

## 2 **Studies on $\tau$ decays at Belle II**

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3 **L. Zani**

4 *INFN of Roma Tre,*

5 *Via della Vasca Navale, 84 Roma, Italy*

6 *E-mail:* [laura.zani@roma3.infn.it](mailto:laura.zani@roma3.infn.it)

7 **ABSTRACT:** Tau leptons are powerful tools to probe physics beyond the Standard Model (SM), which  
8 might be enhanced if the couplings are mass-dependent, thanks to the larger mass of  $\tau$ . The Belle  
9 II experiment, as successor of Belle, is installed at the SuperKEKB asymmetric energy electron-  
10 positron collider and aims at collecting the world's largest sample of tau pair events  $e^+e^- \rightarrow \tau^+\tau^-$ .  
11 Direct searches for new invisible mediators, charged lepton flavor violation in  $\tau$  decays, and tests  
12 of the SM via precision measurements of  $\tau$  lepton properties and couplings are reported in the  
13 following article. The results presented here are based on the data collected by Belle II during  
14 2019-2021.

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26 **1 Introduction**

27 As the only leptons massive enough to decay into hadrons,  $\tau$ s not only allow to investigate the  
28 hadronization mechanism via their hadronic final states, but might preferentially couple to non-SM  
29 physics, through mass-dependent couplings. From the experimental point of view though, taus are  
30 challenging since they can not be detected as long-lived particles, but must instead be reconstructed  
31 from their final-state products, which involve undetectable neutrinos. Furthermore, they allow  
32 searching for charged lepton flavor violation (LFV), which would provide an indisputable proof for  
33 beyond SM physics. Processes involving LFV can occur in the SM via weak interaction charged  
34 currents, due to neutrino oscillations, and are predicted at the level of  $10^{-50}$ , which is beyond the  
35 reach of current and future experiments. Belle II has a unique capability to probe both new invisible  
36 mediators and LFV in  $\tau$  decays. Moreover, it can look for indirect signs of non-SM physics through  
37 high precision measurements of SM fundamental parameters. We report searches for new invisible  
38 particles,  $\tau$  LFV decays and the measurement of the  $\tau$  lepton mass using the data collected by  
39 the Belle II detector [1] at the SuperKEKB asymmetric energy  $e^+e^-$  collider [2]. SuperKEKB  
40 mainly operates at a centre-of-mass energy (c.m.) of 10.58 GeV and adopts a nano-beam scheme  
41 to reach unprecedented instantaneous luminosity. At the time of this conference, the accelerator  
42 has achieved the peak luminosity world record of  $4.7 \times 10^{34} \text{ cm}^{-2}/\text{s}$  and Belle II has so far collected  
43  $424 \text{ fb}^{-1}$  of data, including unique energy scan samples. It is currently in its first long shutdown.

44 **2 Leptons as discovery tools: the experimental challenges**

45 Leptonic production of tau pair processes  $e^+e^- \rightarrow \tau^+\tau^-$  provide a very clean physics environment  
46 and can rely on precise QED predictions to look for physics beyond the SM. The way is two-fold:

one could look for deviations from SM predictions in high precision measurements of very clean and precisely computed observables; a second possibility is instead to search for processes that would be either forbidden or highly suppressed in the SM and whose observation is *per se* a hint of new physics. The first class of measurements are mainly systematically limited and to improve the current results and attain the world's best precision, an excellent understanding of the experiment performance at the fraction of permille level is required. On the other hand, measurements of rare or forbidden processes imply to achieve unprecedented luminosity to collect sufficiently large data sets and devise new analysis techniques to boost signal efficiencies in order to reach sensitivities below the  $10^{-8}$  level.

