons of the electroweak interaction. Many models predict new forces violating
 121 LFU, that could for example imply a new singly-charged scalar singlet [9]. Tau decays allow high
 122 precision tests of LFU by measuring the ratio of the branching fractions of τ decays to muon and to
 123 electron,

$$R_\mu = \frac{\mathcal{B}(\tau \rightarrow \mu \nu_\mu \nu_\tau)}{\mathcal{B}(\tau \rightarrow e \nu_e \nu_\tau)} \quad (7.1)$$

124 The most precise result to date is $R_\mu = 0.9796 \pm 0.0016_{\text{stat}} \pm 0.0036_{\text{sys}}$ provided by Babar [10]. It uses
 125 467 fb^{-1} collision data, for a final 0.4% precision, systematically dominated by the contribution of
 126 the particle identification (PID) and trigger uncertainties. Simulation studies at Belle II show room

127 for a three-fold improvement: by devising dedicated low multiplicity triggers based on calorimeter
 128 information, which provide a better understanding of the kinematic dependency and reduce the
 129 associated systematic uncertainty; by dropping the likelihood-based PID selector for pions and
 130 deploying BDT classifier for lepton identification, which will decrease the probability to wrongly
 131 identify a pion as a lepton to less than 0.1%; eventually, adding the 1×1 -prong decays as signal
 132 signature to increase the size of the analyzed data set. Further studies for the development of the
 133 specific 1×1 topology triggers are still needed, but already with one quarter of Babar data set, Belle
 134 II expected sensitivity achieves the same statistical precision of 0.16%.

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