### 3 Experimental facility: Belle II

The Belle II detector is a multipurpose spectrometer surrounding the  $e^+e^-$  interaction point and providing coverage of more than 90% of the solid angle. The details of the Belle II detector can be found elsewhere [1]. Belle II ensures a very high reconstruction efficiency for neutral particles and excellent resolutions despite the harsh beam background environment, both of which are crucial when dealing with recoiling system and missing-energy final states. Additionally, it is equipped with dedicated low-multiplicity trigger lines at hardware level, mainly based on calorimetric information, which were not available at Belle. Profiting from the well known initial state of  $e^+e^-$  collisions, and its near-hermetic coverage, Belle II has a unique capability to probe signatures involving invisible final states and particles escaping detection. Moreover, the production cross-section for  $e^+e^- \rightarrow \tau^+\tau^-$  events is 0.919 nb at a c.m. energy  $\sqrt{s} = 10.58$  GeV, allowing Belle II to collect large data samples for precision measurements of  $\tau$  lepton properties.

#### 3.1 Typical $\tau$ signatures in $e^+e^-$ collisions

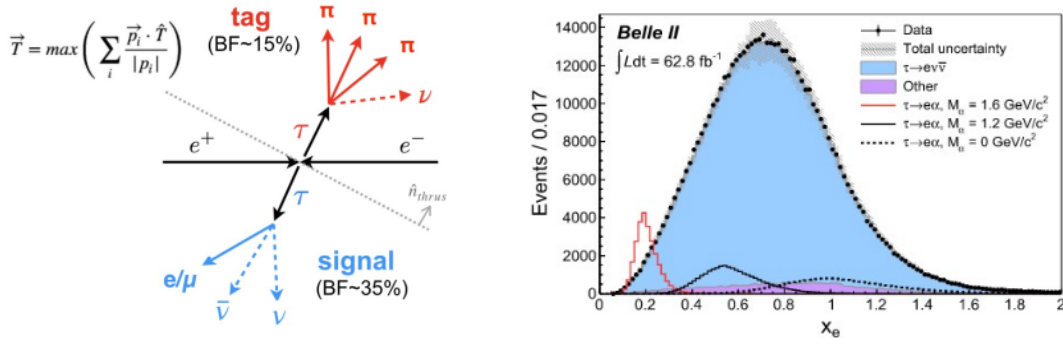
In  $e^+e^- \rightarrow \tau^+\tau^-$  processes, tau candidates are produced back-to-back in c.m. system. Their decay products are well separated into two opposite hemispheres, defined by the plane perpendicular to the thrust axis  $\hat{\mathbf{n}}_T$ , which is the vector maximizing the quantity

$$T = \max_{\hat{\mathbf{n}}_T} \left( \frac{\sum_i |\mathbf{p}_i \cdot \hat{\mathbf{n}}_T|}{\sum_i |\mathbf{p}_i|} \right), \quad (3.1)$$

where  $\mathbf{p}_i$  is the momentum of the final state particle  $i$ , including both charged and neutral particles. According to the number of charged particles in each hemisphere, consistently with charge conservation in  $\tau$  decays, two main topologies can be selected: the  $3 \times 1$ -prong decays, as schematically shown in the left drawing of Figure 1, with three charged particles on one side and only one in the opposite hemisphere; or the  $1 \times 1$ -prong decays. Requiring the reconstructed tracks to match one of these topology classes is a powerful way to suppress the main background from continuum  $e^+e^- \rightarrow q\bar{q}$  processes and enhance signal purity when reconstructing  $e^+e^- \rightarrow \tau^+\tau^-$  events.

### 4 Searches for a new invisible boson $\alpha$ in $\tau$ decays

Decays of  $\tau$  leptons to new LFV bosons are postulated in many models [3]. The process searched for in this study is  $e^+e^- \rightarrow \tau^+(\rightarrow \ell\alpha)\tau^-(\rightarrow \pi^+\pi^-\pi^-\nu)$  and its charge conjugated, where the first  $\tau$



**Figure 1.** On the left, a typical 3x1-prong decay in  $e^+e^- \rightarrow \tau^+\tau^-$  events, where one tau decays into three charged particles and the other one into one charged particle, is depicted. On the right, the distribution of the normalized lepton energy  $x_\ell$  for the electron channel in the search for  $\tau \rightarrow \ell\alpha$  is shown. Data are the black dots and the simulation is the stack filled histograms.

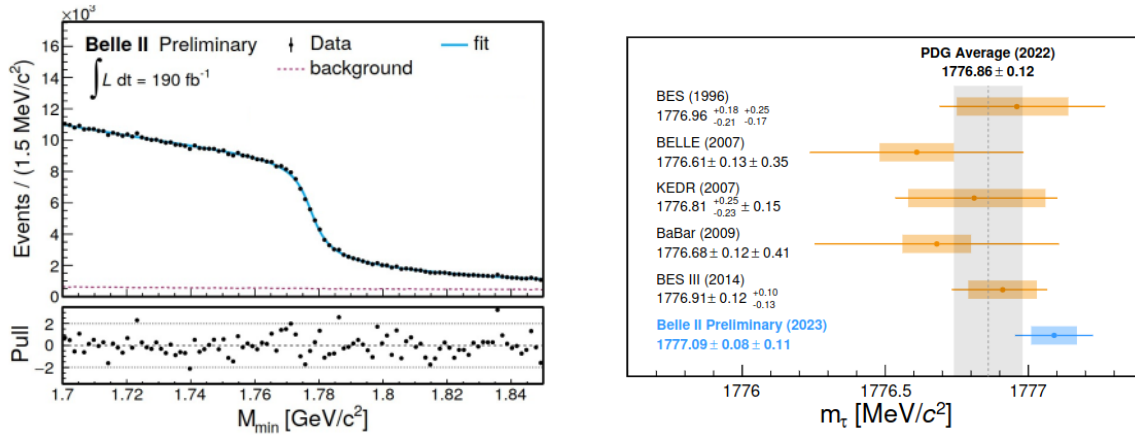
82 is defined as the signal and the second one as the tag. The signal  $\tau$  is searched for in its decay to a  
 83 new invisible boson  $\alpha$ , accompanied by a lepton  $\ell = e, \mu$ , therefore  $3 \times 1$ -prong events are selected.  
 84 The signal  $\tau$  rest-frame is approximated using as energy half the collision energy  $\sqrt{s}/2$  and as  
 85 momentum direction the opposite to the one of the reconstructed tag  $\tau$ . We exploit the kinematic  
 86 features of the signal process as a two-body decay to discriminate it from the background, by looking  
 87 for a narrow peak in the distribution of the normalized lepton energy in the c.m. frame, reported  
 88 in the right plot of Figure 1, over a smooth contribution coming from the irreducible background  
 89 of  $\tau \rightarrow \ell\nu\bar{\nu}_\ell$  processes. In absence of any signal excess in  $63 \text{ fb}^{-1}$  data, 95% CL upper limits are  
 90 computed on the ratio of branching fractions  $\mathcal{B}(\tau \rightarrow \ell\alpha)$  normalized to  $\mathcal{B}(\tau \rightarrow \ell\nu\bar{\nu}_\ell)$  [4]. This  
 91 analysis provides limits between 2-14 times more stringent than the previous one set by ARGUS [5].

## 92 5 Direct searches for LFV $\tau \rightarrow \ell\phi$ decays

93 Possible new mediators may enhance the branching fraction for  $\tau$  LFV decays  $\tau^- \rightarrow \ell^-\phi$  up  
 94 to observable levels of  $10^{-11} - 10^{-8}$ , and accommodate for flavor anomalies observed in lepton  
 95 flavor universality tests with  $B$  decays [6]. In contrast to previous searches for  $\tau^- \rightarrow \ell^-\phi$  decays  
 96 performed at Belle [7] on  $e^+e^- \rightarrow \tau^+\tau^-$  events, we apply for the first time an *untagged* approach.  
 97 Only the signal  $\tau$  decay to a  $\phi$  meson candidate and a lepton, either muon or electron, is explicitly  
 98 reconstructed and the other  $\tau$  is not required to decay to any specific known final state. Event  
 99 kinematics features and signal properties are used in a BDT classifier to suppress the background,  
 100 with twice the signal efficiency for the muon mode with respect to previous analyses. Yields are  
 101 extracted with a Poisson counting experiment approach from windows peaking at the known  $\tau$  mass  
 102 and at zero in the 2D plane of  $(M_\tau, \Delta E_\tau)$ , respectively, where  $\Delta E_\tau$  is the difference between the  
 103 reconstructed energy of the signal  $\tau$  in the c.m. frame and half the collision energy. We find no  
 104 significant excess and set 90% CL upper limits on the branching fractions of  $\mathcal{B}_{\text{UL}}(\tau \rightarrow e\phi) =$   
 105  $23 \times 10^{-8}$  and  $\mathcal{B}_{\text{UL}}(\tau \rightarrow \mu\phi) = 9.7 \times 10^{-8}$  [8].

## 106 6 Measurement of the $\tau$ lepton mass

107 Lepton properties are fundamental parameters of the SM and need to be measured with the highest  
 108 precision. Belle II is suitable to measure several  $\tau$  lepton properties. By applying the pseudo-  
 109 mass  $M_{min}$  technique to reconstructed  $e^+e^- \rightarrow \tau^+\tau^-$  events from  $190 \text{ fb}^{-1}$  data, we provide the  
 110 world's most precise measurement of the  $\tau$  mass  $M_\tau$ . The measured value is extracted from a fit to  
 111 the endpoint of the distribution  $M_{min} = \sqrt{M_{3\pi}^2 + 2(\sqrt{s}/2 - E_{3\pi}^*)(E_{3\pi}^* - P_{3\pi}^*)}$ , which is computed  
 112 from events where the signal  $\tau$  is reconstructed in its decays to three charged pions and the other  $\tau$   
 113 decaying into one charged particle. The distributions of the pseudomass in simulation and data is  
 114 shown in the left plot of Figure 2. An excellent control of the systematic sources, dominated by the  
 115 calibration of the beam energies and the charged-particle momentum scale, is required to reduce  
 116 the total systematic uncertainty to  $0.11 \text{ MeV}/c^2$ , achieving the most precise measurement to date of  
 the  $\tau$  lepton mass of  $1777.09 \pm 0.08_{\text{stat}} \pm 0.11_{\text{sys}}$  [9].



**Figure 2.** On the left, the spectrum of the reconstructed pseudomass in data (black dots) and the superimposed fit (solid blue line) are shown. The bottom inset plot displays differences between data and fit result divided by the statistical uncertainties. On the right, the summary of the most precise measurements of the tau mass to date, compared to the world average (gray band) and this work result (blue text).

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## 118 7 Prospects on lepton flavor universality tests

119 Lepton flavor universality (LFU) in SM assumes all three leptons have equal coupling strength to  
 120 the charged gauge bosons of the electroweak interaction. Many models predict new forces violating  
 121 LFU, that could for example imply a new singly-charged scalar singlet [10]. Tau decays allow high  
 122 precision tests of LFU by measuring the ratio of the branching fractions of  $\tau$  decays to muon and to  
 123 electron,

$$R_\mu = \frac{\mathcal{B}(\tau \rightarrow \mu \nu_\mu \nu_\tau)}{\mathcal{B}(\tau \rightarrow e \nu_e \nu_\tau)} \quad (7.1)$$

124 The most precise result to date is  $R_\mu = 0.9796 \pm 0.0016_{\text{stat}} \pm 0.0036_{\text{sys}}$  provided by Babar [11]. It uses  
 125  $467 \text{ fb}^{-1}$  collision data, for a final 0.4% precision, systematically dominated by the contribution of  
 126 the particle identification (PID) and trigger uncertainties. Simulation studies at Belle II show room

127 for a several improvements: by devising dedicated low multiplicity triggers based on calorimeter  
128 information, which provide a better understanding of the kinematic dependency and reduce the  
129 associated systematic uncertainty; by dropping the likelihood-based PID selector for pions and  
130 deploying BDT classifier for lepton identification, which will decrease the probability to wrongly  
131 identify a pion as a lepton to less than 0.1%; eventually, adding the  $1 \times 1$ -prong decays as signal  
132 signature to increase the size of the analyzed data set. Further studies for the development of the  
133 specific  $1 \times 1$  topology triggers are still needed, but already with one quarter of Babar data set, Belle  
134 II expected sensitivity achieves the same statistical precision of 0.16%.

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